

Gravitational Aspects of Solitons in Cosmology

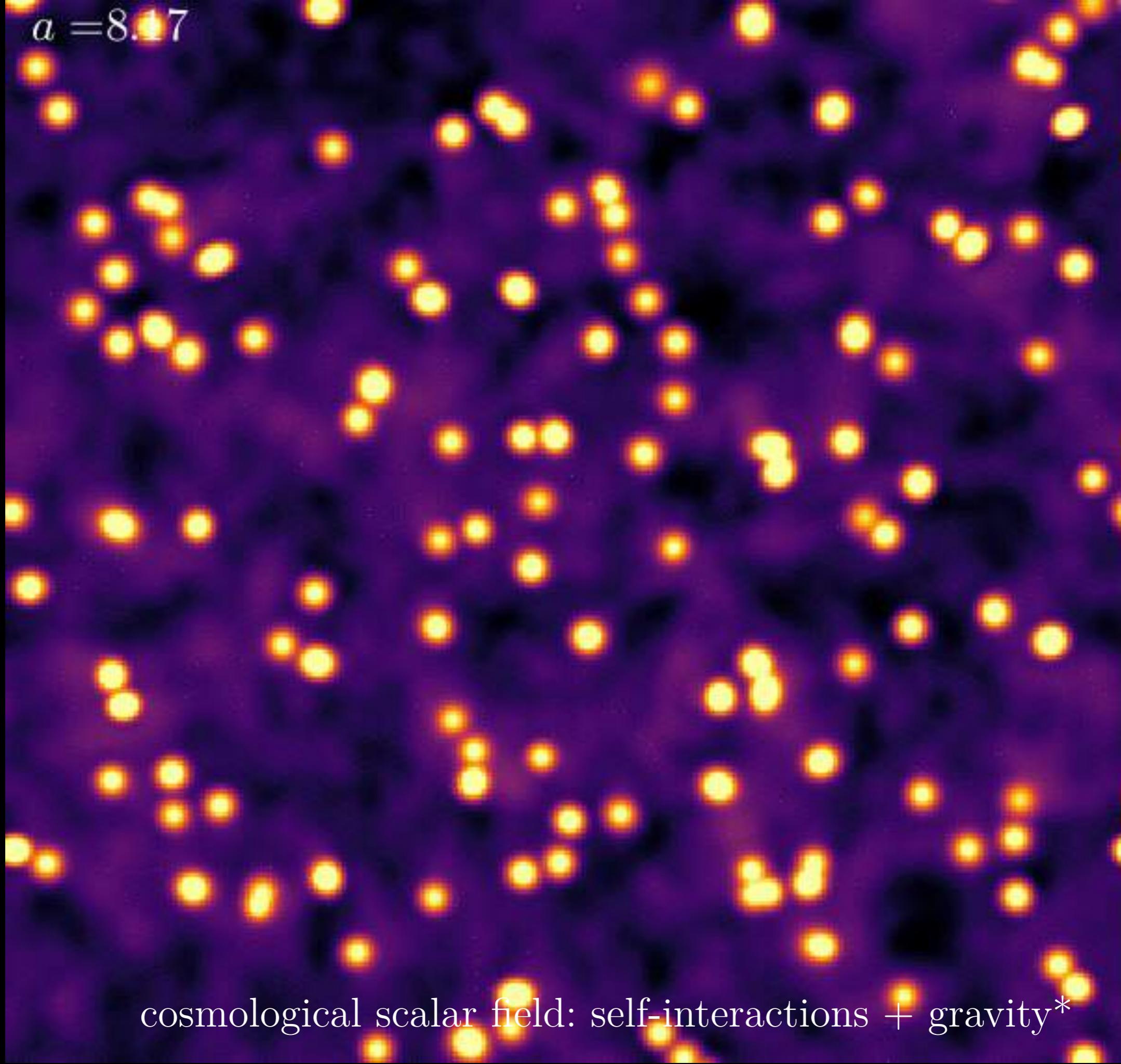


Mustafa A. Amin
Feb 25, 2019, SITP Colloquium



RICE

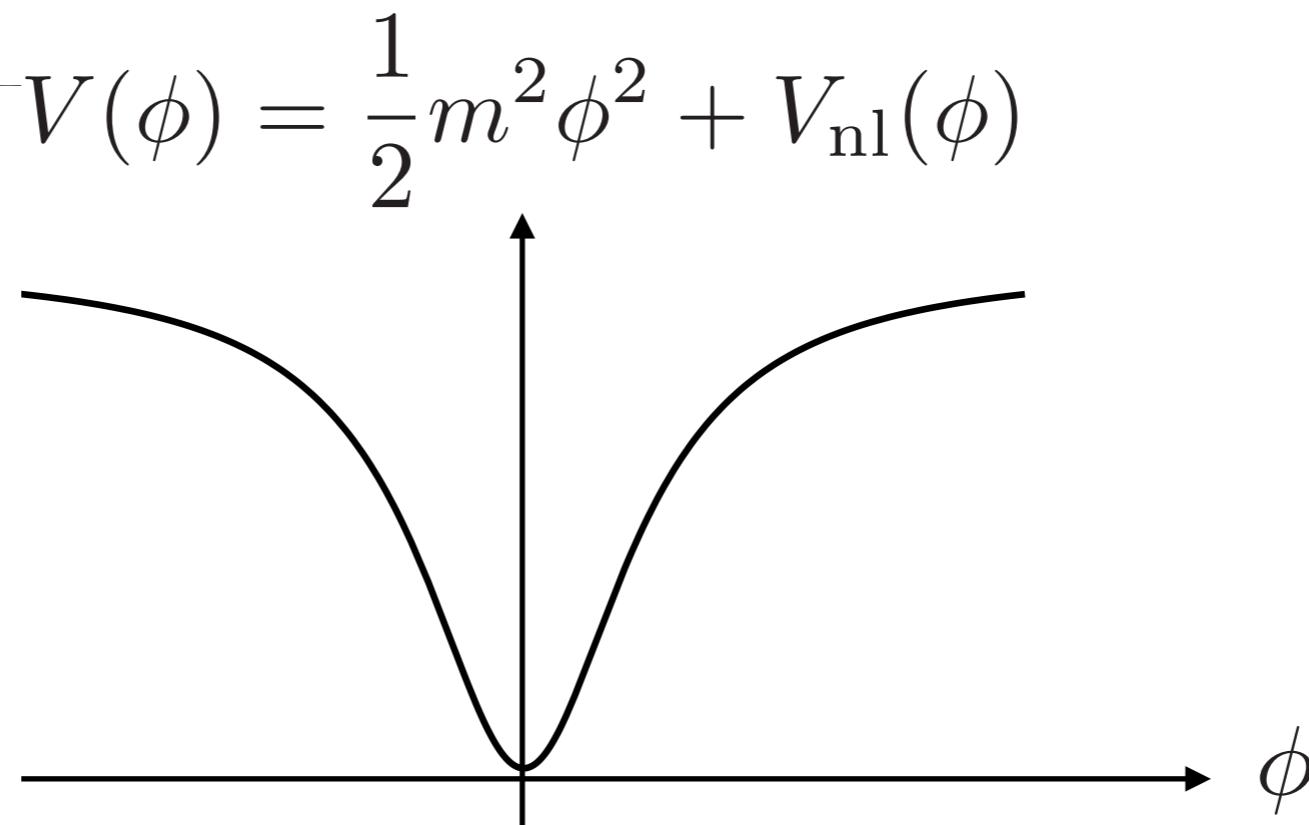
$a = 8.17$



cosmological scalar field: self-interactions + gravity*

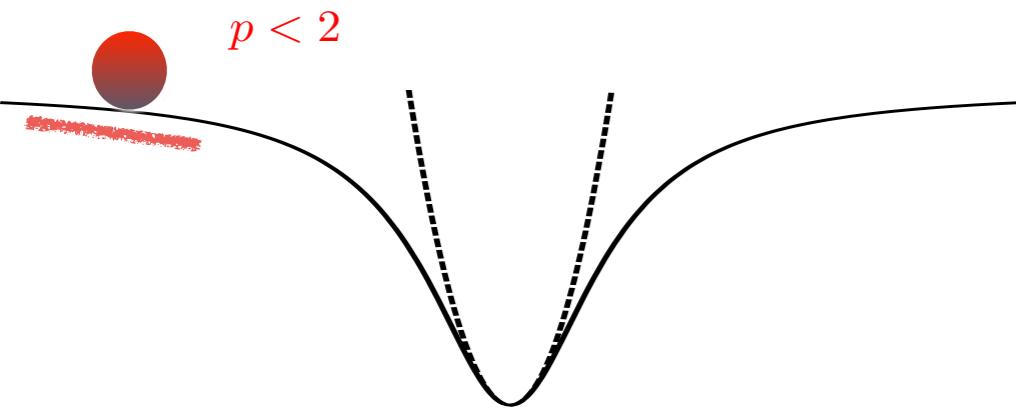
cosmological scalar fields self-interaction + gravity

$$S = \int d^4x \sqrt{-g} \left[\frac{m_{\text{pl}}^2}{2} R - \frac{1}{2} (\partial\phi)^2 - V(\phi) \right]$$



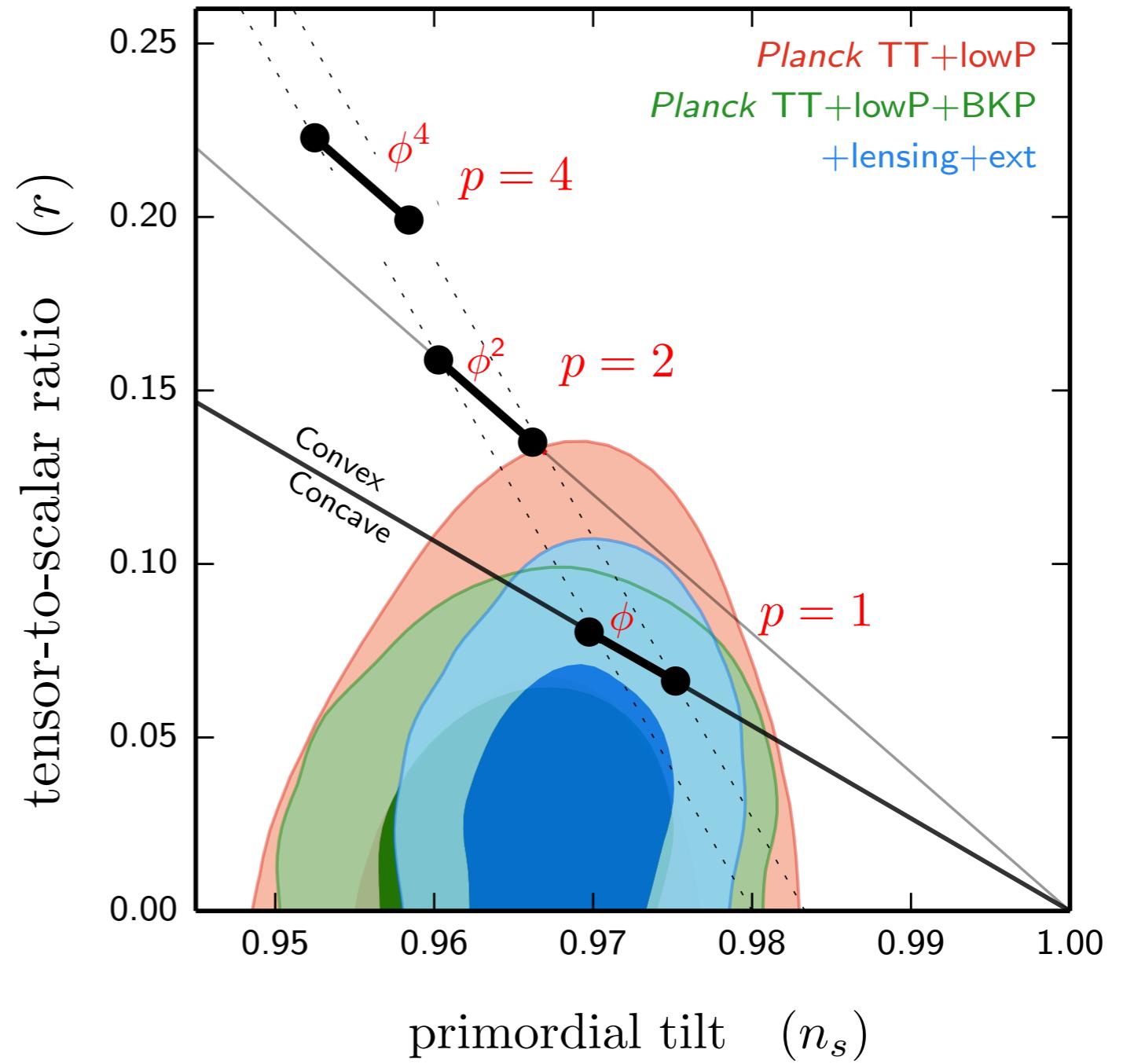
inflation: post-inflationary dynamics

$$V(\phi) \propto \phi^p$$

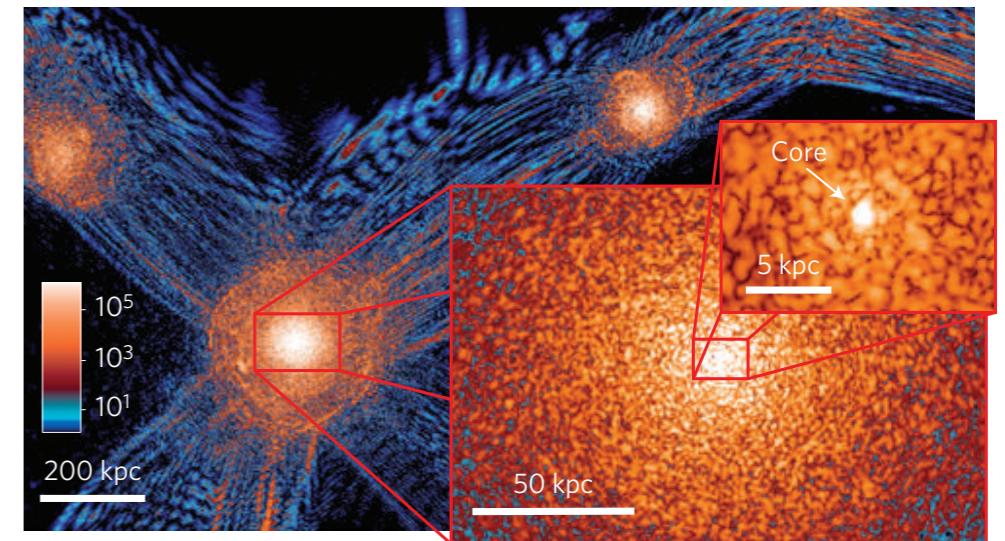
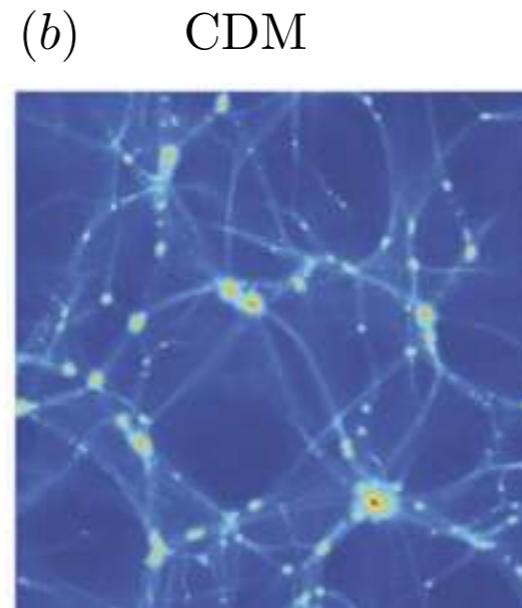
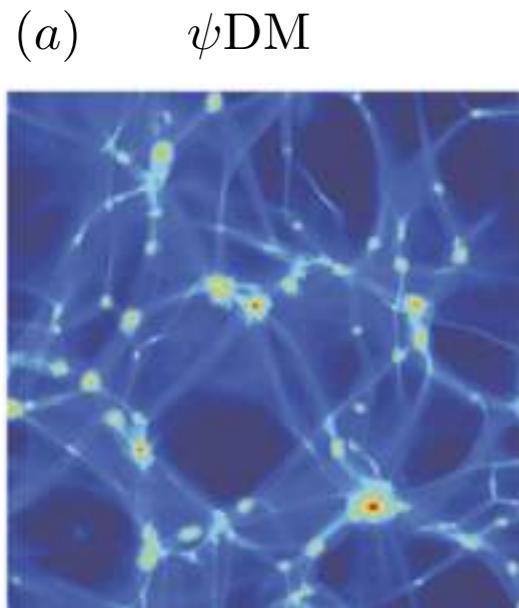


for example:

Starobinsky Inflation (1979)
Silverstein & Westphal (2008)
Kallosh & Linde (2013)



dark matter: axion-like fields



Schive et. al (2014)

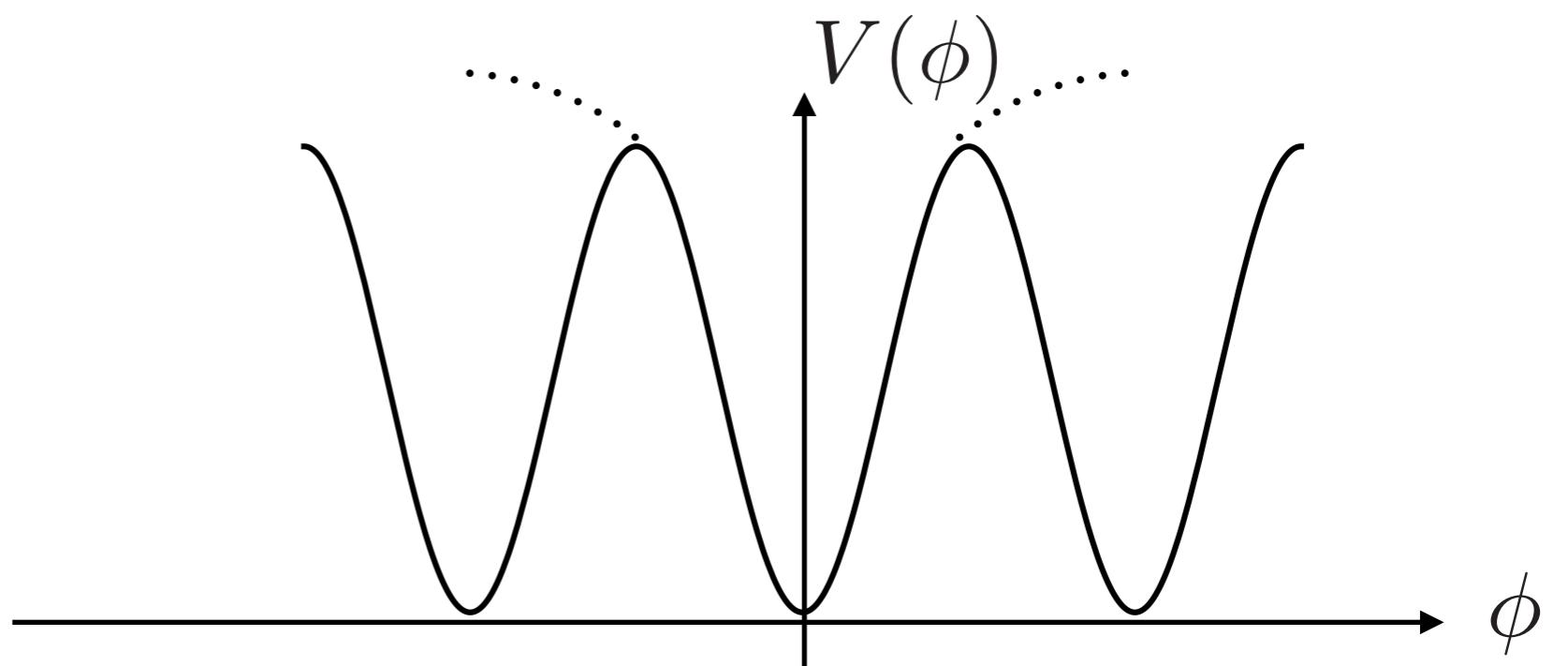
for example:

Peccei & Quinn (1977)

Hu, Barkana & Gruzinov (2000)

Arvanitaki et. al (2009)

Hui et. al (2016)



implications

- eq. of state after inflation ?
 - stochastic gravitational wave-generation ?
 - constrained by N_{eff} or direct detection
 - primordial black hole formation ?
-

- distinguishability from WIMPS ?
- compact objects
 - eg. sources of gravitational waves ?

eq. of motion for cosmological fields

$$\square\phi + V'(\phi) = 0$$

nonlinear Klein-Gordon eq.

$$G_{\mu\nu} = \frac{1}{m_{\text{pl}}^2} T_{\mu\nu}$$

Einstein Eq.

