

# Small Scale Structure in Vector Dark Matter Mustafa A. Amin RICE









## distinct phenomenology in (ultra)light vector dark matter

 new class of polarized vector solitons with macroscopic spin



Zhang, Jain MA (2021) Adshead & Lozanov (2021) Jain and MA (2021)

interference patterns, and halo density profiles 

MA, Jain, Karur & Mocz (2022)



## distinct phenomenology in (ultra)light vector dark matter

 new class of polarized vector solitons with macroscopic spin



Adshead & Lozanov (2021) Jain and MA (2021) Zhang, Jain MA (2021)

• interference patterns, and halo density profiles

MA, Jain, Karur & Mocz (2022)

files (2022)

MA, Jain, Karur & Mocz (2022)



### formation mechanisms?



Gorghetto, Hardy, March-Russell, Song & West (2022)

#### misaligned scalar + self-interactions?



#### gravitational in late universe







### motivation & introduction

#### dark matter mass ?



image credit: E. Ferreira



### dark matter spin?





### light, bosonic wave dark matter



$$\lambda_{\rm dm}^3 \lambda_{\rm dB}^3 \sim 10^{23} \left(\frac{10^{-5} \,\mathrm{eV}}{m}\right)^4 \sim 10^{83} \left(\frac{10^{-20} \,\mathrm{eV}}{m}\right)^4$$

$$_{\rm dB} \sim 10^3 {\rm cm} \times \left(\frac{10^{-5} {\rm eV}}{m}\right) \left(\frac{10^{-3} c}{v}\right) \sim 1 {\rm pc} \left(\frac{10^{-20} {\rm eV}}{m}\right) \left(\frac{10^{-3} c}{v}\right)$$





#### vector dark matter

 $S = \int d^4x \sqrt{-g} \left[ -\frac{1}{4} g^{\mu\alpha} g^{\nu\beta} \mathcal{G}_{\mu\nu} \mathcal{G}_{\alpha\beta} + \frac{1}{2} \frac{m^2 c^2}{\hbar^2} g^{\mu\nu} W_{\mu} W_{\nu} + \frac{c^3}{16\pi G} R + \dots \right]$ 

 $\mathcal{G}_{\mu\nu} = \partial_{\mu}W_{\nu} - \partial_{\nu}W_{\mu}$ 

$$S = \int d^4x \sqrt{-g} \left[ -\frac{1}{2} \partial_\mu \phi \partial^\mu \phi + \frac{1}{2} \frac{m^2 c^2}{\hbar^2} \right]$$



#### vector case



#### scalar case



#### non-relativistic limit = multicomponent Schrödinger-Poisson

$$S = \int \mathrm{d}^4 x \, \sqrt{-g} \Big[ -\frac{1}{4} g^{\mu\alpha} g^{\nu\beta} \, \mathcal{G}_{\mu\nu} \mathcal{G}_{\alpha\beta} + \frac{1}{2} \frac{n}{2} \Big]$$

$$\mathcal{G}_{\mu\nu} = \partial_{\mu}W_{\nu} - \partial_{\nu}W_{\mu}$$

#### non-relativistic limit

$$\boldsymbol{W}(t,\boldsymbol{x}) \equiv \frac{\hbar}{\sqrt{2mc}} \Re \left[ \boldsymbol{\Psi}(t,\boldsymbol{x}) e^{-imc^2 t/\hbar} \right]$$

$$\mathcal{S}_{nr} = \int \mathrm{d}t \mathrm{d}^3 x \left[ \frac{i\hbar}{2} \mathbf{\Psi}^{\dagger} \dot{\mathbf{\Psi}} + \mathrm{c.c.} - \frac{\hbar^2}{2m} \nabla \mathbf{\Psi}^{\dagger} \cdot \nabla \mathbf{\Psi} + \frac{1}{8\pi G} \Phi \nabla^2 \Phi - m \Phi \mathbf{\Psi}^{\dagger} \mathbf{\Psi} \right]$$





split in "fast" and "slow" parts

Adshead & Lozanov (2021) Jain & MA (2021)

#### non-relativistic limit = multicomponent Schrödinger-Poisson

$$[\mathbf{\Psi}]_i = \psi_i \text{ with } i = 1, 2, 3 \text{ vector}$$
  
 $i\hbar \frac{\partial}{\partial t} \mathbf{\Psi} = -\frac{\hbar^2}{2m} \nabla^2 \mathbf{\Psi} + m \Phi \mathbf{\Psi}$ 

$$[\Psi]_i = \psi_i$$
 with  $i = 1$  scalar

#### case

)

#### $\nabla^2 \Phi = 4\pi G m \, \Psi^\dagger \Psi$

#### case

generalization to s spin-s field field in Jain & MA (2021)



$$N = \int d^{3}x \Psi^{\dagger} \Psi, \quad \text{and} \quad M = mN, \qquad \text{(particle number and rest mass)}$$
$$E = \int d^{3}x \Big[ \frac{\hbar^{2}}{2m} \nabla \Psi^{\dagger} \cdot \nabla \Psi - \frac{Gm^{2}}{2} \Psi^{\dagger} \Psi \int \frac{d^{3}y}{4\pi |\boldsymbol{x} - \boldsymbol{y}|} \Psi^{\dagger}(\boldsymbol{y}) \Psi(\boldsymbol{y}) \Big], \qquad \text{(energy)}$$
$$\boldsymbol{S} = \hbar \int d^{3}x \, i \Psi \times \Psi^{\dagger}, \qquad \text{(spin angular momentum)}$$
$$\boldsymbol{L} = \hbar \int d^{3}x \, \Re \left( i \Psi^{\dagger} \nabla \Psi \times \boldsymbol{x} \right). \qquad \text{(orbital angular momentum)}$$

$$[\mathbf{\Psi}]_i = \psi_i \text{ with } i =$$



### vector vs. scalar DM: two key differences

# interference polarized solitons



#### wave interference



 $|\Psi_a(\boldsymbol{x}) + \Psi_b(\boldsymbol{x})|^2 \neq |\Psi_a(\boldsymbol{x})|^2 + |\Psi_b(\boldsymbol{x})|^2$ 



 $|\Psi_a(\boldsymbol{x}) + \Psi_b(\boldsymbol{x})|^2 = |\Psi_a(\boldsymbol{x})|^2 + |\Psi_b(\boldsymbol{x})|^2$ 





### reduced wave interference in VDM

$$\boldsymbol{\Psi}_{a}(\boldsymbol{x}) = \boldsymbol{\epsilon}_{a}^{(s)} e^{i\boldsymbol{k}_{a}\cdot\boldsymbol{x}}$$

$$|\Psi_{a}(\boldsymbol{x}) + \Psi_{b}(\boldsymbol{x})|^{2} = 2\left(1 + \Re\left[\epsilon_{a}^{(s)\dagger} \cdot \epsilon_{a}^{(s)}e^{-i(\boldsymbol{k}_{a}-\boldsymbol{k}_{b})\cdot\boldsymbol{x}}\right]\right) = 2\left(1 + \operatorname{int}_{(s)}\right)\right)$$

$$\sqrt{\langle \operatorname{int}_{(s)}^{2} \rangle} = \frac{1}{\sqrt{2(2s+1)}}$$

$$\frac{\sqrt{\langle \operatorname{int}_{(0)}^{2} \rangle}}{\sqrt{\langle \operatorname{int}_{(0)}^{2} \rangle}} = \frac{1}{\sqrt{3}}$$

$$\sqrt{\langle \operatorname{int}_{(s)}^2 \rangle} = \frac{1}{\sqrt{2(2s+1)}}$$

s = 0 for SDM and s = 1 for VDM

MA, Jain, Karur & Mocz (2022)







#### solitons in massive spin-0 (scalar fields)

 $\Psi_{\rm sol}(t,\boldsymbol{x}) = \psi_{\rm sol}(\mu,\boldsymbol{x})e^{i\mu c^2 t/\hbar}$ 

$$M_{\rm sol} \approx 60.7 \frac{m_{\rm pl}^2}{m} \sqrt{\frac{\mu}{m}}$$
$$R_{\rm sol} \sim \sqrt{\frac{m}{\mu}} \frac{\hbar}{mc}$$