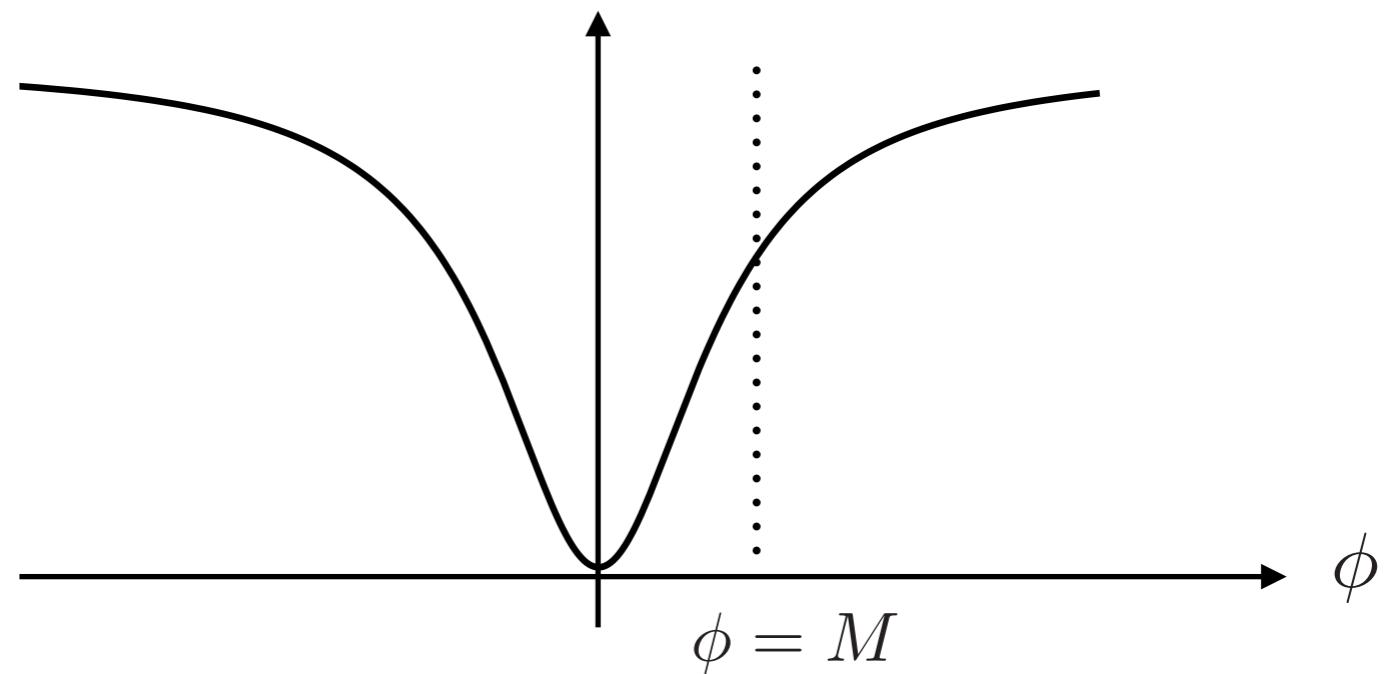
$$V(\phi) = \frac{1}{2}m^2\phi^2 + V_{\text{nl}}(\phi)$$

examples:

$$V(\phi) = \frac{m^2 M^2}{2} \tanh^2 \left(\frac{\phi}{M} \right)$$

$$V(\phi) = m^2 M^2 \left[\sqrt{1 + \frac{\phi^2}{M^2}} - 1 \right]$$

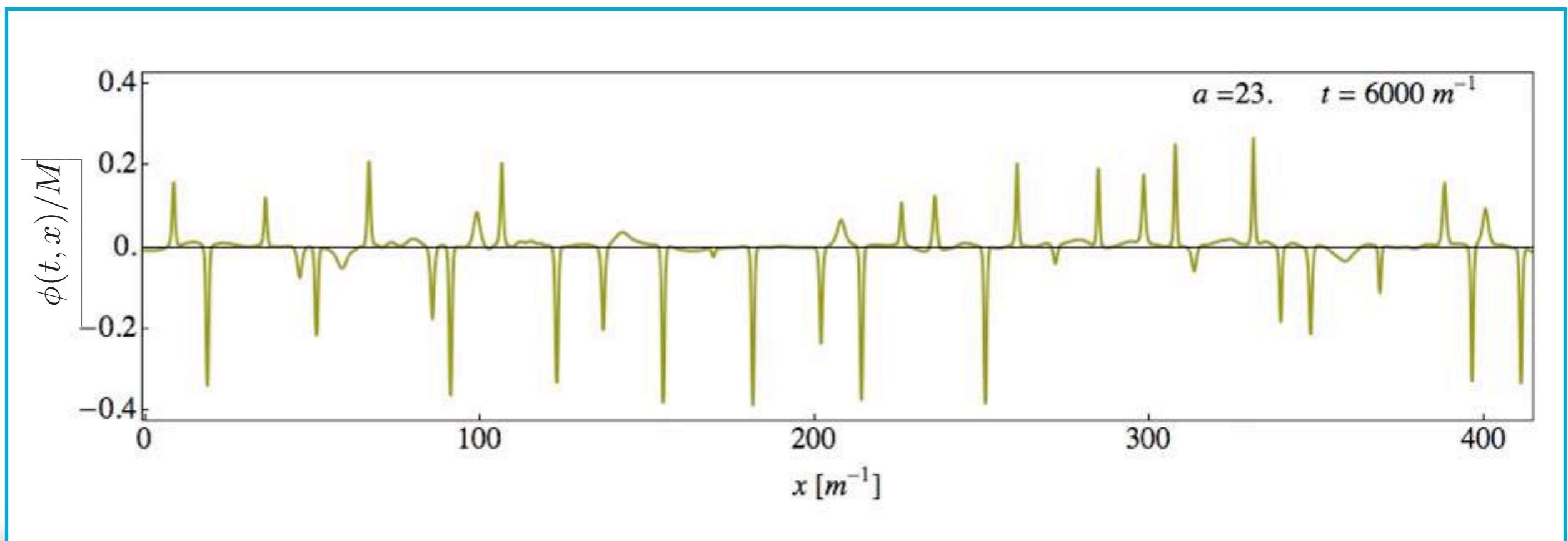
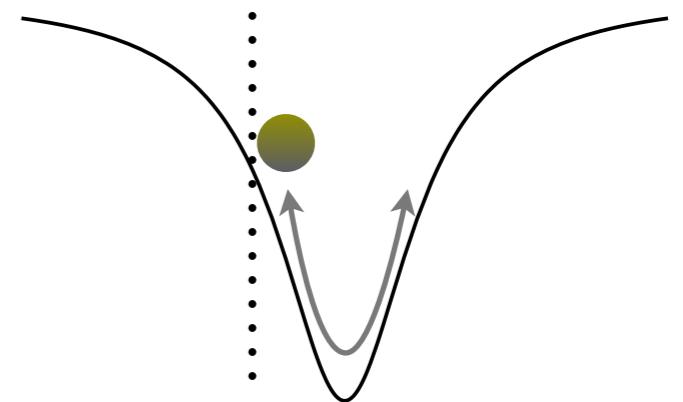
$$V(\phi) = m^2 M^2 \left[1 - \cos \frac{\phi}{M} \right]$$



dynamics of oscillating fields

expansion ✓
self-interactions ✓
gravitational int. ✗

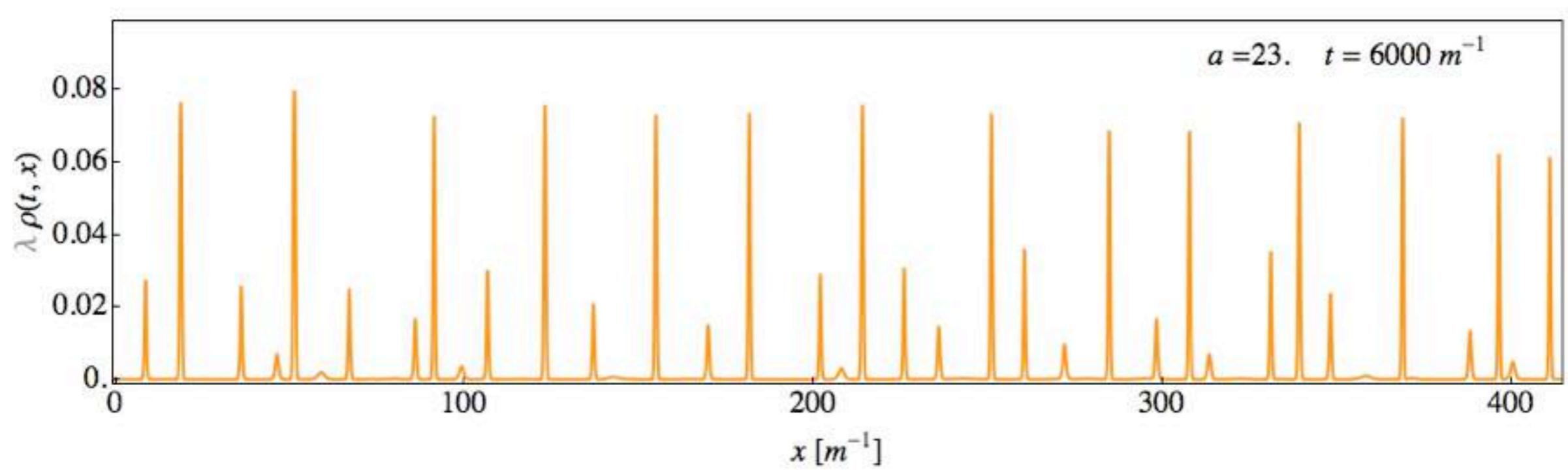
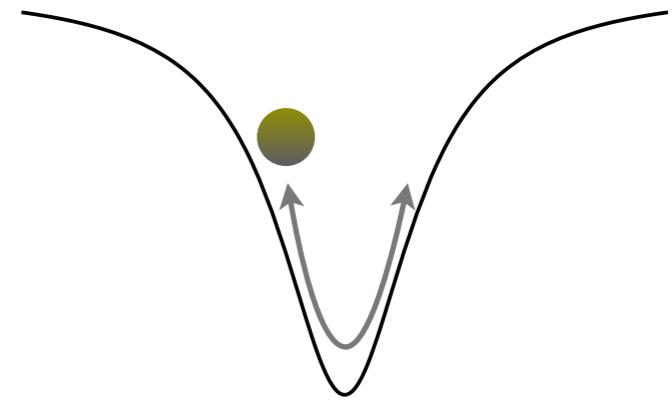
$$\square\phi = V'(\phi)$$



MA (2010)
Khlopov, Malomed & Zeldovich (1985)

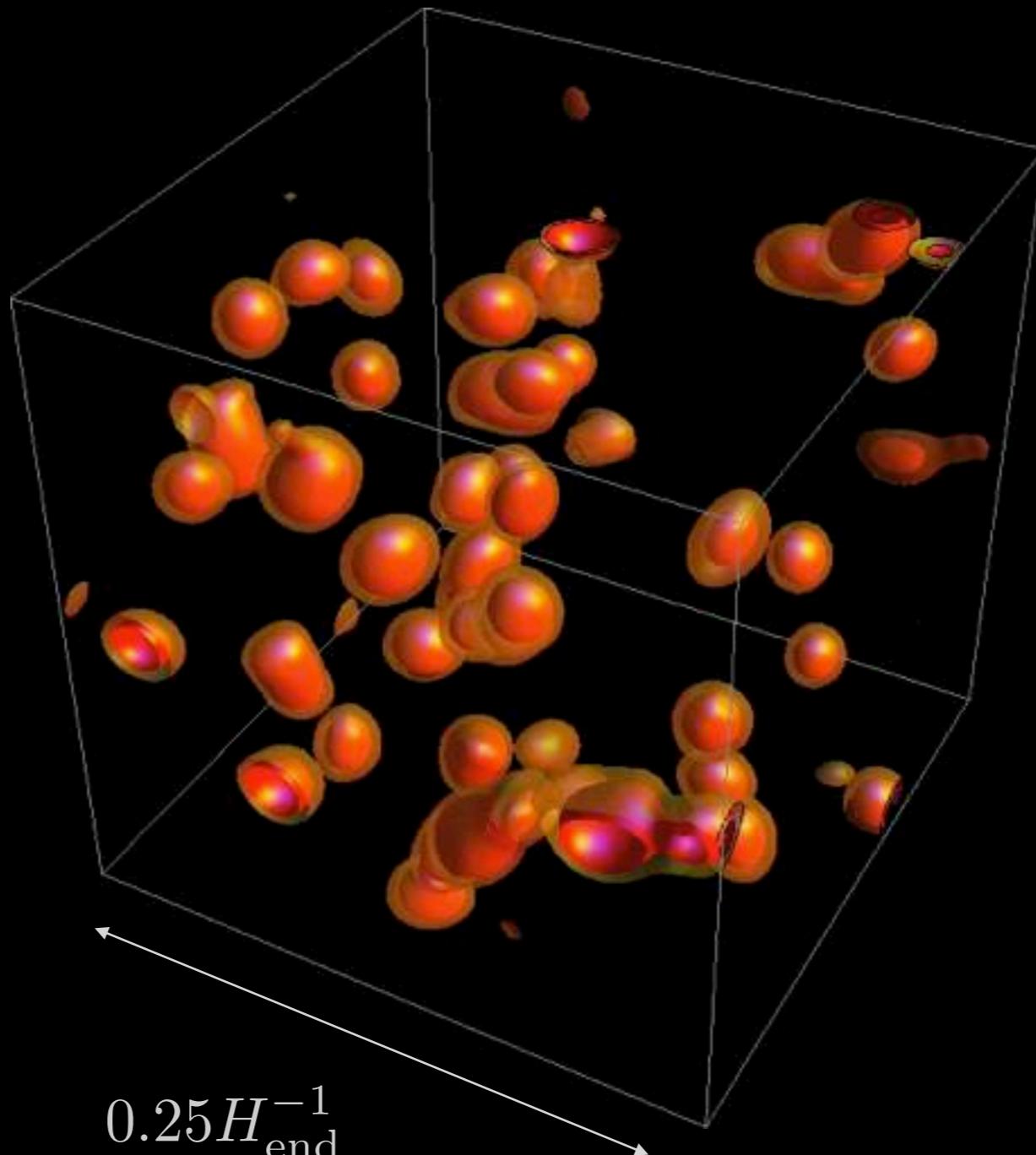
dynamics of oscillating fields

$$\square\phi = V'(\phi)$$



soliton formation in relativistic fields

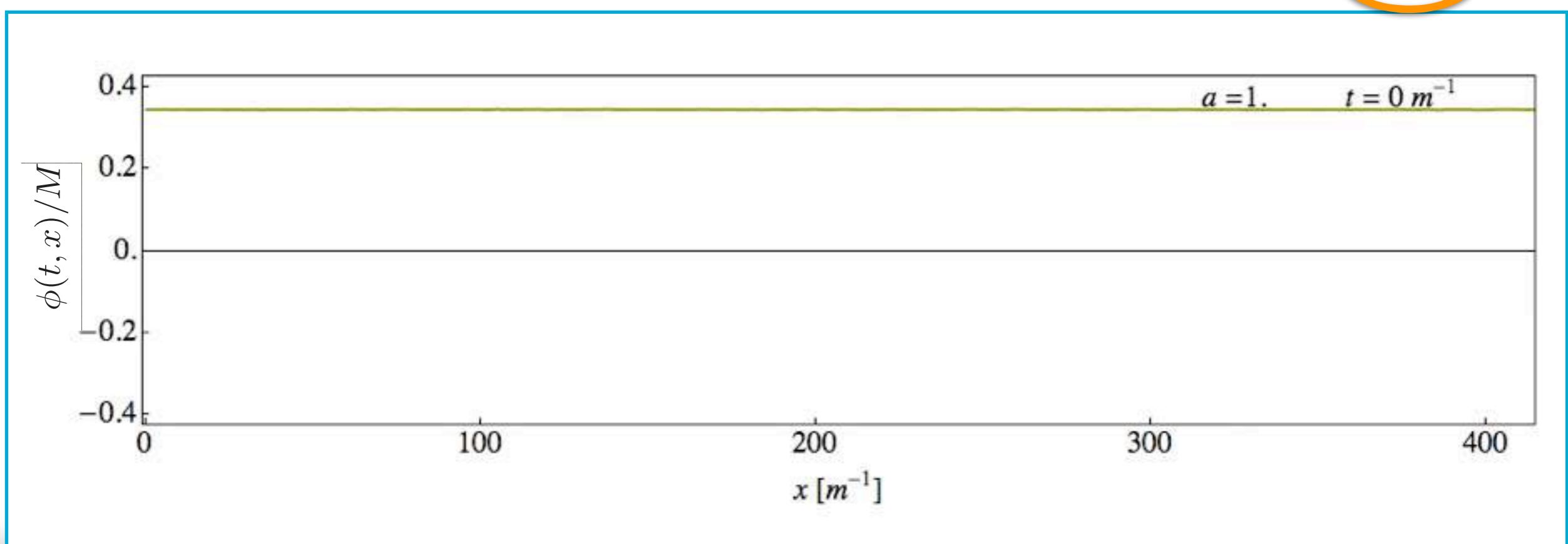
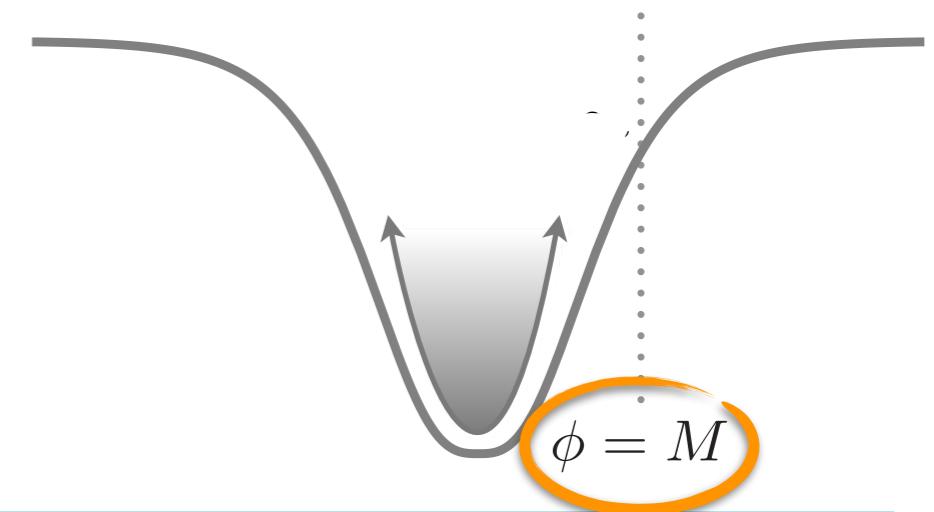
- expansion ✓
- self-interactions ✓
- gravitational int. ✗



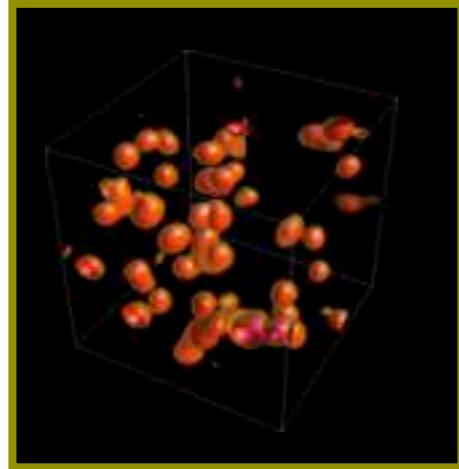
MA, Easter, Finkel, Flaucher & Hertzberg (2011)

condition for rapid fragmentation ?

$$\frac{\text{growth-rate of fluctuations}}{\text{expansion rate}} \sim \frac{m_{\text{pl}}}{M} \gg 1$$



solitons?



existence and stability:

MA (2013)

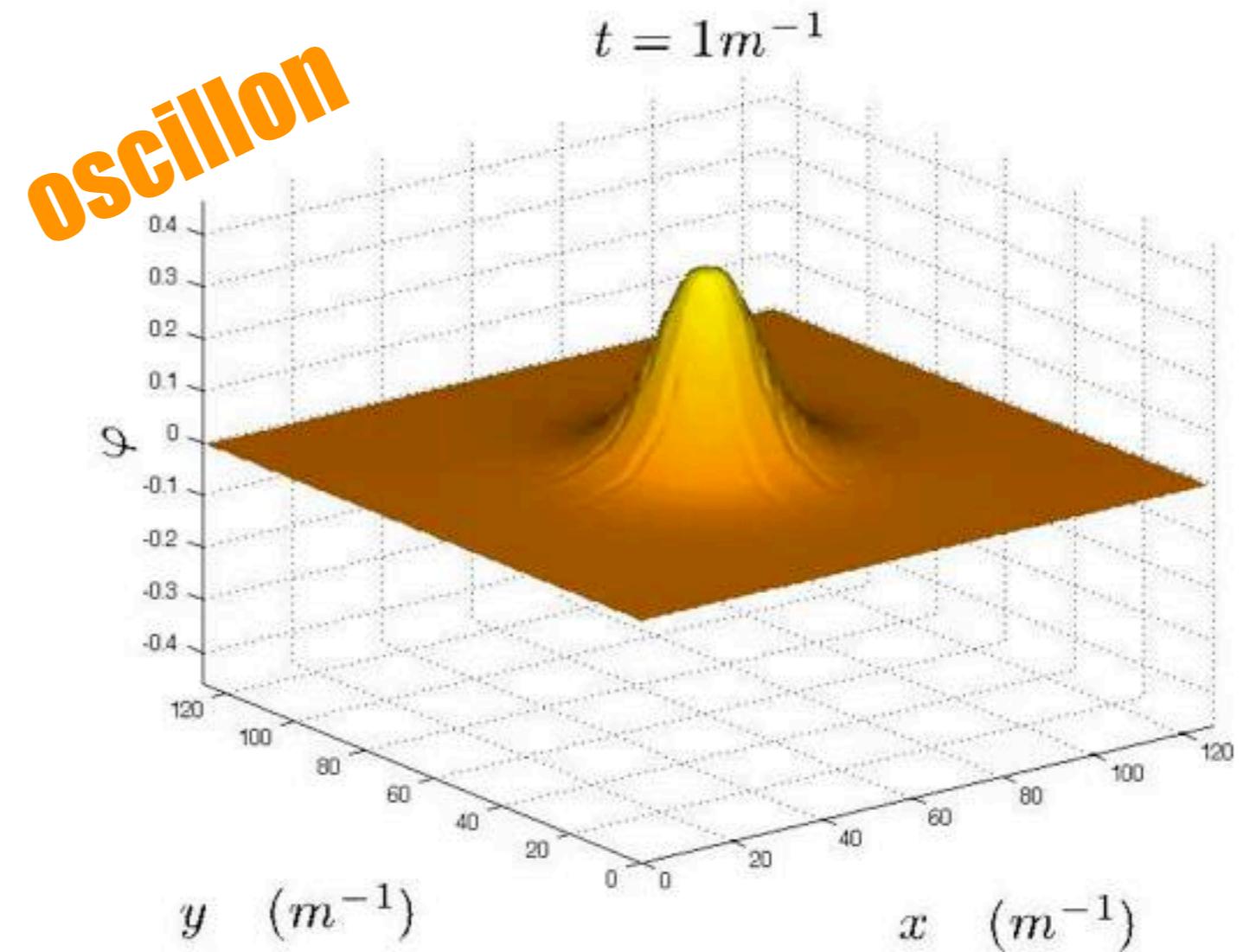
MA & Shirokoff (2010)

Segur & Kruskal (1987)

Hertzberg (2011)

Mukaido et. al (2016,17)

(1) oscillatory (2) spatially localized (3) **very long lived**



Bogolubsky & Makhankov (1976), Gleiser (1994), Copeland et. al (1995)

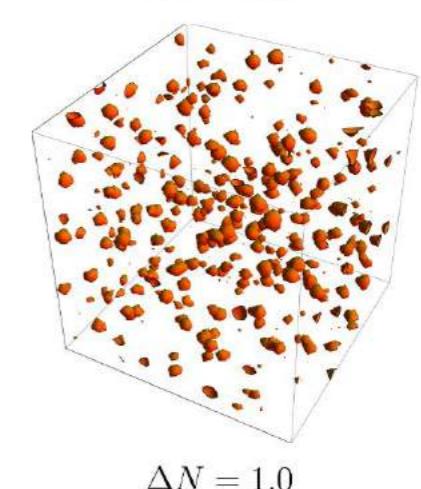
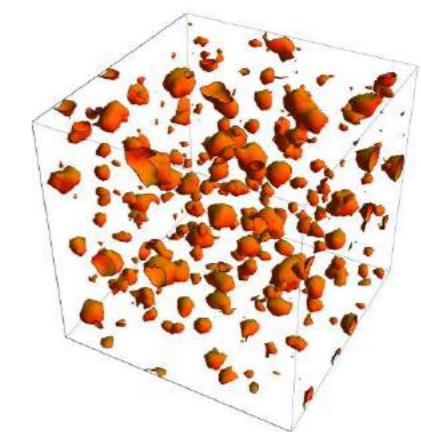
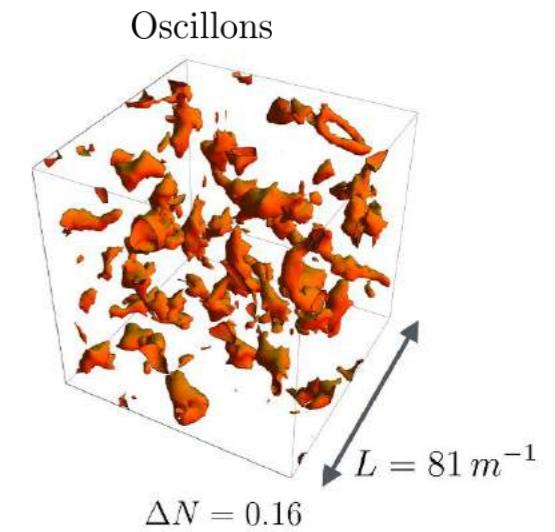
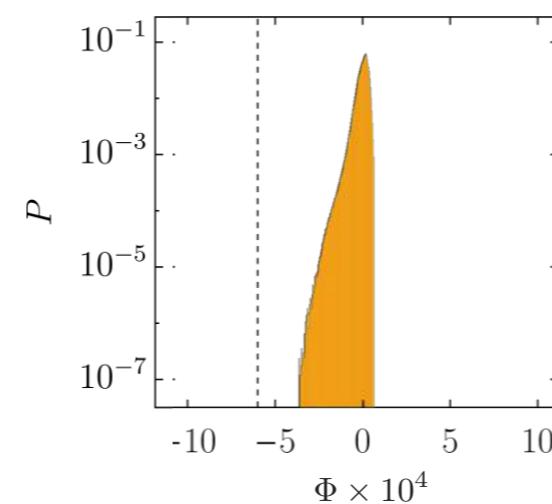
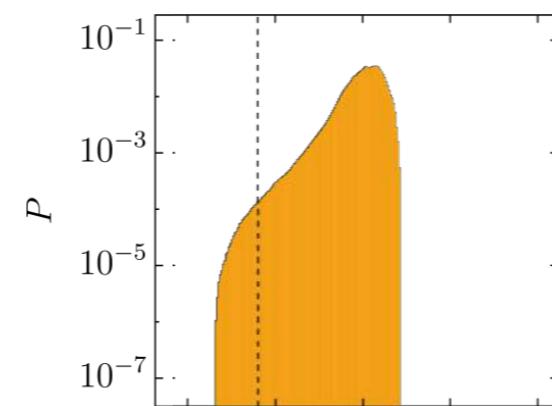
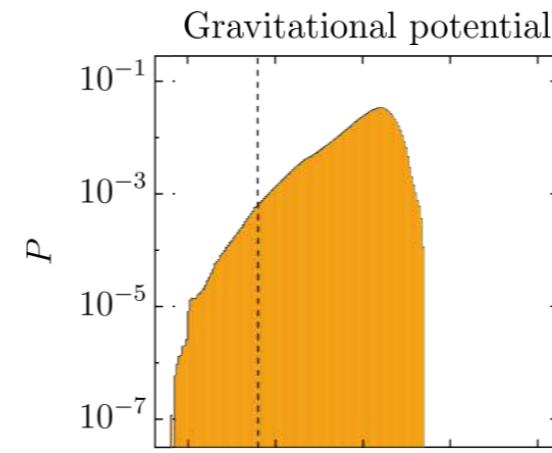
expansion ✓
self-interactions ✓
gravitational int. ✗

passively calculated gravitational potential

$$\Phi \lesssim \text{few} \times 10^{-3}$$



w/ K. Lozanov (2019)



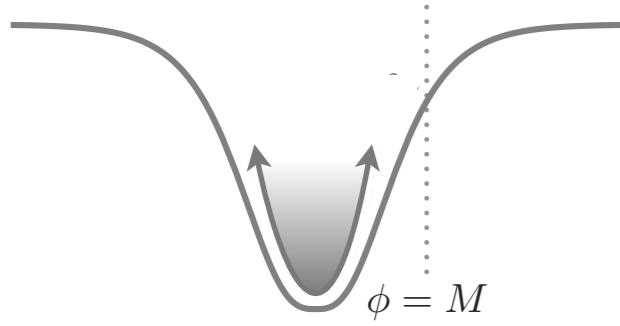
Time

expansion ✓

self-interactions ✓

gravitational int. ✗

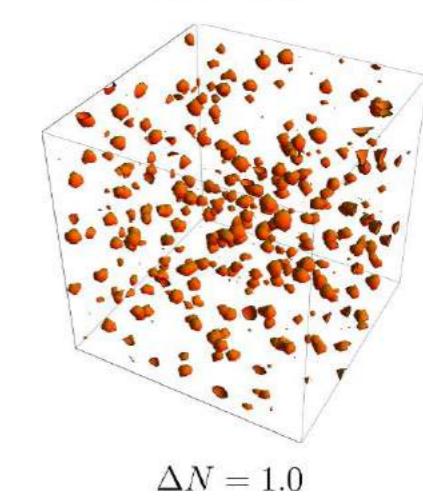
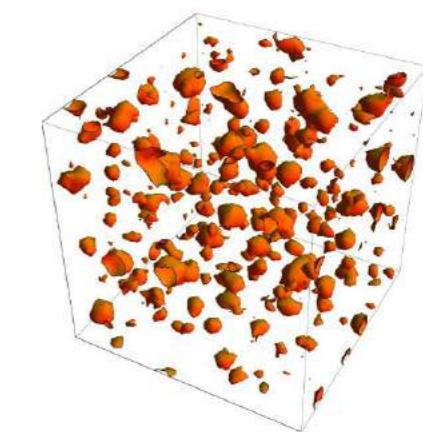
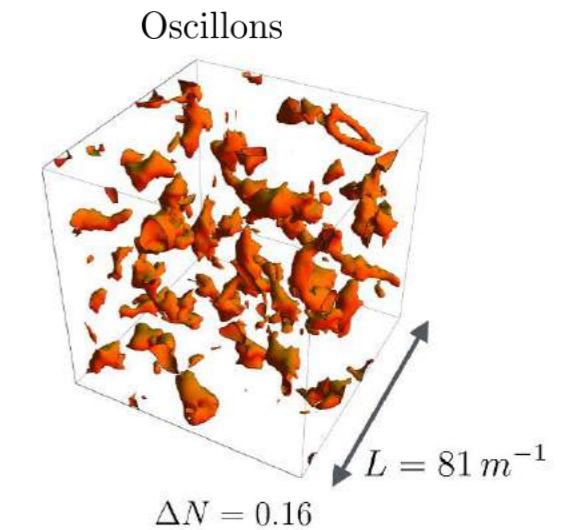
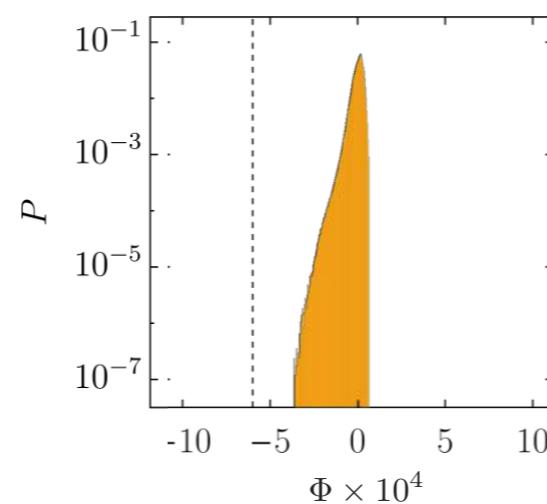
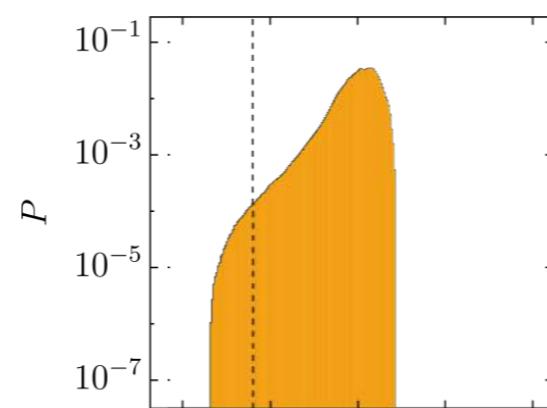
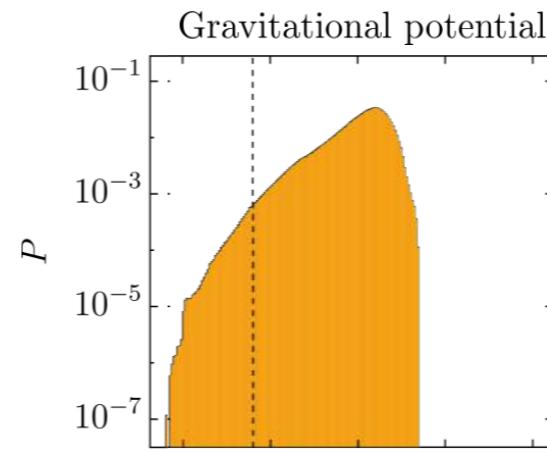
passively calculated gravitational potential



$$|\Phi|_{\text{sol}} \lesssim 10 \times \left(\frac{M}{m_{\text{pl}}} \right)^2$$

growth-rate of fluctuations
expansion rate

$$\sim \frac{m_{\text{pl}}}{M} \gg 1$$

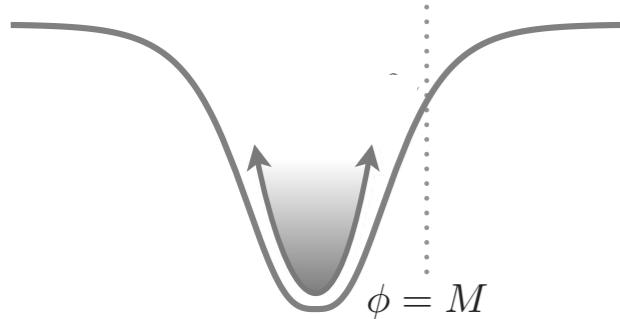


expansion ✓

self-interactions ✓

gravitational int. ✗

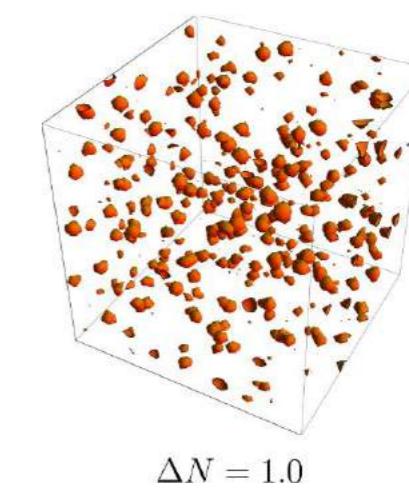
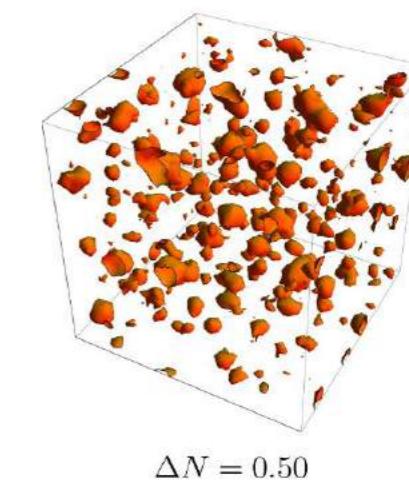
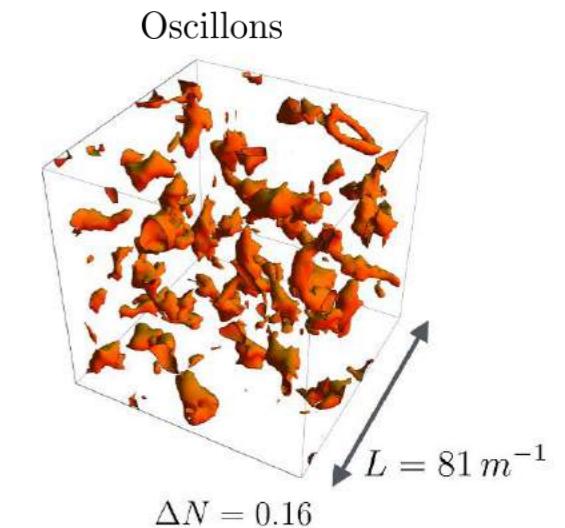
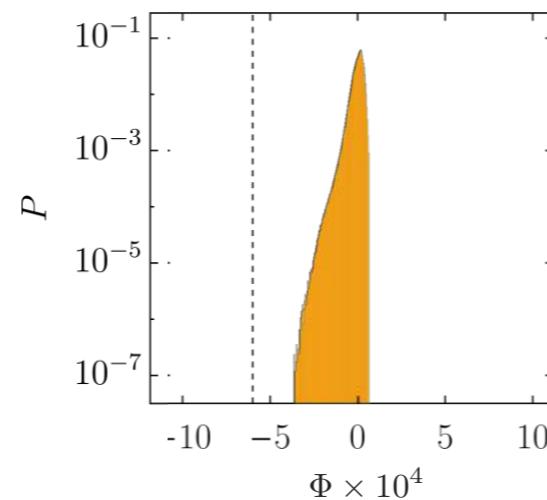
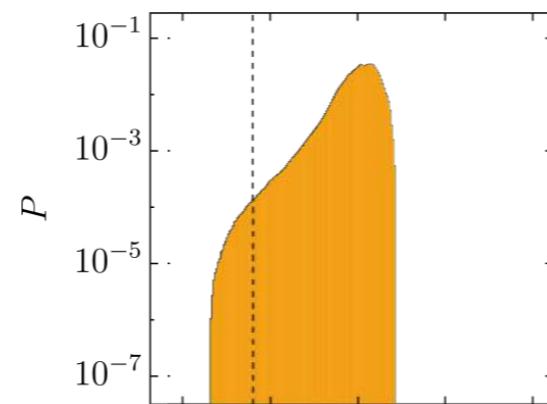
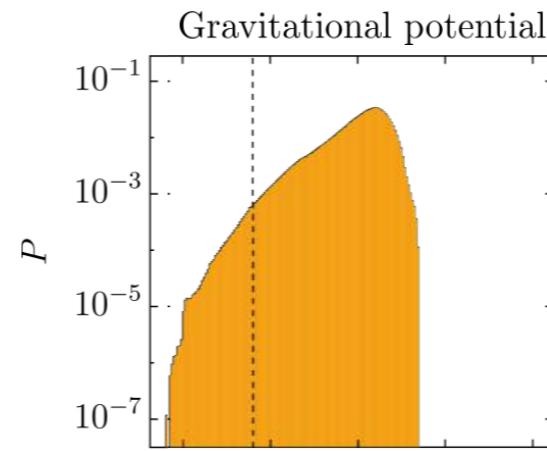
passively calculated gravitational potential



$$|\Phi|_{\text{sol}} \lesssim 10 \times \left(\frac{M}{m_{\text{pl}}}\right)^2$$
$$\Phi \lesssim \text{few} \times 10^{-3}$$

not easy to form black holes from individual solitons*

$$\frac{\text{growth-rate of fluctuations}}{\text{expansion rate}} \sim \frac{m_{\text{pl}}}{M} \gg 1$$



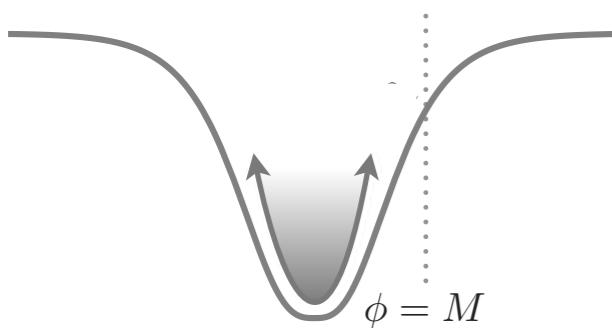
Time

expansion ✓

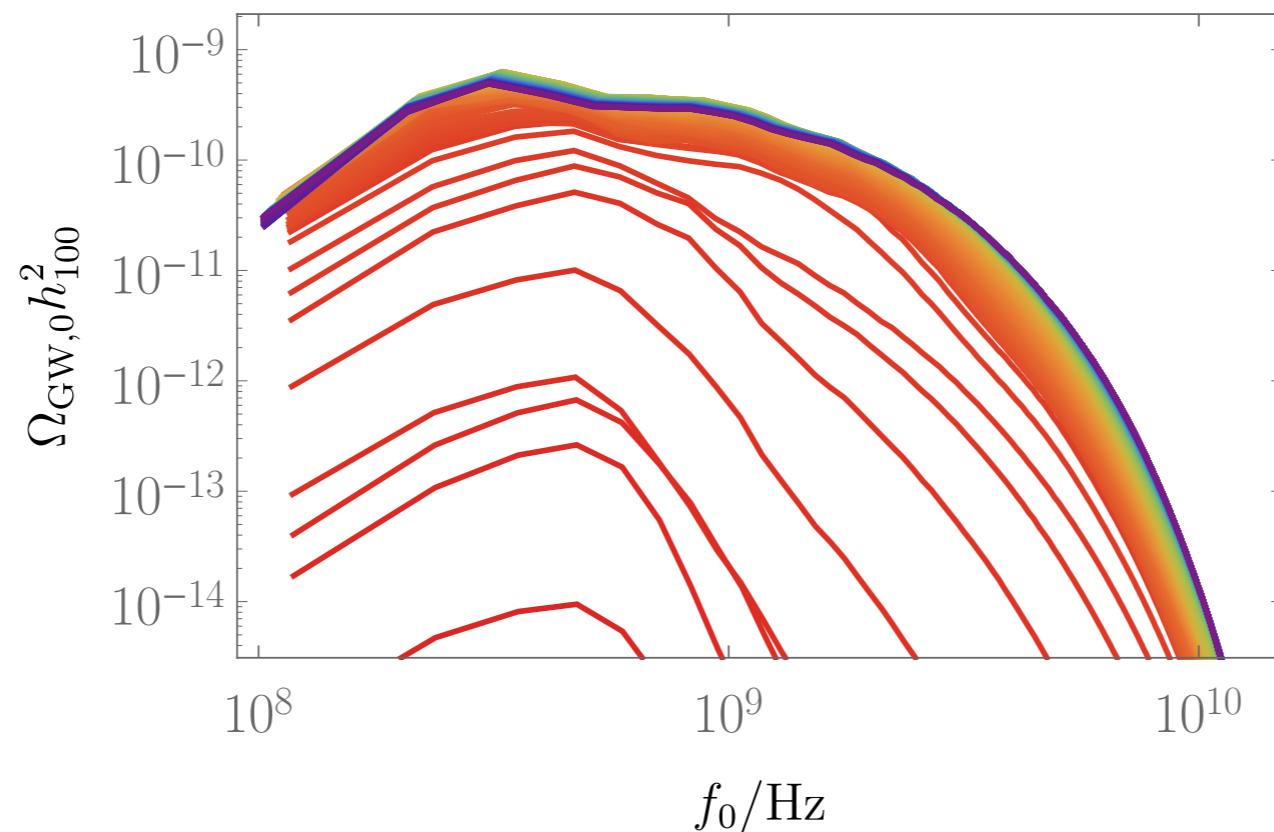
self-interactions ✓

gravitational int. ✗

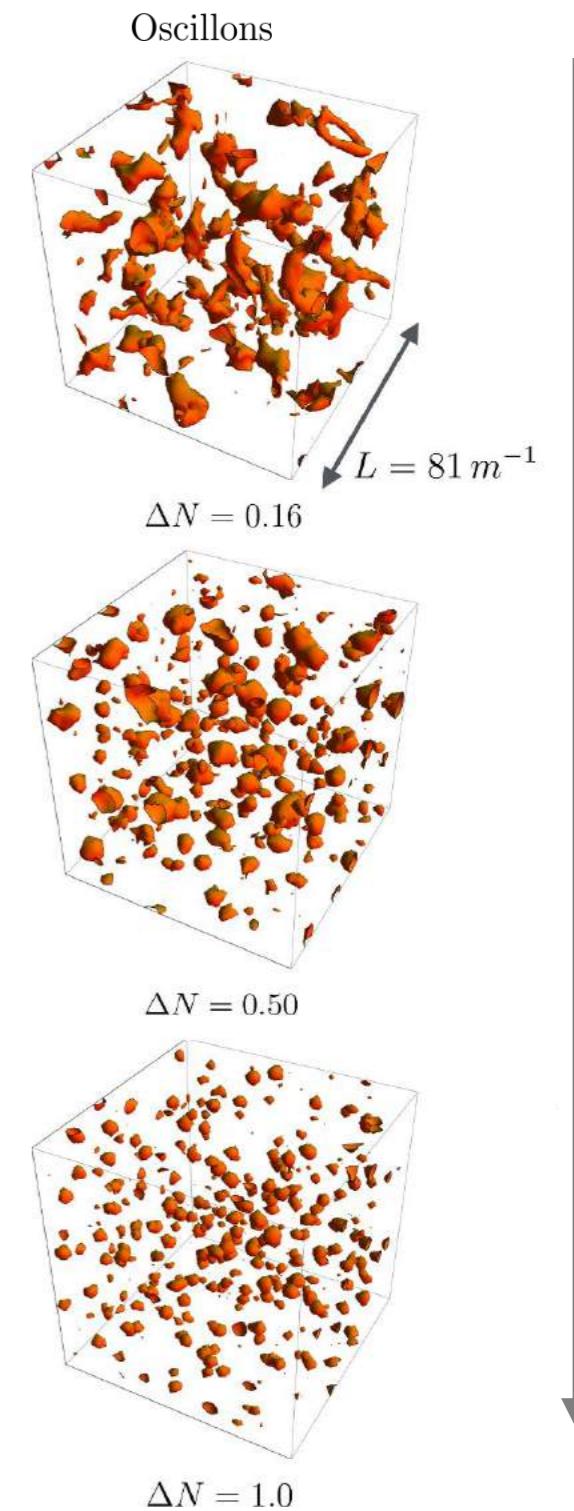
passively calculated gravitational waves