 $\mu/m \ll 1$ 





lowest energy state for fixed NN = # of particles in soliton

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#### "polarized" vector solitons

$$\Psi_{\rm sol}(t, \boldsymbol{x}) = \psi_{\rm sol}(\mu, r) e^{i\mu c^2 t/\hbar} \boldsymbol{\epsilon}$$

 $\epsilon^{\dagger}\epsilon = 1$ 



 $\boldsymbol{W}(t,\boldsymbol{x}) \equiv \frac{\hbar}{\sqrt{2mc}} \Re \left[ \boldsymbol{\Psi}(t,\boldsymbol{x}) e^{-imc^2 t/\hbar} \right]$ 



$$\boldsymbol{\epsilon}_{1,\hat{z}}^{(0)} = \begin{pmatrix} 0\\0\\1 \end{pmatrix}$$

$$\boldsymbol{\epsilon}_{1,\hat{z}}^{(\pm 1)} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1\\ \pm i\\ 0 \end{pmatrix}$$



#### "polarized" vector solitons

$$\Psi_{\rm sol}(t, \boldsymbol{x}) = \psi_{\rm sol}(\mu, r) e^{i\mu c^2 t/\hbar} \boldsymbol{\epsilon} \qquad \boldsymbol{\epsilon}^{\dagger} \boldsymbol{\epsilon} = 1$$

$$\mathbf{S}_{sol} \approx i(\boldsymbol{\epsilon} \times \boldsymbol{\epsilon}^{\dagger}) 60.7 \frac{m_{pl}^2}{m^2} \sqrt{\frac{\mu}{m}} \hbar = i(\boldsymbol{\epsilon} \times \boldsymbol{\epsilon}^{\dagger}) \frac{M_s}{m^2}$$

macroscopic spin

$$S_{\rm tot}/\hbar = \lambda N \hat{z}$$

N = # of particles in soliton

 $\boldsymbol{W}(t,\boldsymbol{x}) \equiv \frac{\hbar}{\sqrt{2mc}} \Re \left[ \boldsymbol{\Psi}(t,\boldsymbol{x}) e^{-imc^2 t/\hbar} \right]$ 



 $S_{\rm tot} = 0\hat{z}$ 

 $\boldsymbol{S}_{\mathrm{tot}} = \hbar \frac{M_{\mathrm{sol}}}{2} \hat{z}$ 

Jain & MA (2021)









### "polarized" vector solitons

$$\Psi_{\rm sol}(t, \boldsymbol{x}) = \psi_{\rm sol}(\mu, r) e^{i\mu c^2 t/\hbar} \boldsymbol{\epsilon} \qquad \boldsymbol{\epsilon}^{\dagger} \boldsymbol{\epsilon} = 1$$

$$\mathbf{S}_{sol} \approx i(\boldsymbol{\epsilon} \times \boldsymbol{\epsilon}^{\dagger}) 60.7 \frac{m_{pl}^2}{m^2} \sqrt{\frac{\mu}{m}} \hbar = i(\boldsymbol{\epsilon} \times \boldsymbol{\epsilon}^{\dagger}) \frac{M_s}{m^2}$$

macroscopic spin  $S_{\rm tot}/\hbar = \lambda N \hat{z}$ N = # of particles in soliton

- all lowest energy for fixed N
- bases for partially-polarized solitons



 $\boldsymbol{W}(t,\boldsymbol{x}) \equiv \frac{\hbar}{\sqrt{2mc}} \Re \left[ \boldsymbol{\Psi}(t,\boldsymbol{x}) e^{-imc^2 t/\hbar} \right]$ 



 $S_{\text{tot}} = 0\hat{z}$ 

 $\boldsymbol{S}_{\mathrm{tot}} = \hbar \frac{M_{\mathrm{sol}}}{2} \hat{z}$ 

 $0 \le |\mathbf{S}_{ ext{tot}}| \le rac{M_{ ext{sol}}}{m}\hbar$ 





### a different higher energy soliton: the "hedgehogs"

earlier literature

$$W_j(\mathbf{x},t) = f(r)\frac{x^j}{r}\cos\omega t$$
,



 $E_{\rm hh}^s > E$ 

 $E_{\rm hh}^{s=1} \approx 0.33E < 0$ 

hedgehogs not ground states









### vector vs. scalar DM: two key differences

# interference polarized solitons

3+1 dimensional simulations





# SDM



0.34  $t/t_{\rm dyn} \longrightarrow$ 





#### Vector & Scalar Dark Matter

1.36

MA, Jain, Karur & Mocz (2022)