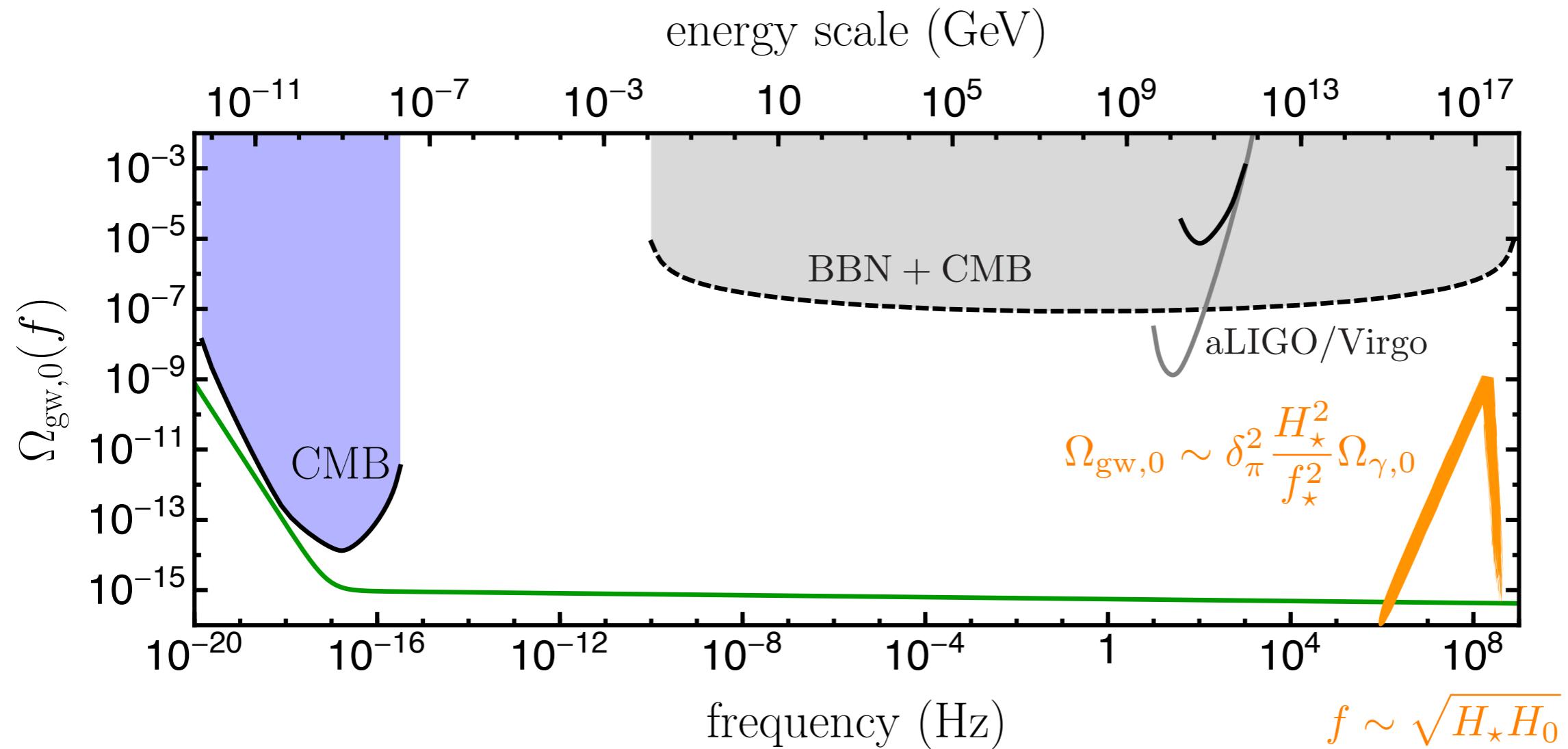
$$\Omega_{\text{GW},0} h_{100}^2 \sim 10^{-6} \left(\frac{M}{m_{\text{Pl}}} \right)^2 \lesssim \mathcal{O}[10^{-9}]$$



$$\frac{\text{growth-rate of fluctuations}}{\text{expansion rate}} \sim \frac{m_{\text{Pl}}}{M} \gg 1$$



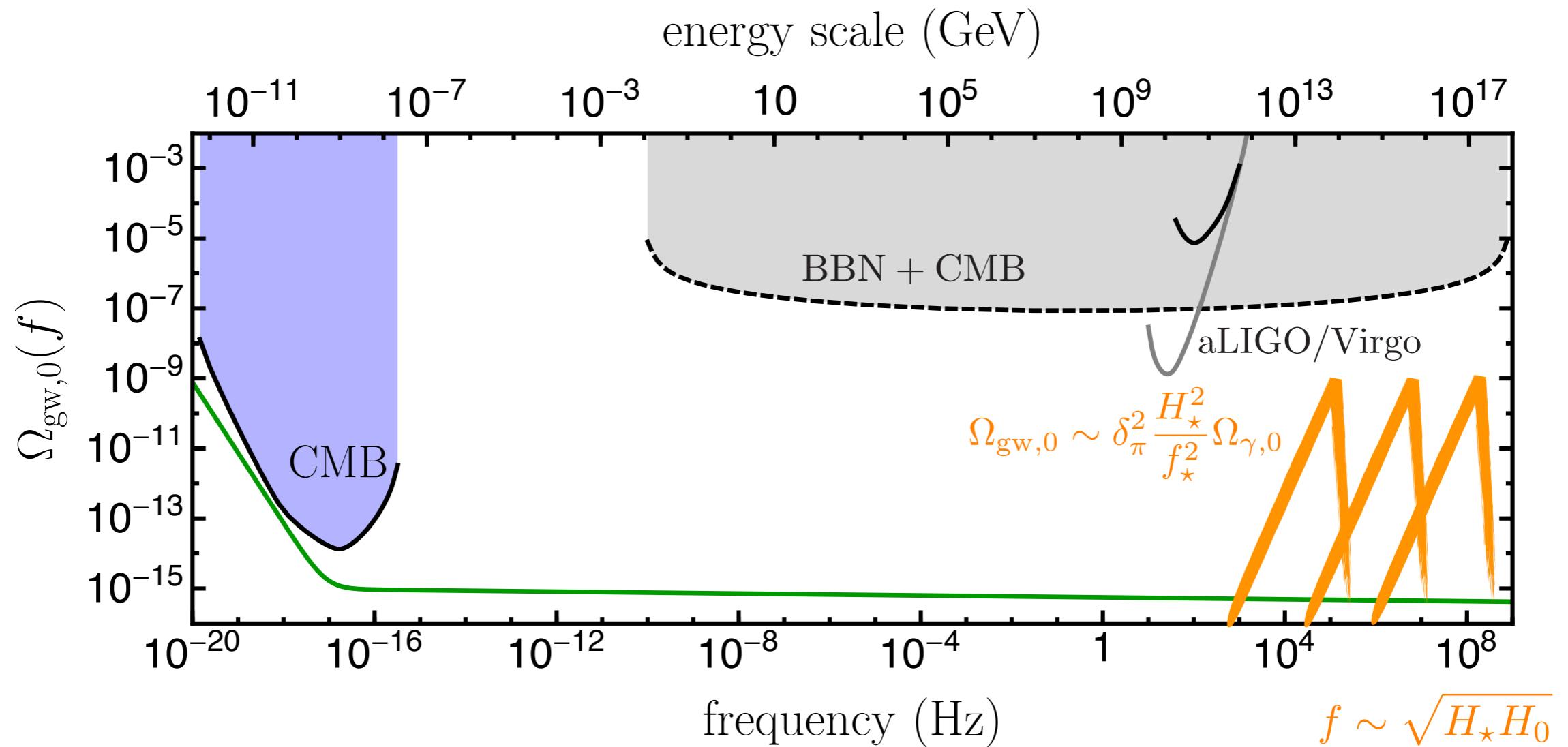
can be constrained ?



$$\Omega_{\text{gw}}(f) \sim \left[e^{-N_{\text{rad}}(1-3w)} a_{\text{eq}} \right] \delta_\pi^2 (H_*/f_*)^2$$

also see Lasky et.al (2016)

can be constrained ?



$$\Omega_{\text{gw},0} h^2 \lesssim 1.12 - 1.68 \times 10^{-7} \quad (\text{CMB} - \text{S4})$$

expansion ✓

self-interactions ✓

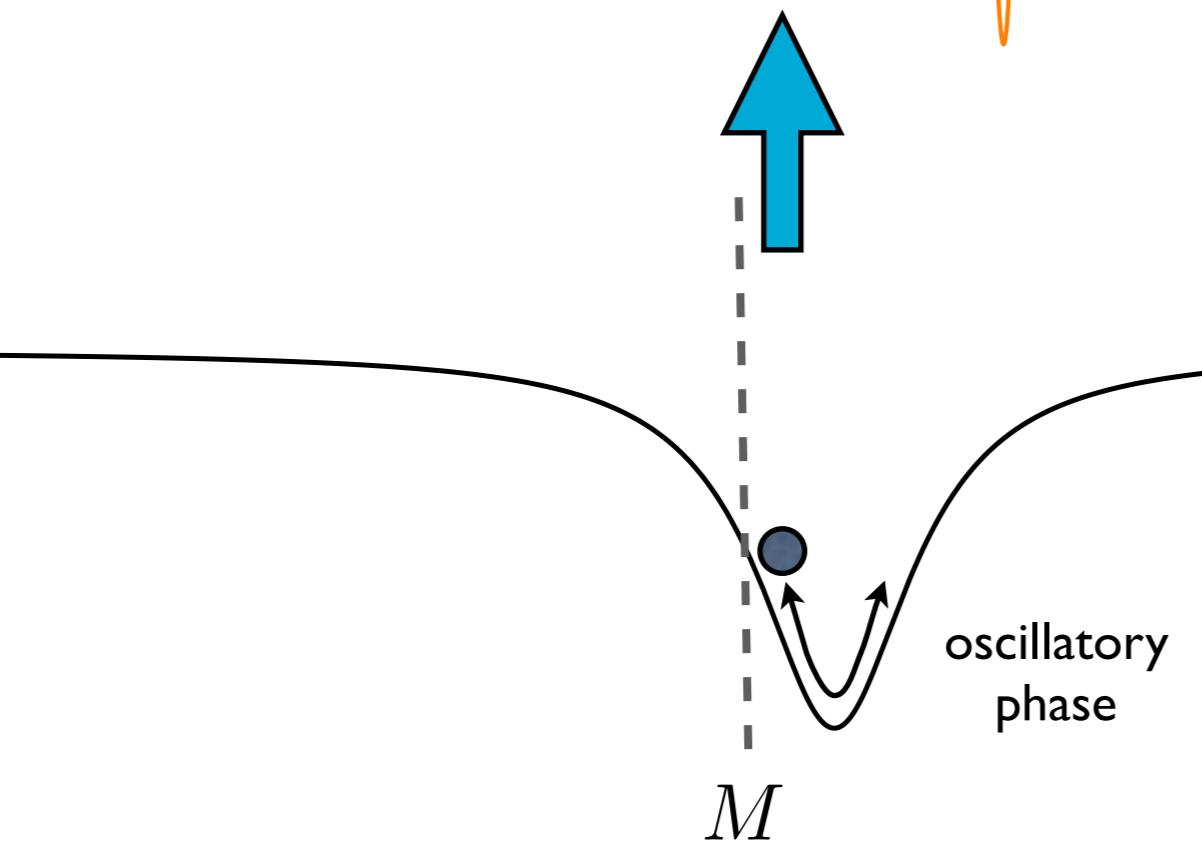
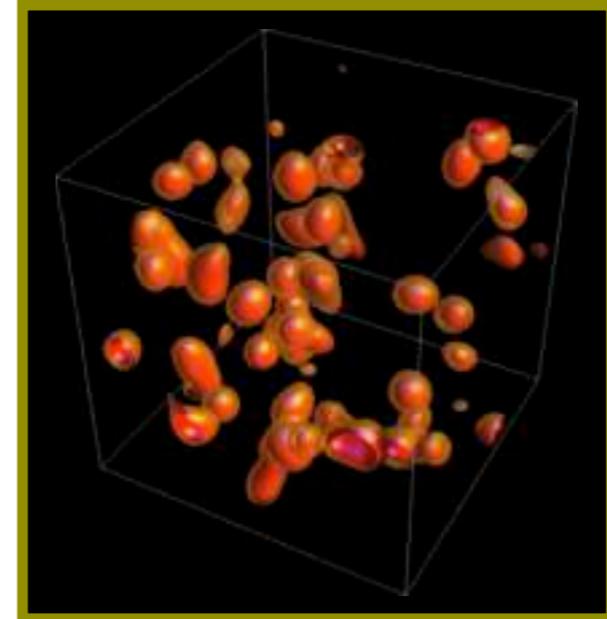
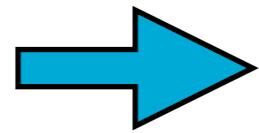
gravitational int. ✗

summary so far ...

$$\frac{\mu_k}{H} \sim \frac{m_{\text{pl}}}{M} \gg 1$$

$$\delta\varphi_k(t) \propto e^{\mu_k t}$$

resonant growth



$$|\Phi|_{\text{sol}} \lesssim 10 \times \left(\frac{M}{m_{\text{pl}}} \right)^2$$

$$\Phi \lesssim \text{few} \times 10^{-3}$$

not easy to form
black holes
from individual
solitons*

$$\Omega_{\text{GW},0} h_{100}^2 \sim 10^{-6} \left(\frac{M}{m_{\text{pl}}} \right)^2 \lesssim \mathcal{O}[10^{-9}]$$

expansion ✓

self-interactions ✓

gravitational interactions

active gravity ?

- gravitational clustering takes time ...
- long time makes it difficult to resolve very fast oscillatory time scale

a way forward ...

- rapid oscillatory behavior of fields (integrate out ?)
- size of solitons and instability length scales $\gtrsim m^{-1}$
- gravity is weak
- **non-relativistic simulations** including local gravitational interactions ?

“non-relativistic” limit

$$\square\phi + V'(\phi) = 0$$

$$G_{\mu\nu} = \frac{1}{m_{\text{pl}}^2} T_{\mu\nu}$$

$$\phi(t, \mathbf{x}) = \frac{\sqrt{2}}{m} \Re[e^{-imt} \psi(t, \mathbf{x})]$$

$$ds^2 = -(1 + 2\Phi)dt^2 + a^2(t)(1 - 2\Phi)d\mathbf{x}^2$$

$$\frac{|\nabla|}{m}, \frac{\partial_t}{m} \ll 1$$

$$|\Phi| \ll 1$$



non-linear Schrodinger eq.



Poisson eq. + Friedmann eq.

$$V(\phi) = \frac{1}{2}m^2\phi^2 + V_{\text{nl}}(\phi)$$

non-relativistic case

$$\left[i \left(\partial_t + \frac{3}{2} H \right) + \frac{1}{2a^2} \nabla^2 - U'_{\text{nl}}(|\psi|^2) - \Phi \right] \psi = 0, \quad \text{nonlinear Schrodinger eq.}$$

$$\phi(t, \mathbf{x}) = \frac{\sqrt{2}}{m} \Re [e^{-imt} \psi(t, \mathbf{x})]$$

non-relativistic case

$$\left[i \left(\partial_t + \frac{3}{2} H \right) + \frac{1}{2a^2} \nabla^2 - U'_{\text{nl}}(|\psi|^2) - \Phi \right] \psi = 0 ,$$

nonlinear Schrodinger eq.

$$\frac{\nabla^2}{a^2} \Phi = \frac{\beta^2}{2} \left[|\psi|^2 + \frac{1}{2a^2} |\nabla \psi|^2 + U_{\text{nl}}(|\psi|^2) \right] - \frac{3}{2} H^2 ,$$

Poisson eq.

$$H^2 = \frac{\beta^2}{3} \overline{\left[|\psi|^2 + \frac{1}{2a^2} |\nabla \psi|^2 + U_{\text{nl}}(|\psi|^2) \right]} ,$$

Friedmann eq.

$$m x^\mu \rightarrow x^\mu$$

$$\frac{\psi}{mM} \rightarrow \psi$$

$$\beta \equiv \frac{M}{m_{\text{pl}}}$$

length/time units

non-linearity

$$\phi(t, \mathbf{x}) = \frac{\sqrt{2}}{m} \Re [e^{-imt} \psi(t, \mathbf{x})]$$

non-relativistic case

$$\left[i \left(\partial_t + \frac{3}{2} H \right) + \frac{1}{2a^2} \nabla^2 - U'_{\text{nl}}(|\psi|^2) - \Phi \right] \psi = 0,$$