### radial density profiles

#### scalar vs. vector dark matter

- less dense & broader core
- smoother transition to  $r^{-3}$  tail





MA, Jain, Karur & Mocz (2022)



### radial density profiles



#### vector vs. scalar dark matter

- shape difference
- smoother transition to  $r^{-3}$  tail

















### 2-point density correlation









### gravitational implications (examples)

# - dynamical heating of stars Church et. al (2021), Dalal & Kravtsov (2022) - cusp-core, diversity

- lensing



#### core-halo mass relation





 $\Xi = |E_{\rm tot}| / (M_{\rm tot}^3 (Gm/\hbar)^2)$ 



 $f_{
m v} = 0.56 \pm 0.03$  $f_{
m s} = 0.61 \pm 0.01$  $M_{\rm f} = f(M_1 + M_2)$  $\frac{M_{\rm core}}{M_{\rm tot}} \propto \Xi^{\log_2 f}$ 



#### Du et. al (2017)



## intrinsic spin







MA, Jain, Karur & Mocz (2022)

spin density  $s = i\hbar\Psi imes\Psi^\dagger$ 

spin  
$$\mathbf{S} = \hbar \int \mathrm{d}^3 x \, i \mathbf{\Psi} \times \mathbf{\Psi}^\dagger$$

 $\boldsymbol{s} = (2mc/\hbar)\boldsymbol{W} \times (\dot{\boldsymbol{W}} - \nabla W_0)$ 



## generation of spin density



#### spin with self-interactions: vector oscillons

 $\mathcal{L} = -\frac{1}{4}\mathcal{G}_{\mu\nu}\mathcal{G}^{\mu\nu} - \frac{m^2}{2}W_{\mu}W^{\mu} + \frac{\lambda}{4}(V_{\mu\nu})^2 + \frac{1}{4}(W_{\mu\nu})^2 + \frac{1}{4$ 

 $S_{\rm tot} \neq 0$ 

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$$W_{\mu}W^{\mu})^2 - \frac{h}{6}(W_{\mu}W^{\mu})^3$$



Zhang, Jain & MA (2022)

$$_{\rm ot} = 0$$

self-interaction supported NOT degenerate in energy spin-spin interaction matters!

$$\mathcal{L} = \Re[i \mathbf{\Psi}^{\dagger} \dot{\mathbf{\Psi}}] - rac{1}{2m} 
abla \mathbf{\Psi}^{\dagger} \cdot 
abla \mathbf{\Psi} - V_{
m nl}(\mathbf{\Psi})$$

$$V_{\rm nl}(\mathbf{\Psi}^{\dagger}, \mathbf{\Psi}) = -\frac{3\lambda}{8m^2} (\mathbf{\Psi}^{\dagger}\mathbf{\Psi})^2 + \frac{5\gamma}{12m^3} (\mathbf{\Psi}^{\dagger}\mathbf{\Psi})^3 + \left[\frac{\lambda}{8m^2} - \frac{\gamma}{4m^3} (\mathbf{\Psi}^{\dagger}\mathbf{\Psi})\right] (\mathbf{S} \cdot \mathbf{S})$$









### implications of non-zero spin solitons?



#### different for higher spin solitons?



Helfer, Lim, Garcia, MA (2016)



above for compact scalar solitons using full numerical GR



## electromagnetic coupling and radiation (axion + photons)

MA, Long, Mou & Saffin (2021)







 $\mathcal{L}_{int} \sim g_{\phi\gamma} \phi F_{\mu\nu} \tilde{F}^{\mu\nu}$ 

MA & Mou (2020)

Motivated by earlier work by Hertzberg & Schippacaise (2018)