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length/time units

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$$\phi(t, \mathbf{x}) = \frac{\sqrt{2}}{m} \Re [e^{-imt} \psi(t, \mathbf{x})]$$

relativistic*

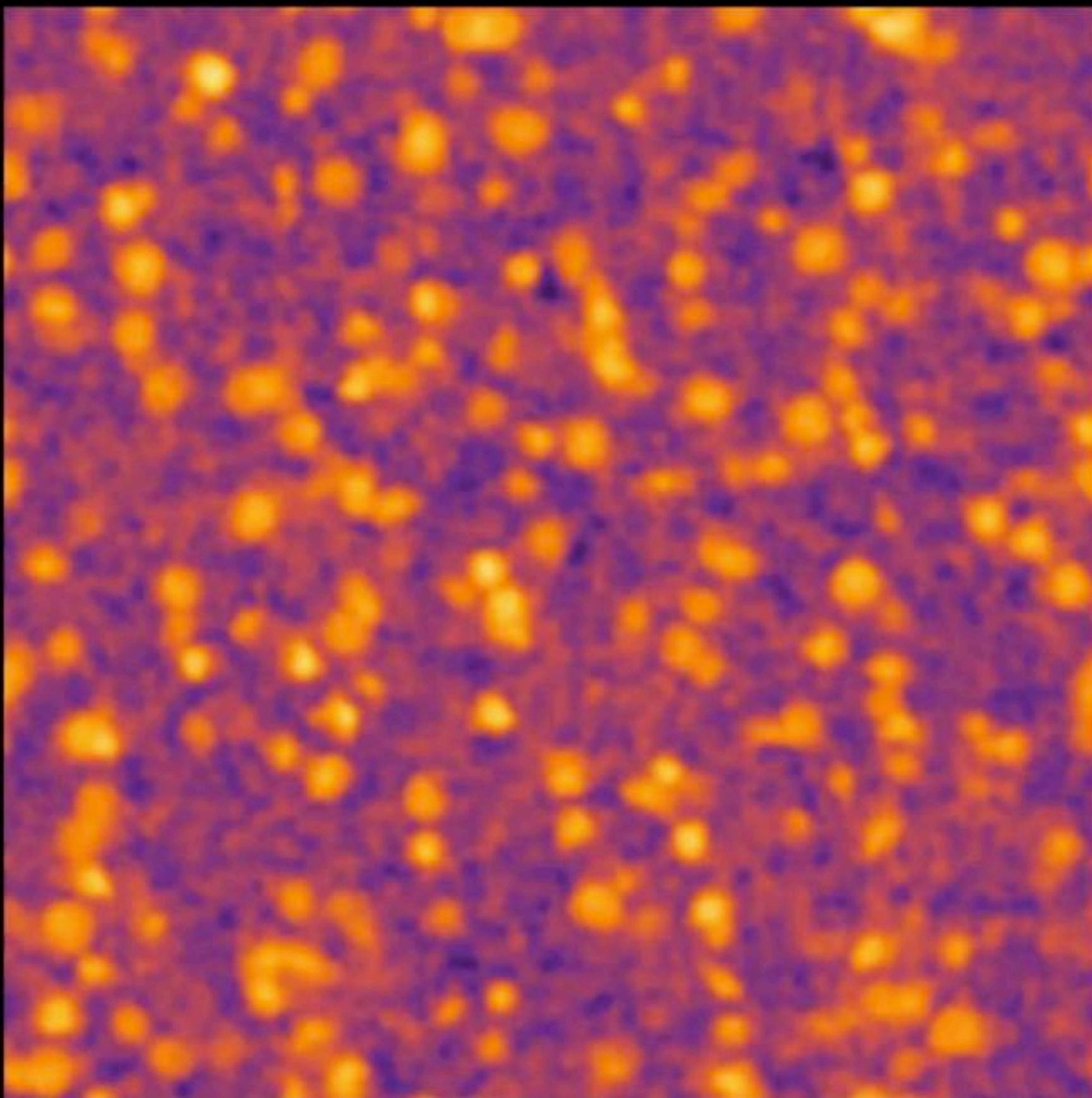


self-interactions



gravitational int. ✗

$a=3.12$



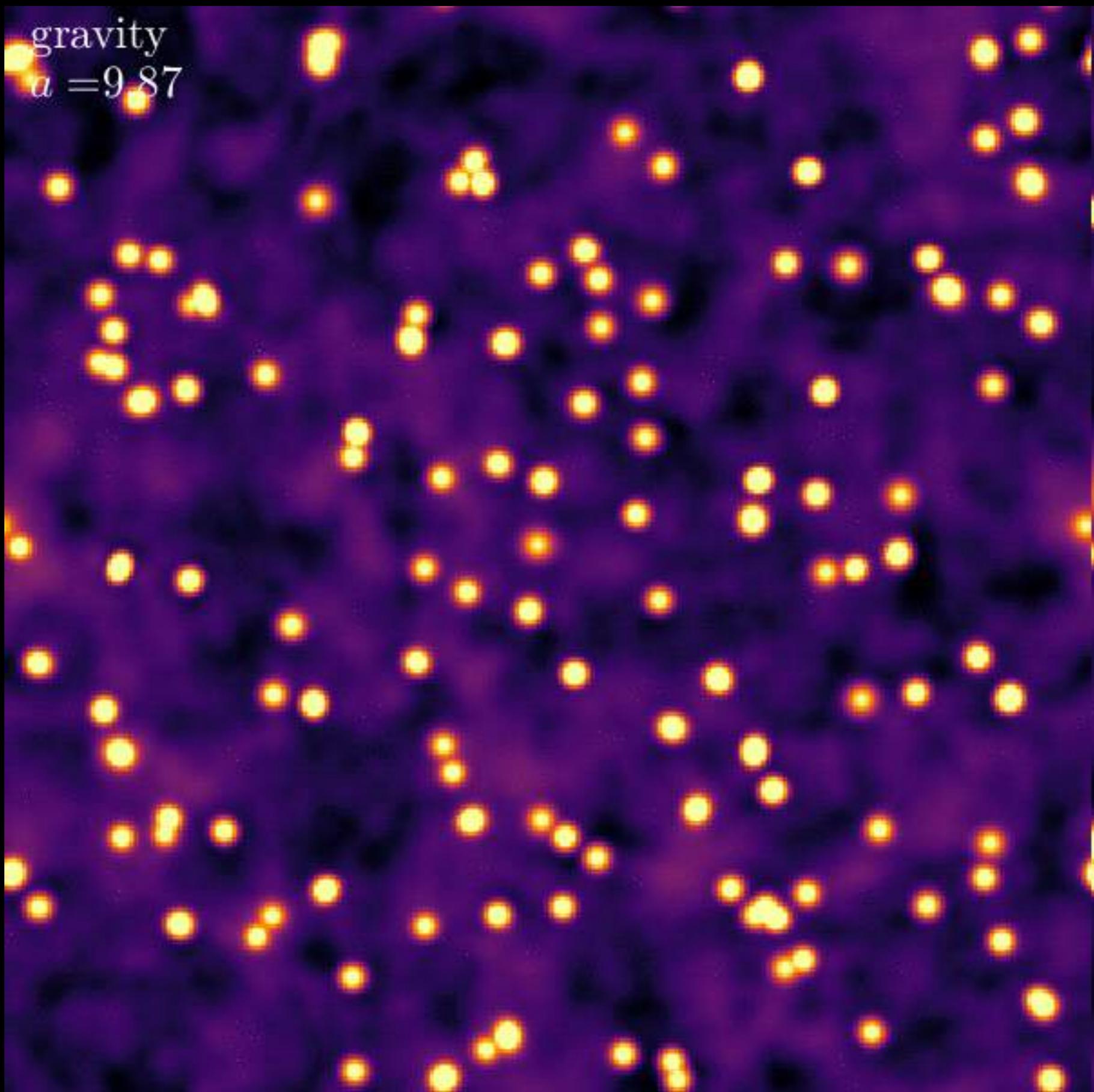
relativistic*



self-interactions



gravitational int.



two linear instabilities

$$\psi(t, \mathbf{x}) = \bar{\psi}(t) \left[1 + \varepsilon \frac{\delta\psi_{\mathbf{k}}(t)}{\bar{\psi}(t)} e^{i\mathbf{k}\cdot\mathbf{x}} \right]$$

$$|\delta\psi_{\mathbf{k}}/\bar{\psi}| \sim e^{\mu_k t}$$

self-interactions

$$k^2 < -4|\bar{\psi}|^2 U''_{\text{nl}}(|\bar{\psi}|^2),$$

$$\mu_k = \left| i \frac{k}{2} \sqrt{k^2 + 4|\bar{\psi}|^2 U''_{\text{nl}}(|\bar{\psi}|^2)} \right|$$

$$\frac{\mu_k}{H} \sim \frac{1}{\beta}$$

gravitational interactions

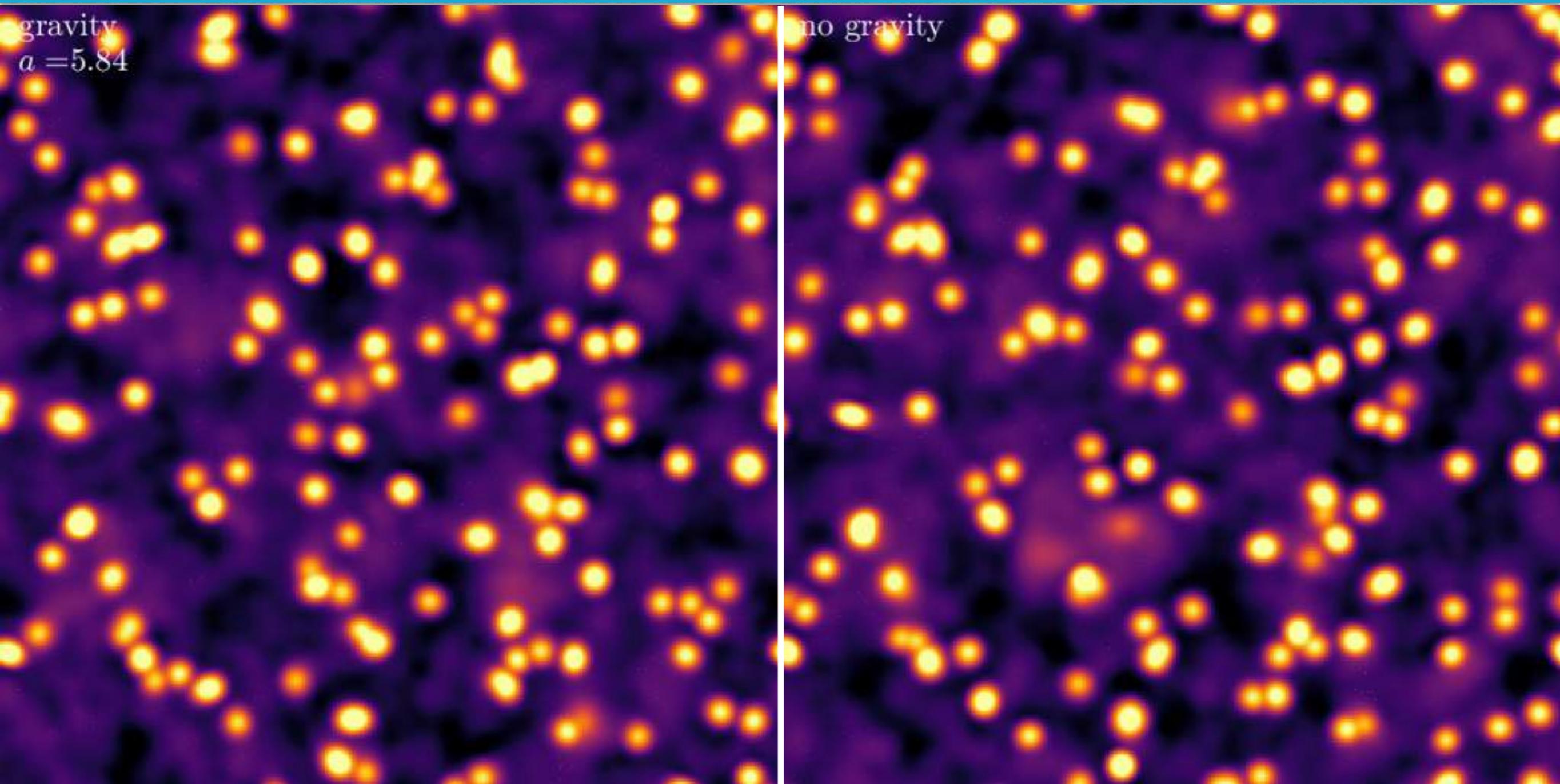
$$k < k_J \approx \sqrt{\sqrt{2}\beta|\bar{\psi}|}$$

$$\mu_k = \sqrt{\frac{1}{2}\beta^2|\bar{\psi}|^2 - \frac{k^4}{4}}$$

$$\frac{\mu_k}{H} \sim 1$$

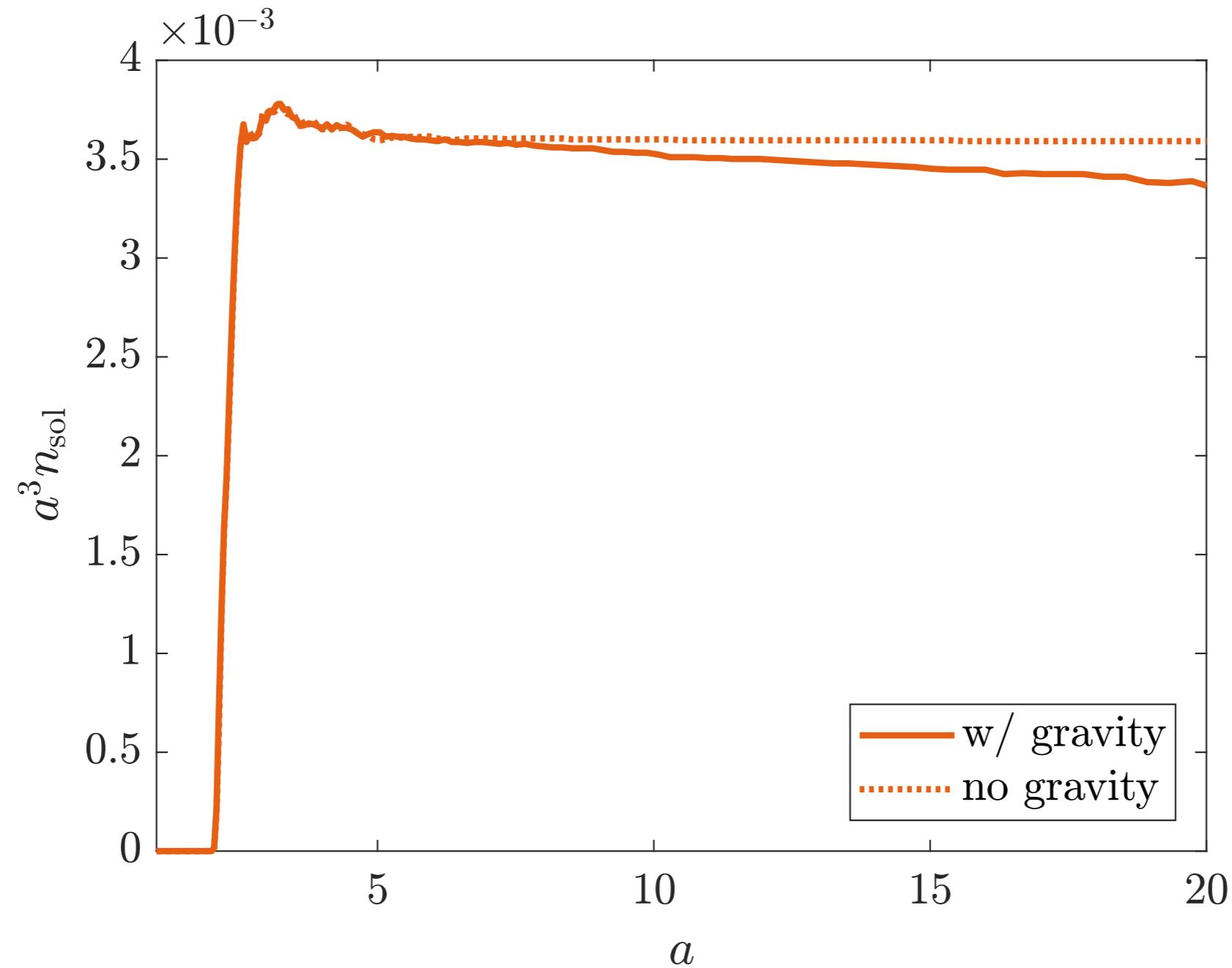
For $\beta \equiv \frac{M}{m_{\text{pl}}} \ll 1$ self-interaction instability dominates

perturbations become nonlinear & form solitons



$$a^3 n_{\text{sol}} \sim (k_{\text{nl}}/2\pi)^3$$

co-moving number density of solitons

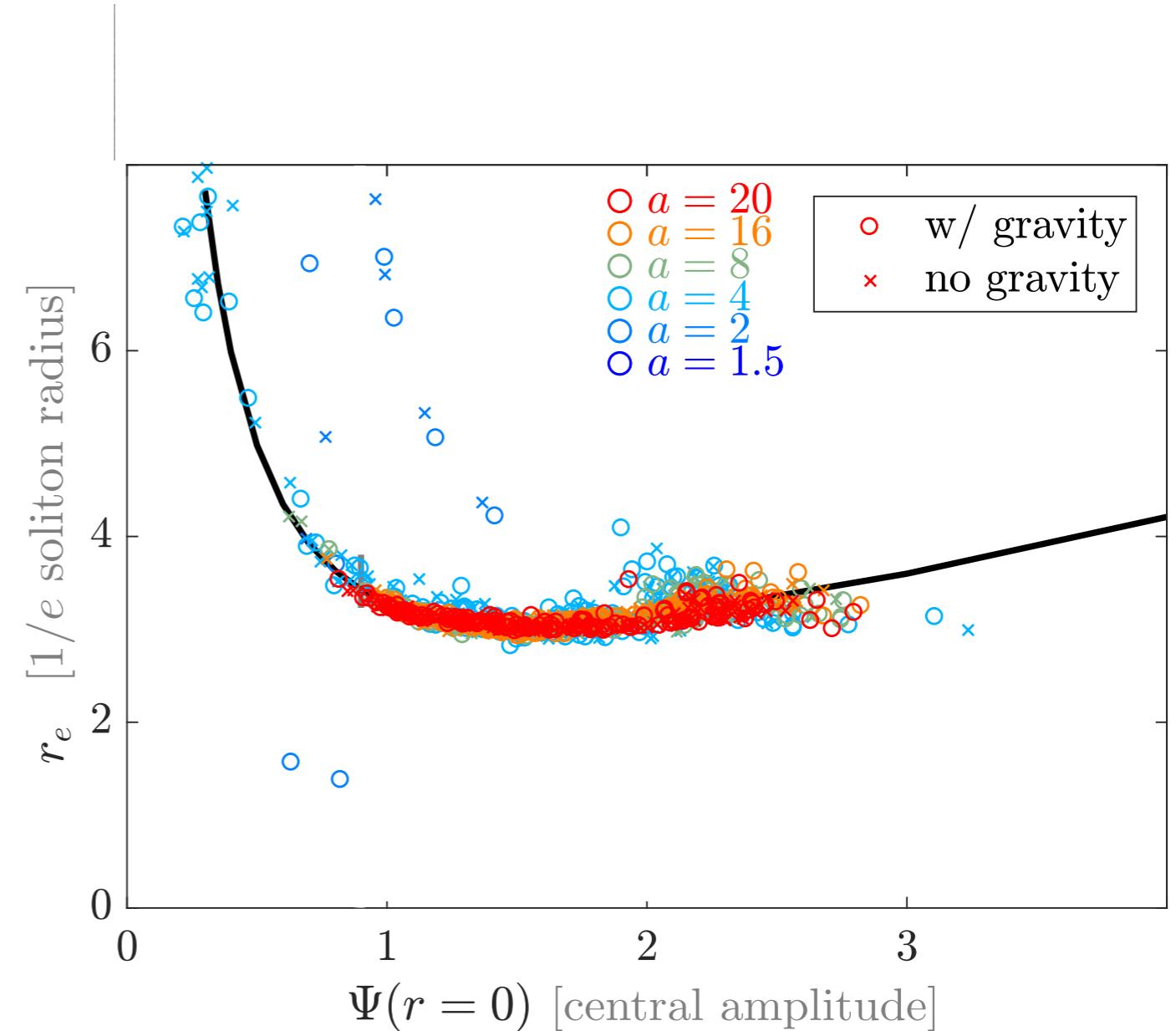
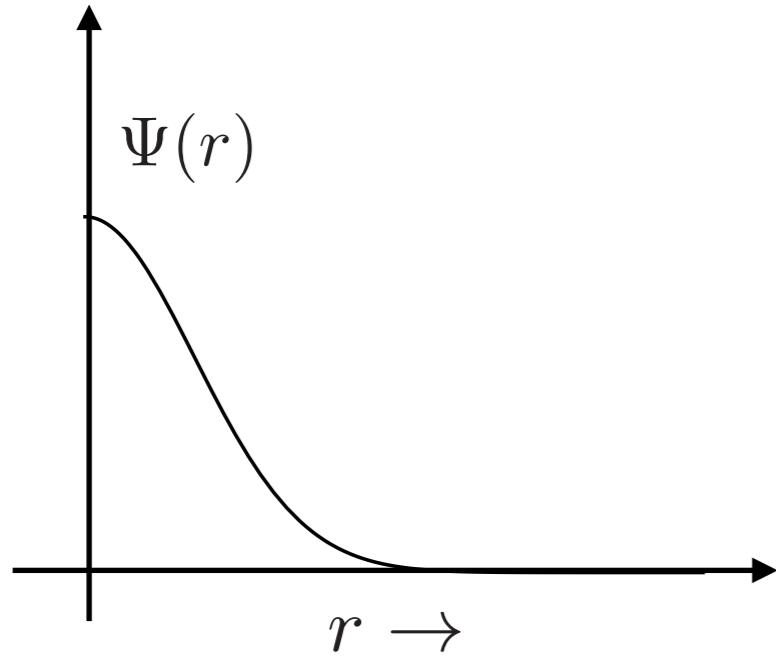


$$\frac{d \ln(a^3 n_{\text{sol}})}{d \ln a} \simeq 0.1$$

few $\times 10^3$ /Hubble volume at time of formation

individual solitons

$$\psi(t, r) = e^{-i\nu t} \Psi(r)$$



$$\phi(t, r) = \sqrt{2} \Re[e^{-it} \psi(t, r)] = \sqrt{2} \Psi(r) \cos[(1 + \nu)t]$$

$$\nu < 0, \quad |\nu| \ll 1$$

individual solitons stability

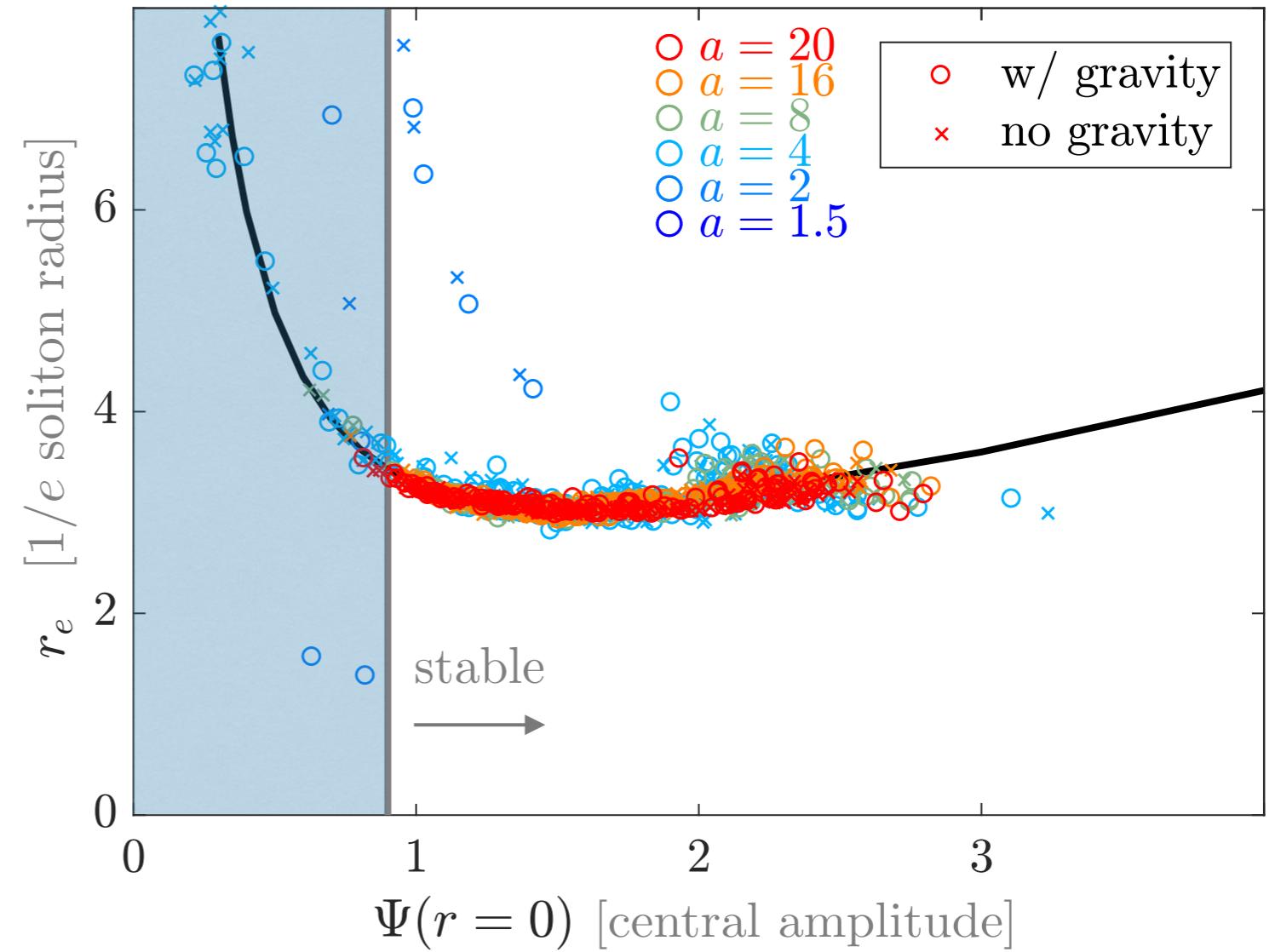
$$\psi(t, r) = e^{-i\nu t} \Psi(r)$$

$$\mathcal{N} \equiv \int d^3r \Psi^2(r)$$

stable iff:

Vakhitov Kolokolov (1973)

$$\frac{d\mathcal{N}}{d(-\nu)} > 0$$

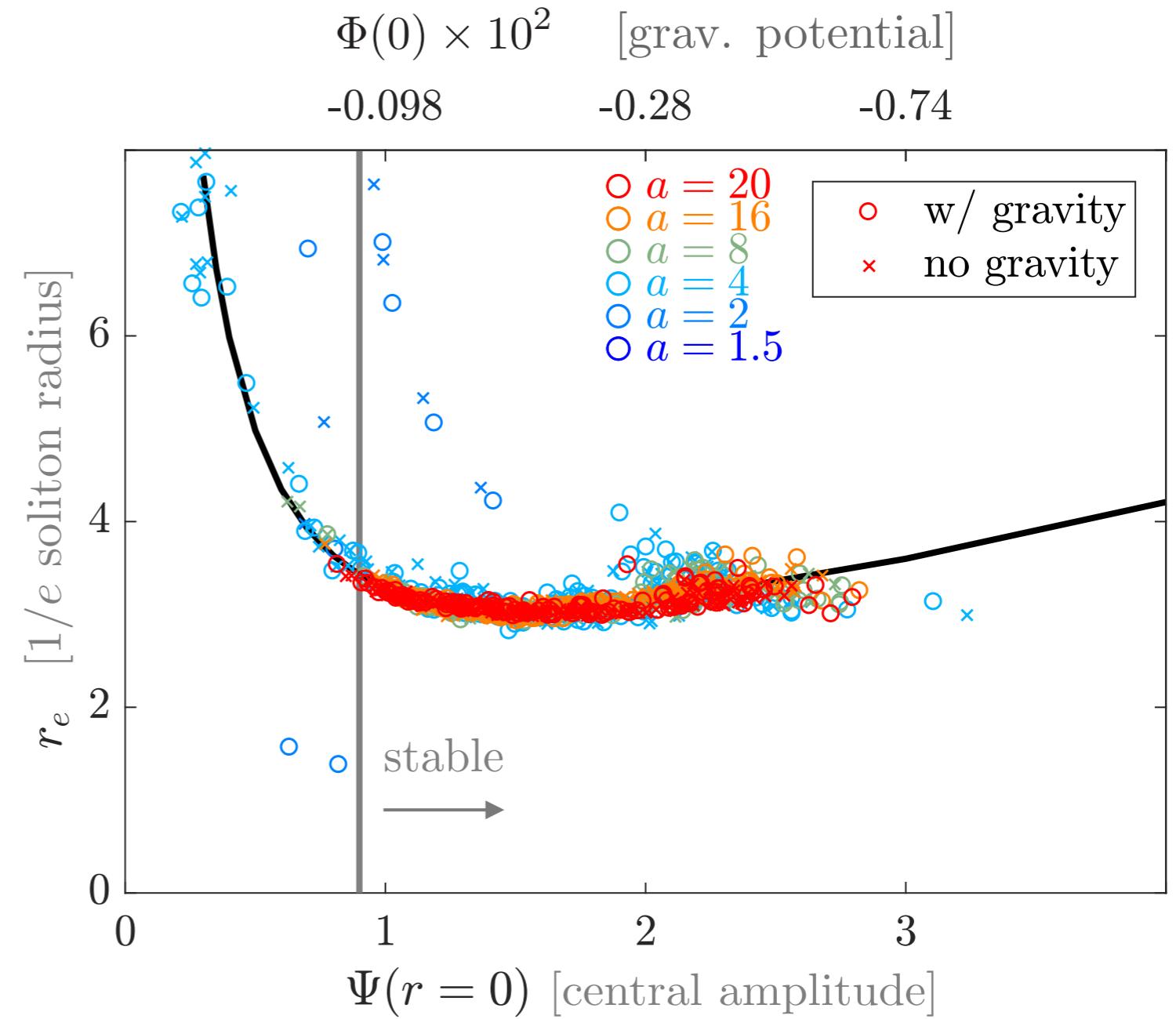


— stability with gravitational interactions needs to be investigated

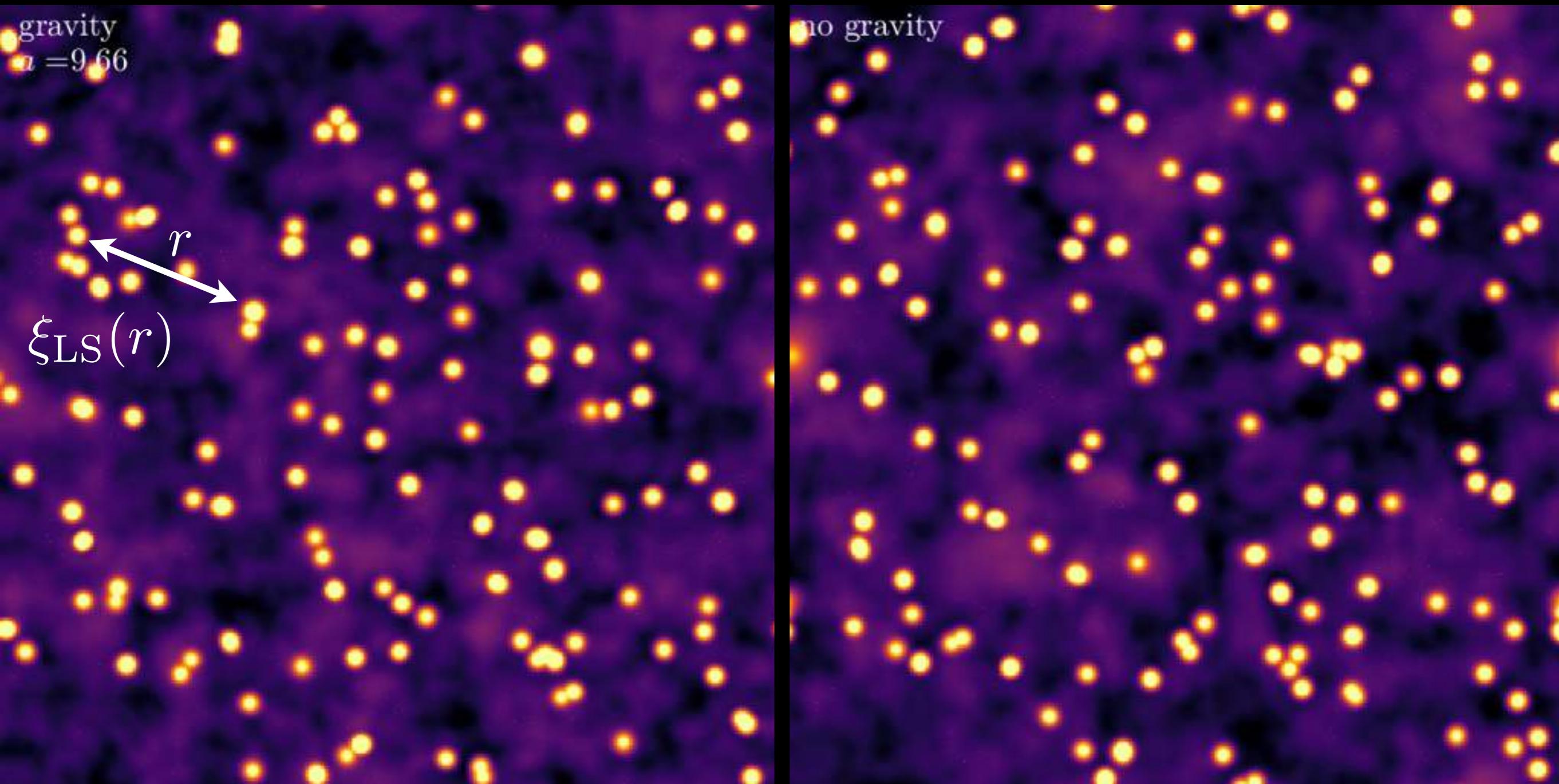
gravity remains weak

$$|\Phi| \sim \mathcal{O}[10^{-3}]$$

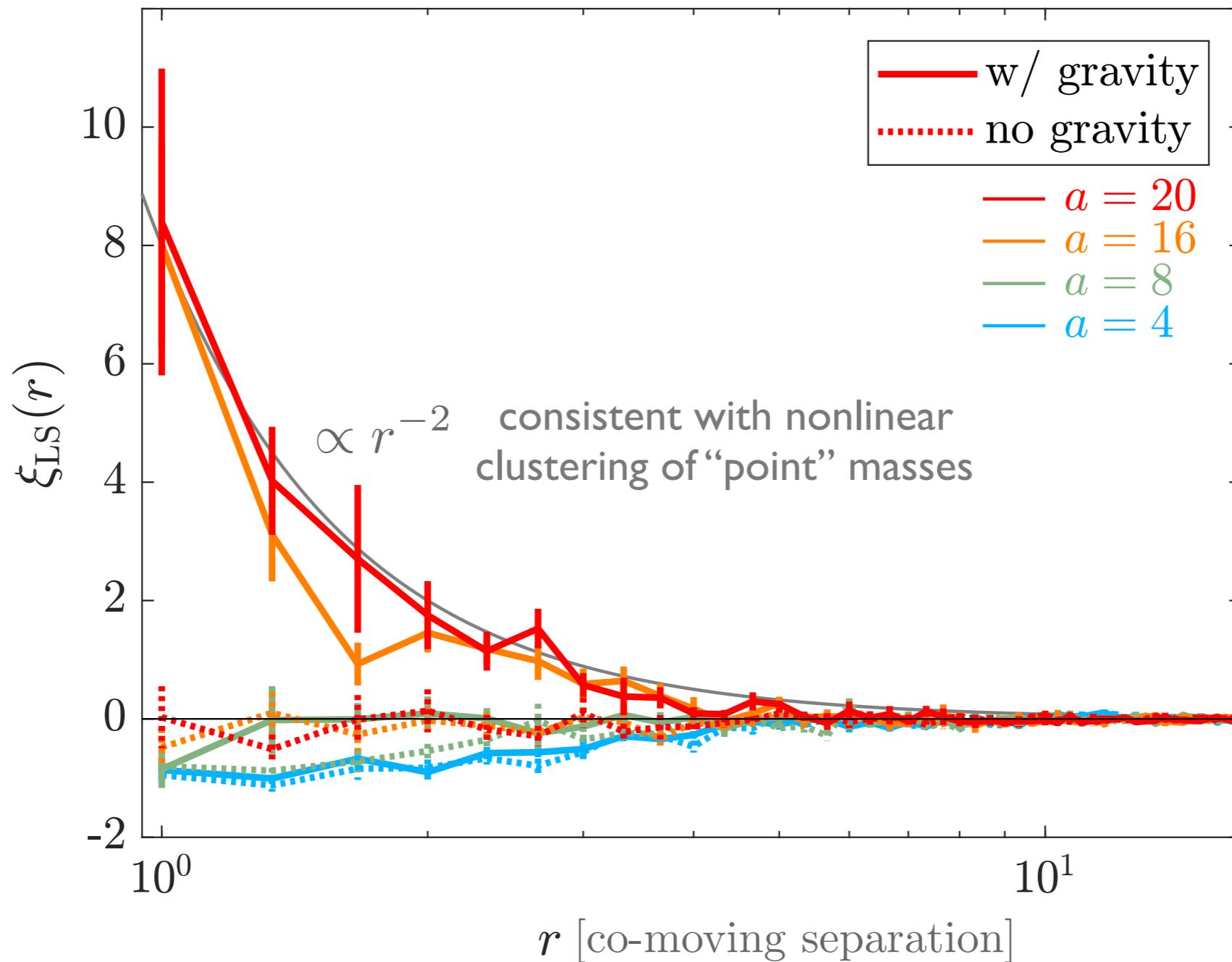
not easy to form black holes
from individual solitons still



gravitational clustering of solitons

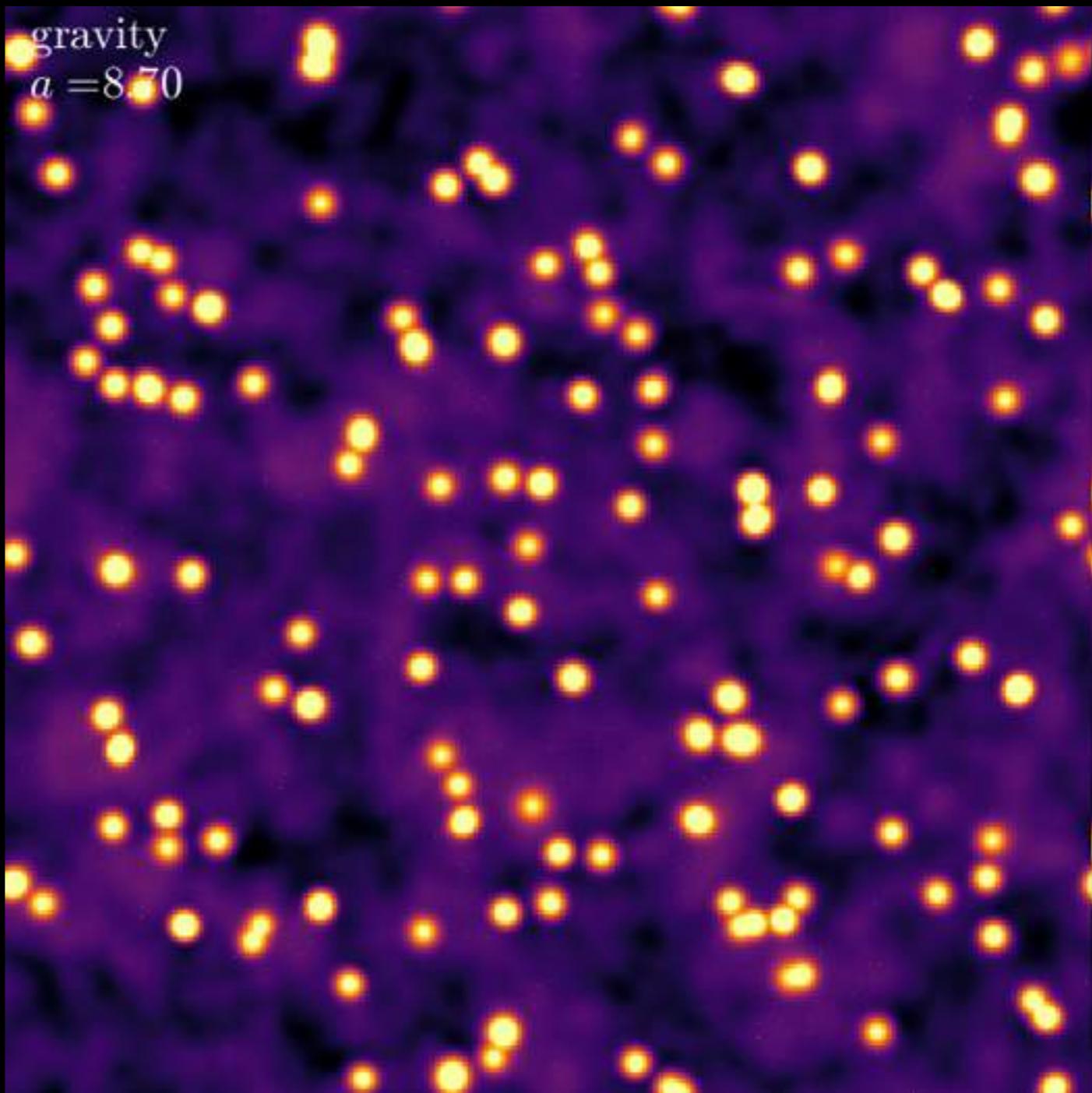
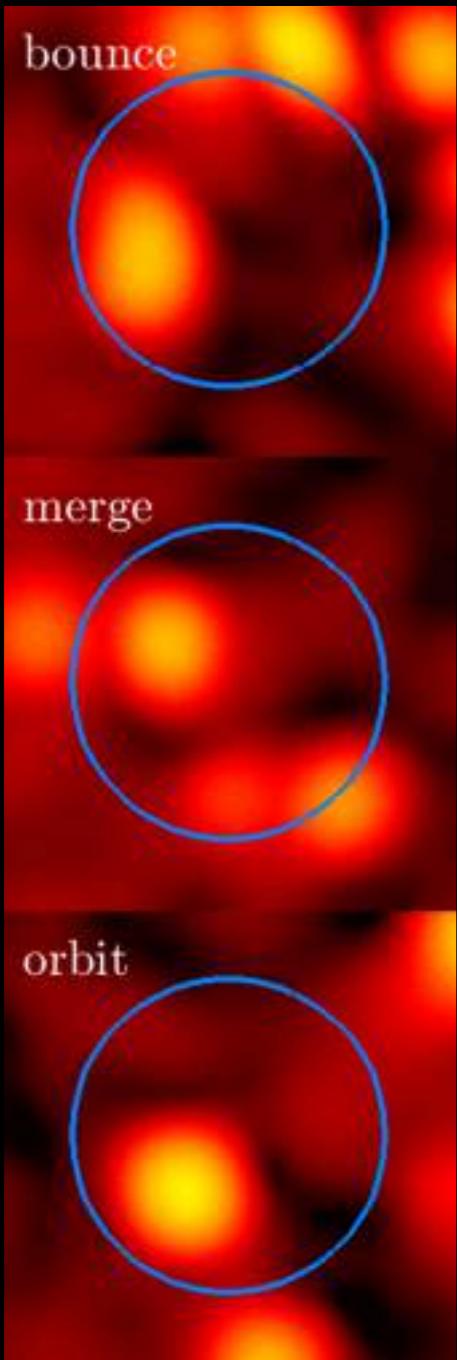