## spin-s + photons: spin of soliton & polarization of photons





 $\mathcal{L}_{int} \sim \begin{cases} g_{W\gamma}^2 W_{\mu} W^{\mu} F_{\alpha\beta} \tilde{F}^{\alpha\beta} \\ g_{H\gamma}^2 (H_{\mu\nu} H^{\mu\nu} - H^2) F_{\alpha\beta} \tilde{F}^{\alpha\beta} \end{cases}$ spin-1 spin-2

 $\sim g_{\mathcal{F}\gamma}^2 \operatorname{Tr}[\mathcal{F}\mathcal{F}] F_{\alpha\beta} \tilde{F}^{\alpha\beta}$ NR limit

 $\epsilon^2 F_{\mu\nu} \mathcal{G}^{\mu\nu}$ 











### formation mechanisms?



Gorghetto, Hardy, March-Russell, Song & West (2022)

#### misaligned scalar + self-interactions?



#### gravitational in late universe

Core 105 103 101 200 kpc 50 kpc Schive et. al (2014)





### gravitational particle production to nonlinear structures

cannot easily do ultralight dark photons 1  

$$\Omega_{\rm vdm} \sim 0.3 \left(\frac{m}{10^{-5} \,{\rm eV}}\right)^{1/2} \left(\frac{H_{\rm inf}}{10^{14} \,{\rm GeV}}\right)^4$$
 10<sup>-2</sup>  
Graham, Mardon, Rajendran (2016) 10<sup>-4</sup>  
Ahmed, Grzadkowski, Socha (2020)  $\mathcal{P}_{\delta}(k)$  10<sup>-6</sup>  
Kolb & Long (2020) 10<sup>-8</sup>

 $10^{-10}$  $10^{-6}$ 

$$M_{\rm sol}(a) \sim 10^{-23} M_{\odot} \left(\frac{a_{\rm eq}}{a}\right)^{3/4} \left(\frac{{\rm eV}}{m}\right)^{3/2}$$
  
 $R_{\rm sol}(a) \sim 10^4 \,{\rm km} \left(\frac{a}{a_{\rm eq}}\right)^{3/4} \left(\frac{{\rm eV}}{m}\right)^{1/2}$ 







#### abundance and detection

$$n = f_s \overline{\rho}(t_0)/M \simeq 10^{20} \mathrm{pc}^{-3} \left(\frac{f_s}{0.05}\right) \left(\frac{\rho_{\mathrm{local}}}{0.5 \,\mathrm{GeV/cm^3}}\right) \left(\frac{0.1 M_J^{\mathrm{eq}}}{M}\right) \left(\frac{m}{\mathrm{eV}}\right)^{3/2}$$

$$\Gamma \simeq n\pi R^2 v_{\rm rel}$$
$$\simeq \frac{0.1}{\rm yr} \left(\frac{m}{\rm eV}\right)^{1/2} \left(\frac{0.1M_J^{\rm eq}}{M}\right)^3 \left(\frac{v_{\rm rel}}{10^{-3}}\right) \left(\frac{f_s}{0.05}\right) \left(\frac{\rho_{\rm local}}{0.5\,{\rm GeV/cm^3}}\right)$$

Gorghetto et. al (2022)





Dror, Harigaya, and Narayan (2018) Bastero-Gil, Santiago, Ubaldi, Vega-Morales (2018)







## generalization to arbitrary spin

#### spin-s field as dark matter



Jain & MA (2021)



#### non-relativistic limit = multicomponent Schrödinger-Poisson

spin-s fields as light dark matter



-phenomenology/numerical simulations

- interference  $\sim 1/(2s+1)$ 

2s+1 component

Schrödinger

Jain & MA (2021)



### extremally polarized solitons

spin-s fields as dark matter



Jain & MA (2021)

macroscopic spin  $S_{tot}/\hbar = \lambda N \hat{z}$ N = # of particles in soliton



#### non-topological solitons spatially localized, coherently oscillating, long-lived





spatially localized

coherently oscillating (components)

exceptionally long-lived

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#### scalar

#### vector

#### tensor



### the near future to do list ...



### future possibilities ...

- formation and survival mechanisms
- BH superradiance with higher spin fields\* (already done)
- dynamical friction
- vortices
- condensation time scales
- lifetimes of higher spin solitons
- initial power spectrum of fluctuations in the fields

#### - direct/indirect detection (interaction with baryons? photons? kinetic mixing)