gravitational clustering of solitons

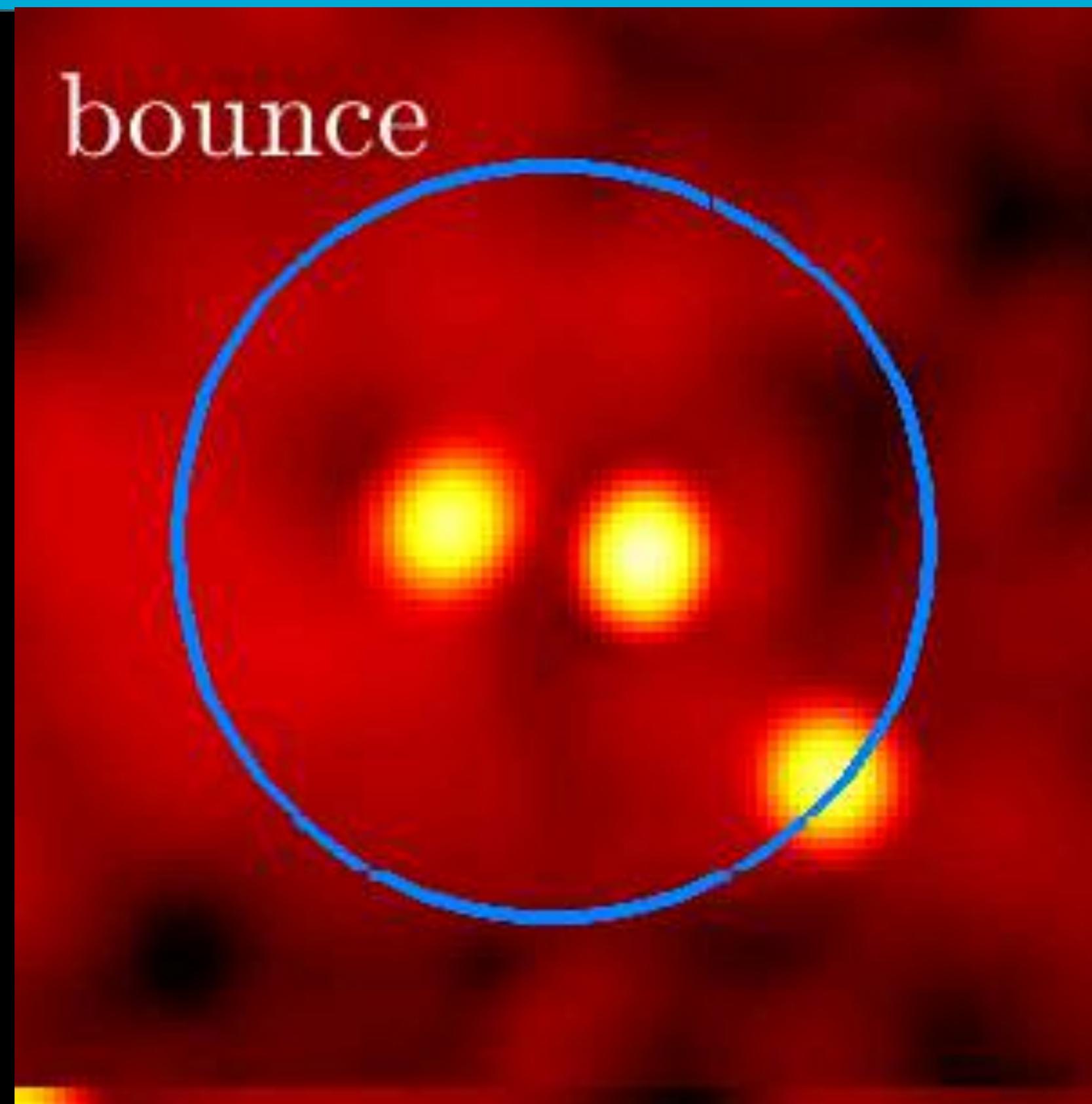


*theoretical arguments for r^2 in Saslaw 1980

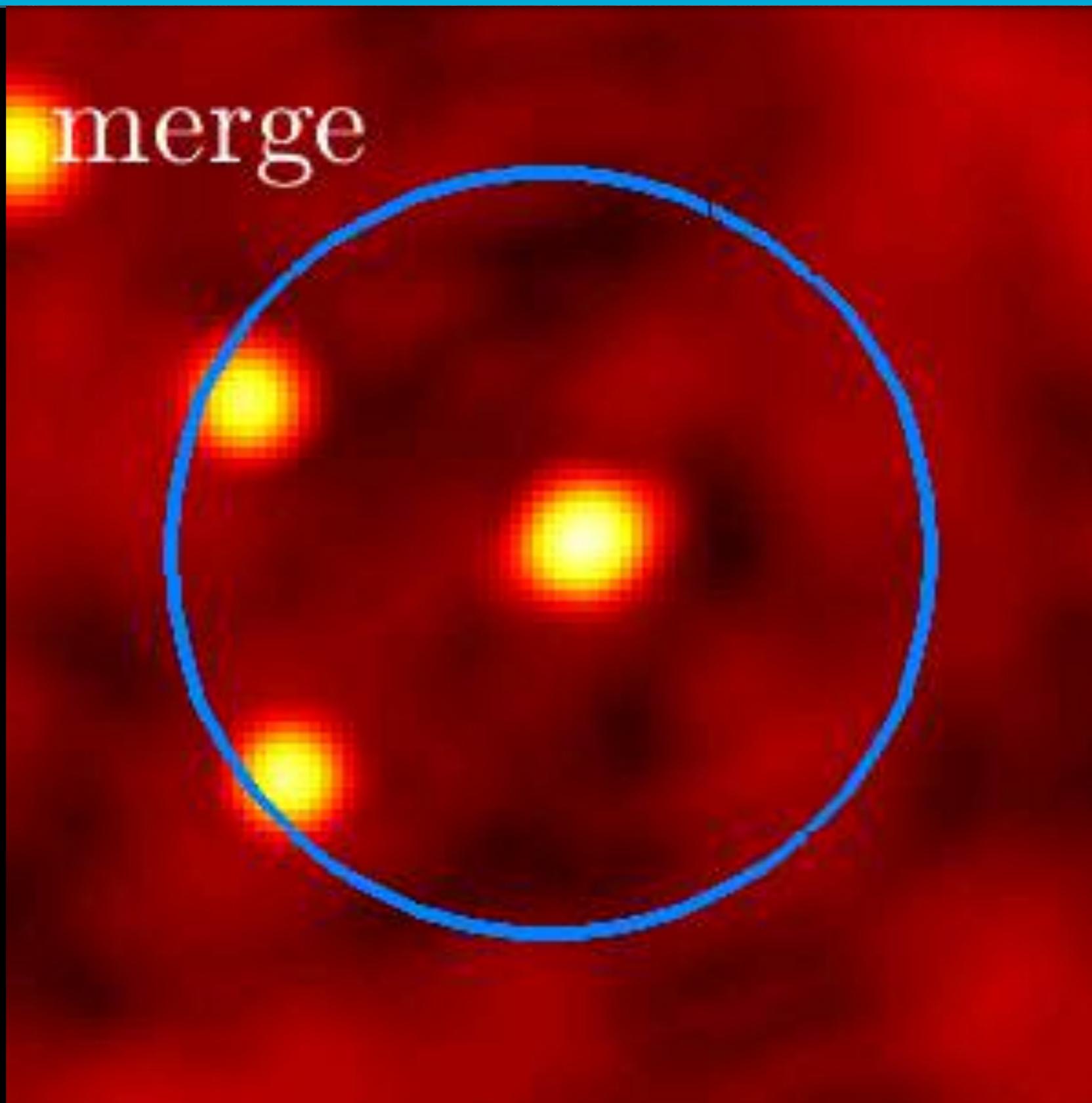
strong interactions



strong interactions: bounce

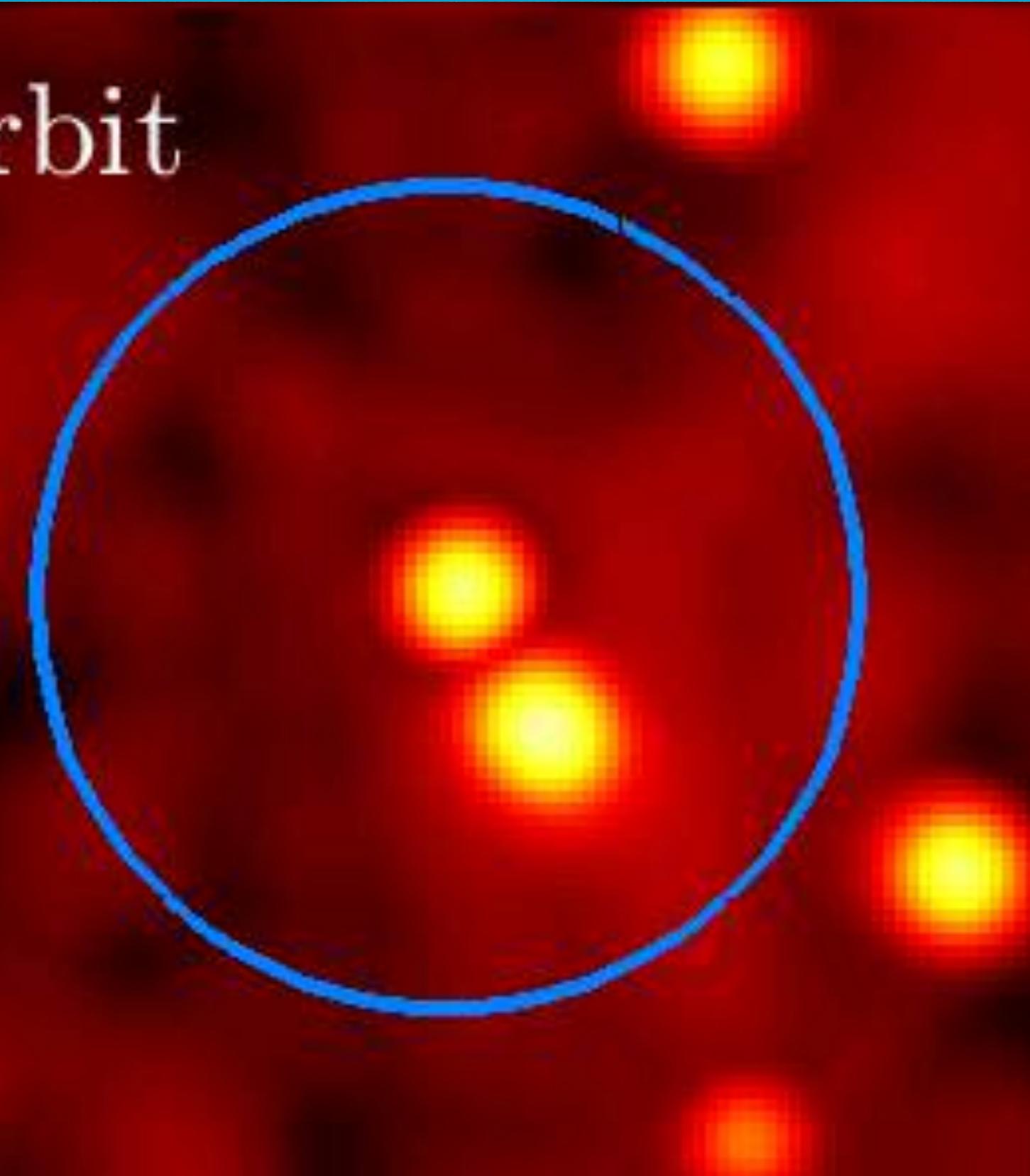


strong interactions: mergers



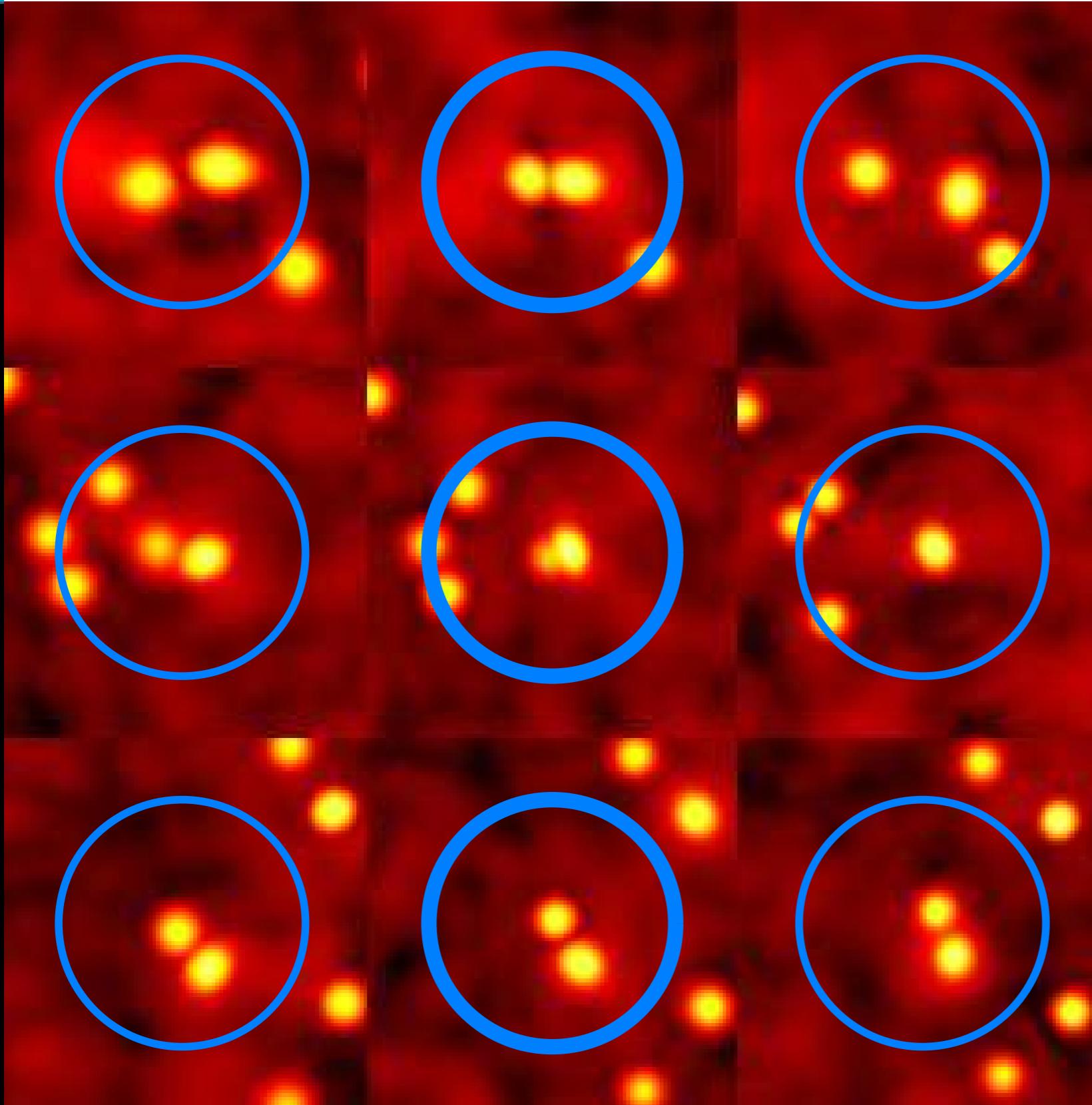
strong interactions: orbit

orbit



strong interactions

bounce



merger

“binary”

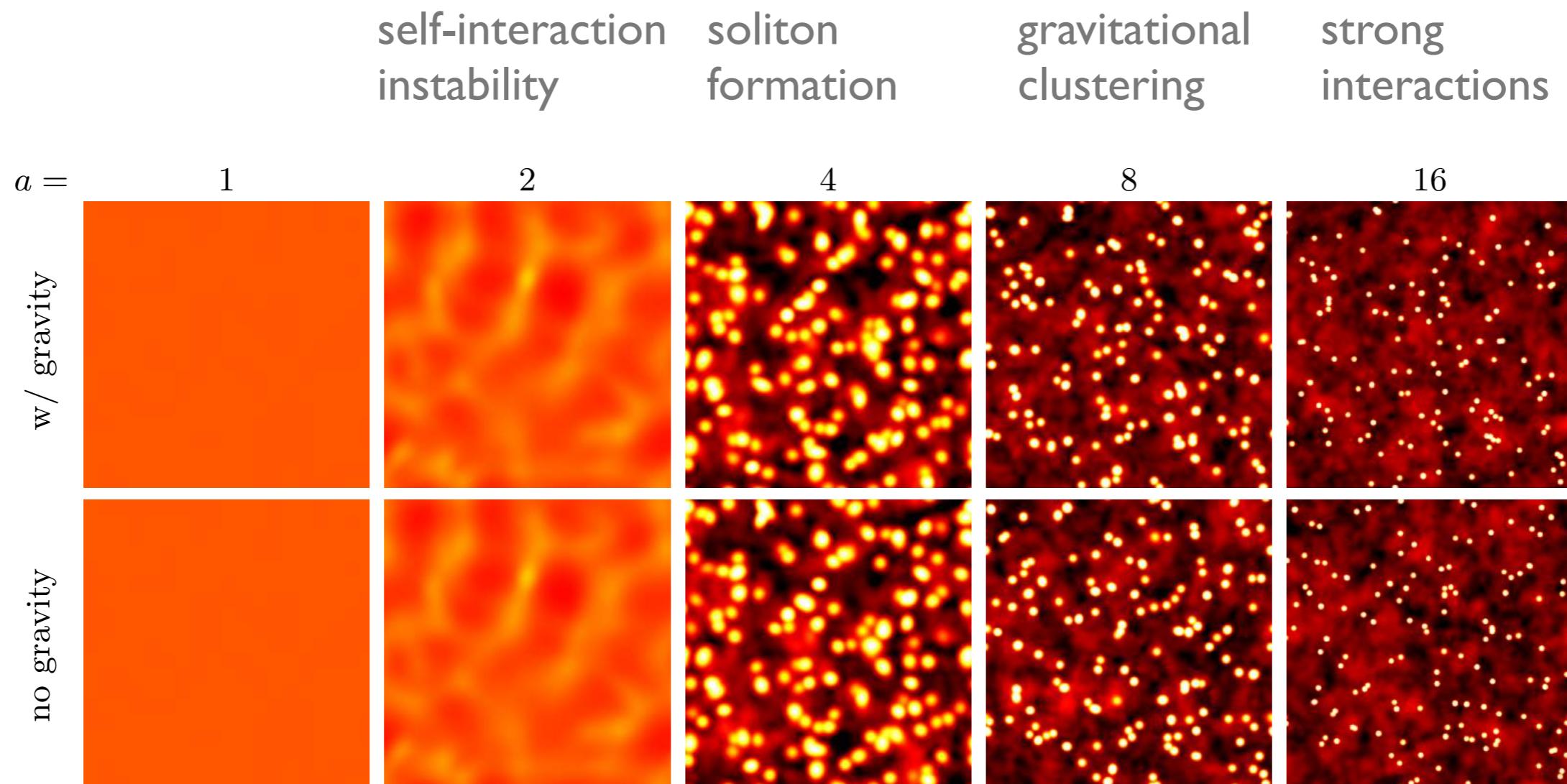
$$\psi_a(t, \mathbf{x}) = \Psi_a(\mathbf{x}) e^{-i\nu_a t + \theta_a}$$

$$|\theta_1 - \theta_2| \simeq \pi$$

$$|\theta_1 - \theta_2| \simeq 0$$

- * need better understanding
- * & simulations

summary with gravitational interactions



things that need more work

- relativistic - non-relativistic connection
- long term state of the strongly interacting soliton gas — how does probability of PBH formation increases ?
- better understanding of the strong interactions

exploring individual interactions further with full numerical GR

$$S = \int d^4x \sqrt{-g} \left[\frac{m_{\text{pl}}^2}{2} R - \frac{1}{2} (\partial\phi)^2 - V(\phi) \right]$$

$$V(\phi) = \frac{1}{2} m^2 \phi^2 + V_{\text{nl}}(\phi)$$

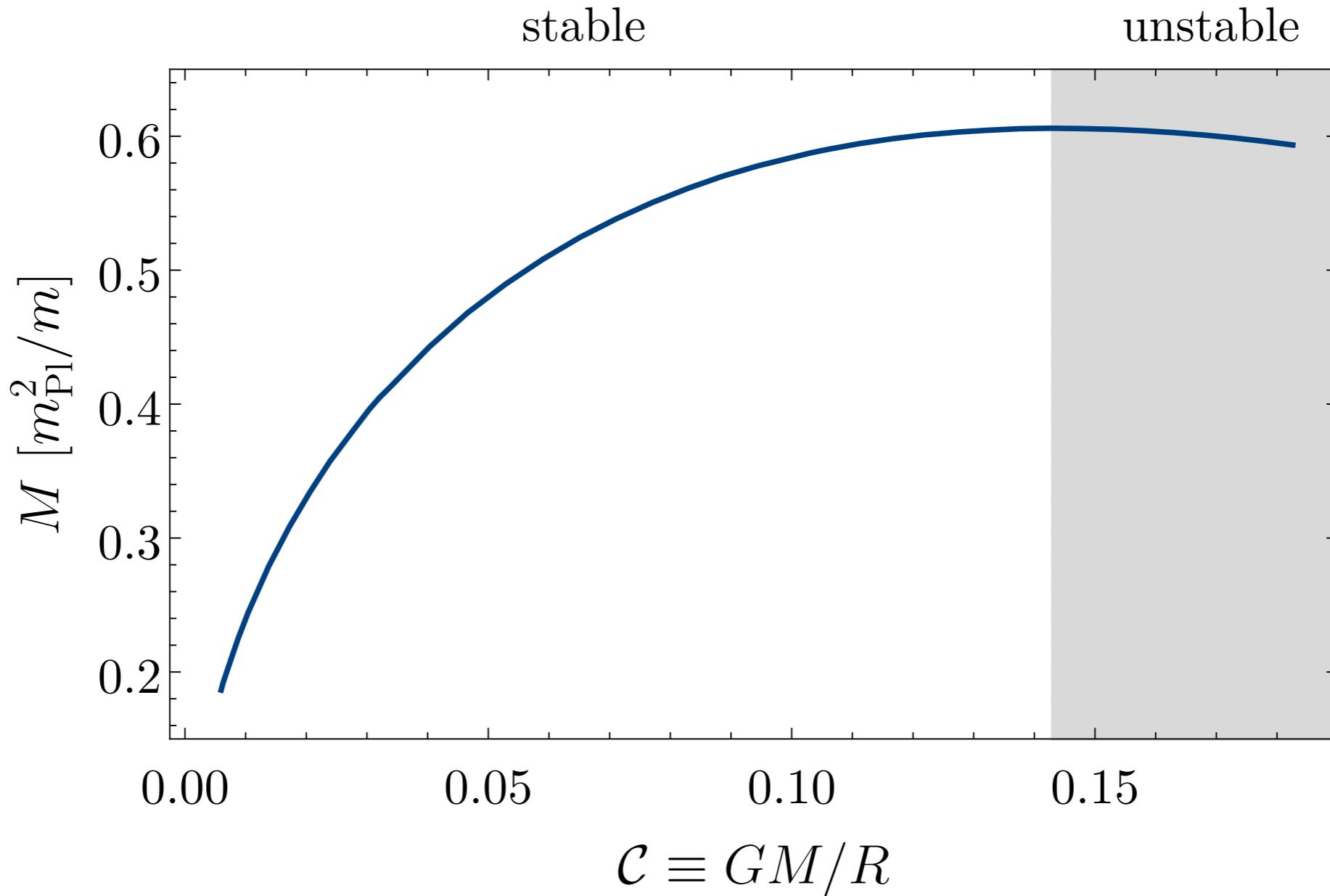
ignore self-interactions

interested in gravitational wave emission from ultra-compact solitons

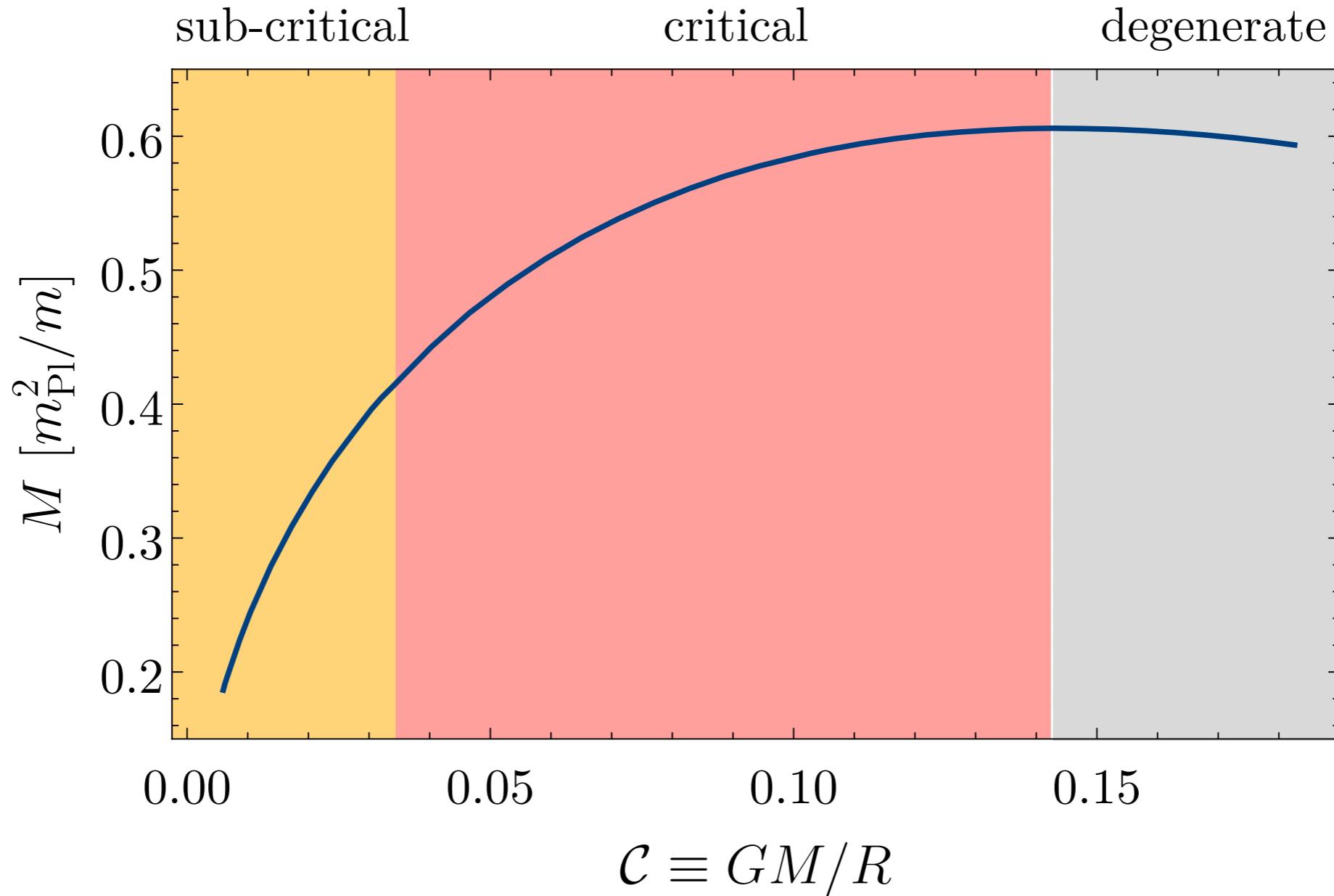
Helper, Lim, Garcia & MA (2018)



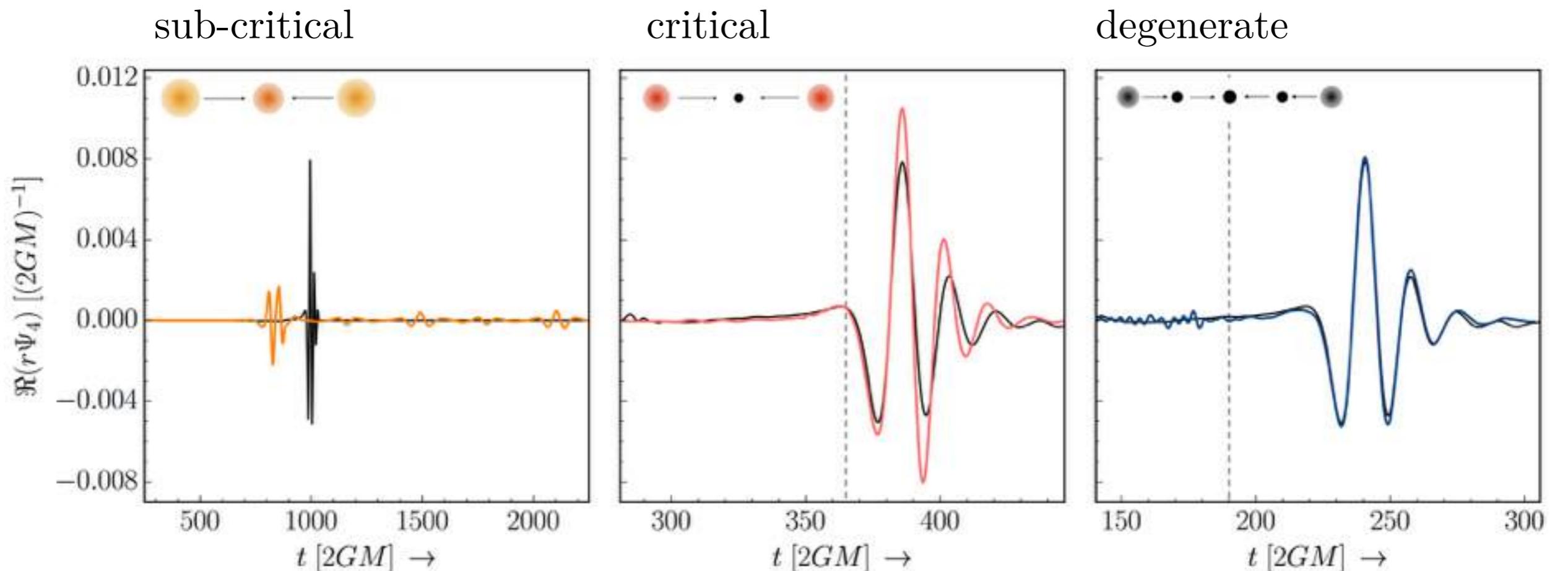
ultra-compact soliton



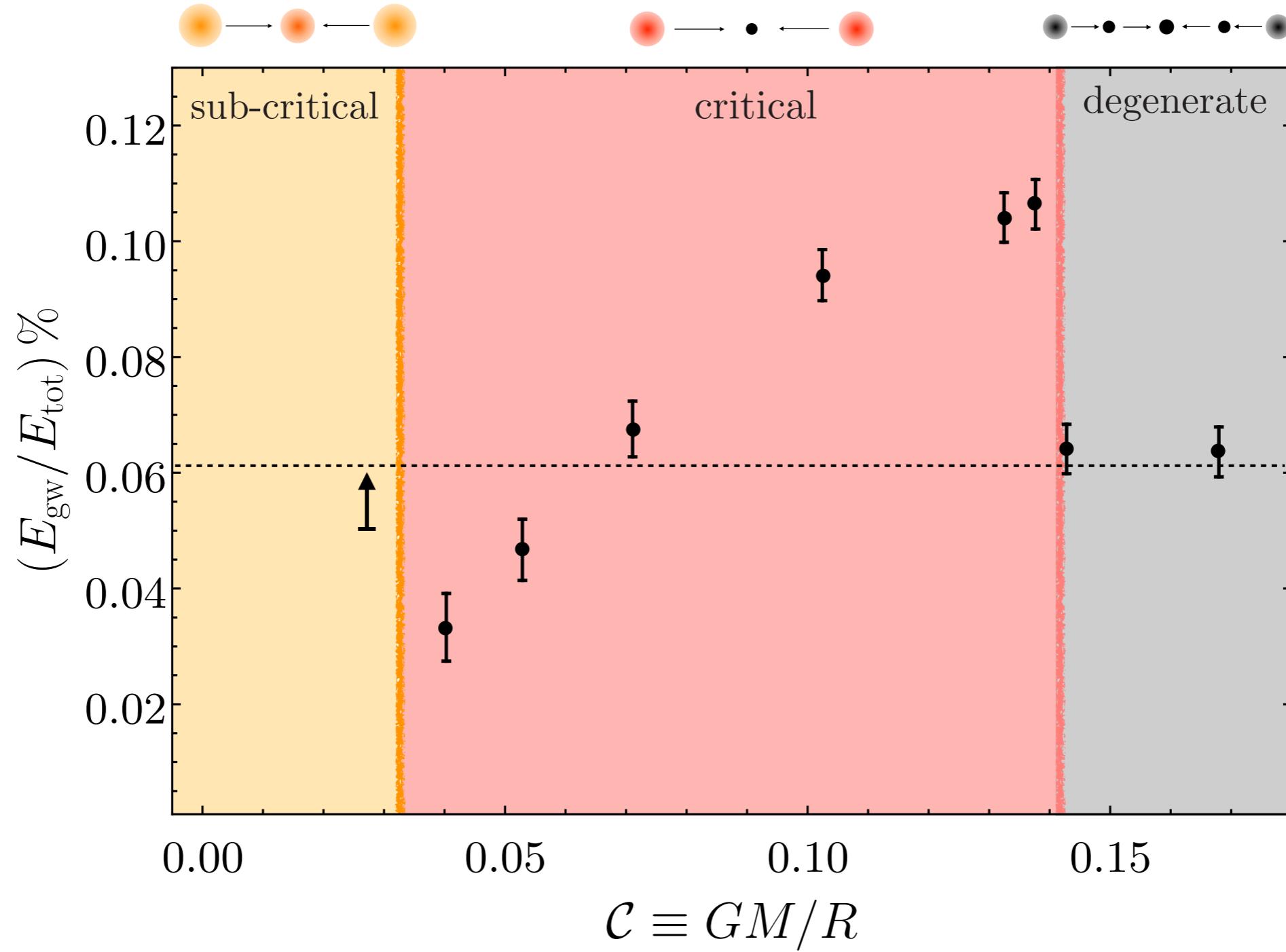
ultra-compact soliton



gravitational waves from ultra-compact soliton collisions



more energy in g-waves than corresponding mass BHs

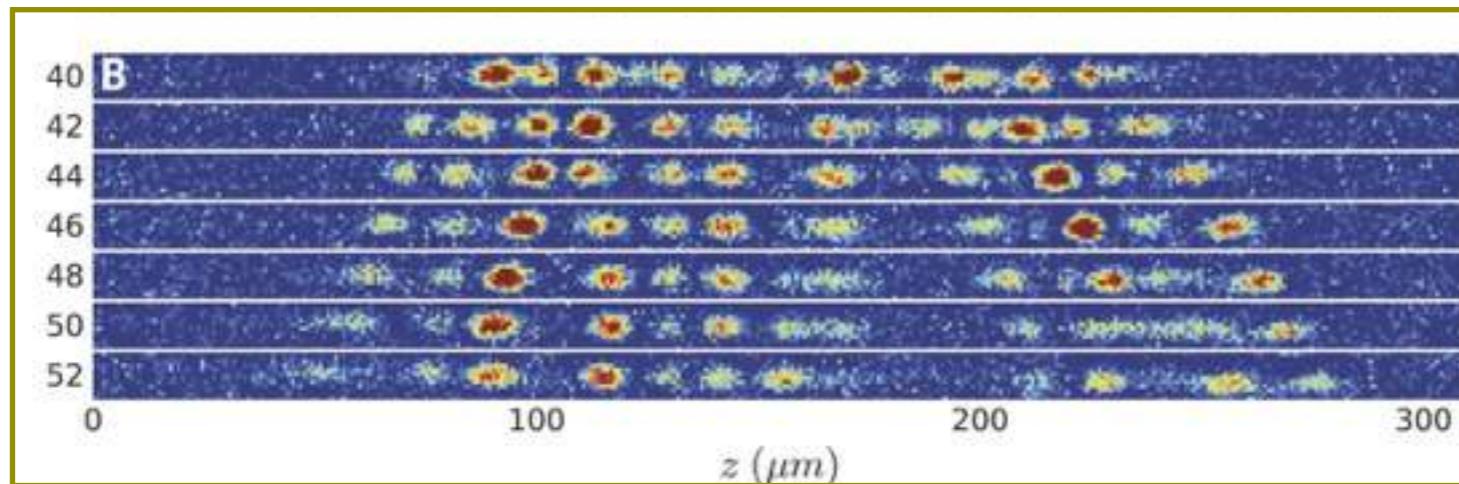


important caveats

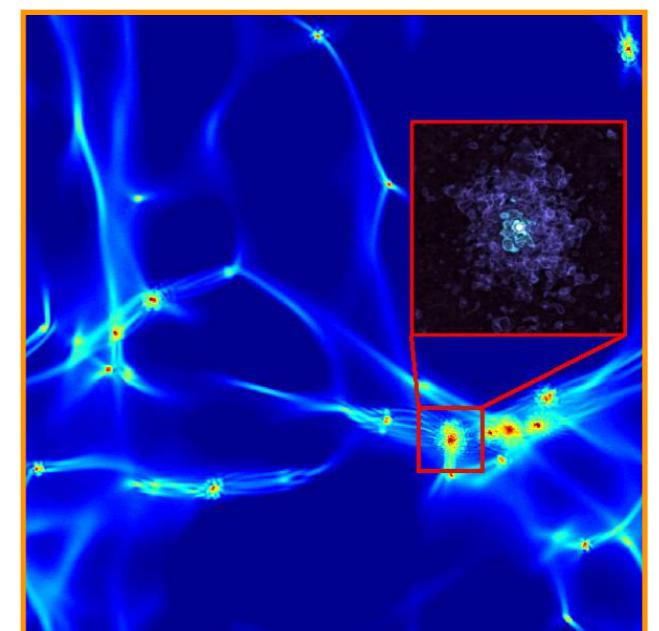
- solitons are in phase (out of phase ones can bounce)
- how likely are these to form ?
- head-on collisions: inspirals might change the answers (working on it — Helfer et. al)

things I did not discuss

- axions in the early universe [there are differences with case discussed]
 - axitons — axion mini clusters ...
- dynamics in late-time ultra-light axions
 - solitons at centers of halos/galaxies
 - small-scale structure of CDM (include baryons — in progress)
 - dynamical friction
- solitons in Bose-Einstein condensates



Nguyen, Luo & Hulet (2017)



Mocz et. al ...

relativistic* ✓

self-interactions ✓

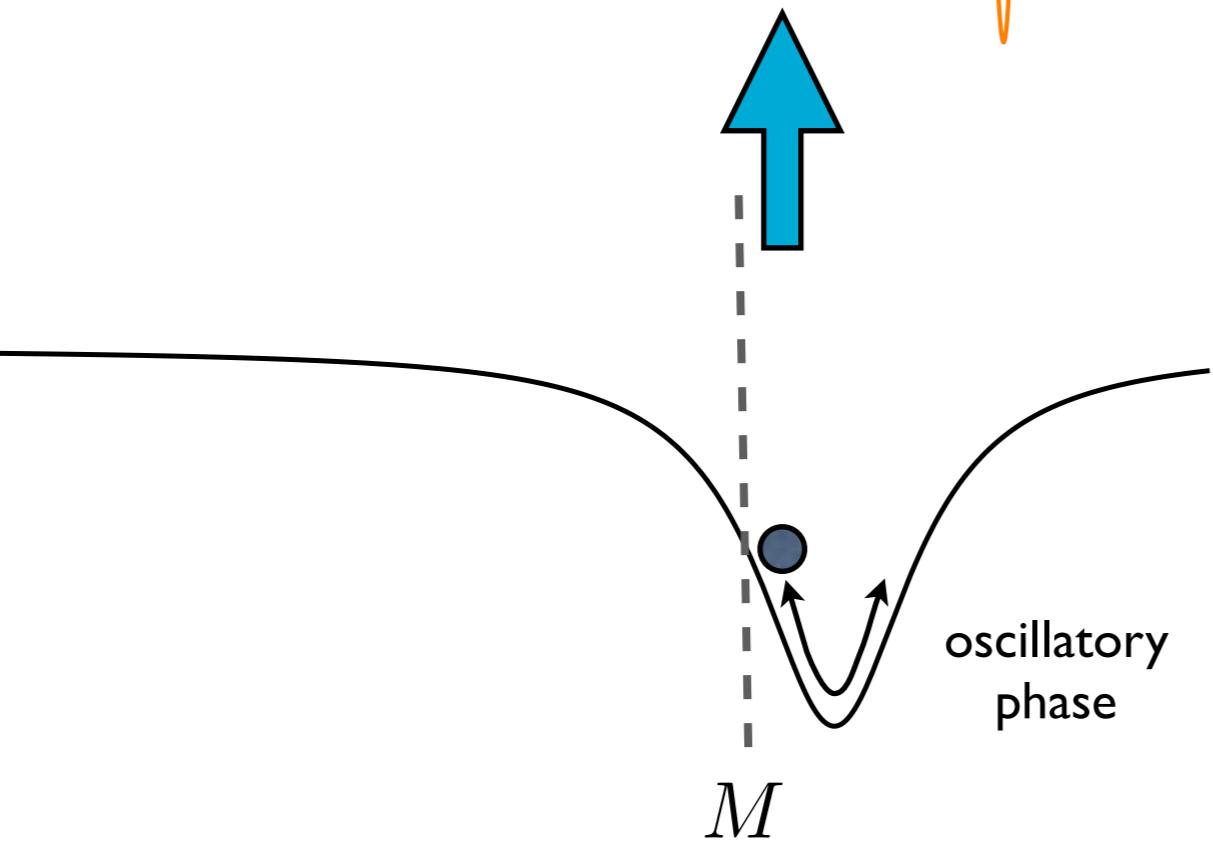
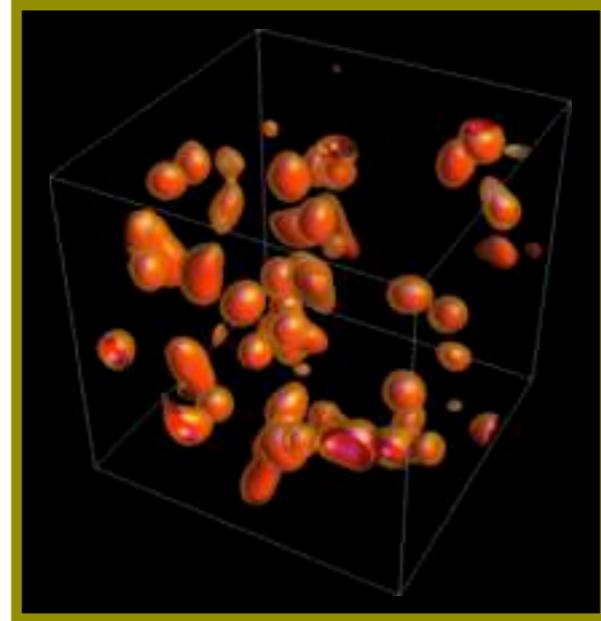
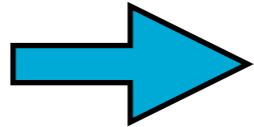
gravitational int. ✗

summary I

$$\frac{\mu_k}{H} \sim \frac{m_{\text{pl}}}{M} \gg 1$$

$$\delta\varphi_k(t) \propto e^{\mu_k t}$$

resonant growth



$$|\Phi|_{\text{sol}} \lesssim 10 \times \left(\frac{M}{m_{\text{pl}}} \right)^2$$

$$\Phi \lesssim \text{few} \times 10^{-3}$$

not easy to form
black holes
from individual
solitons*

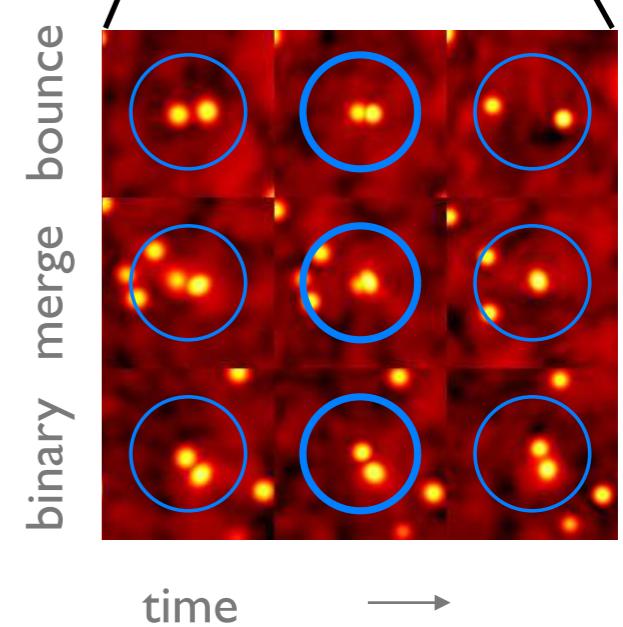
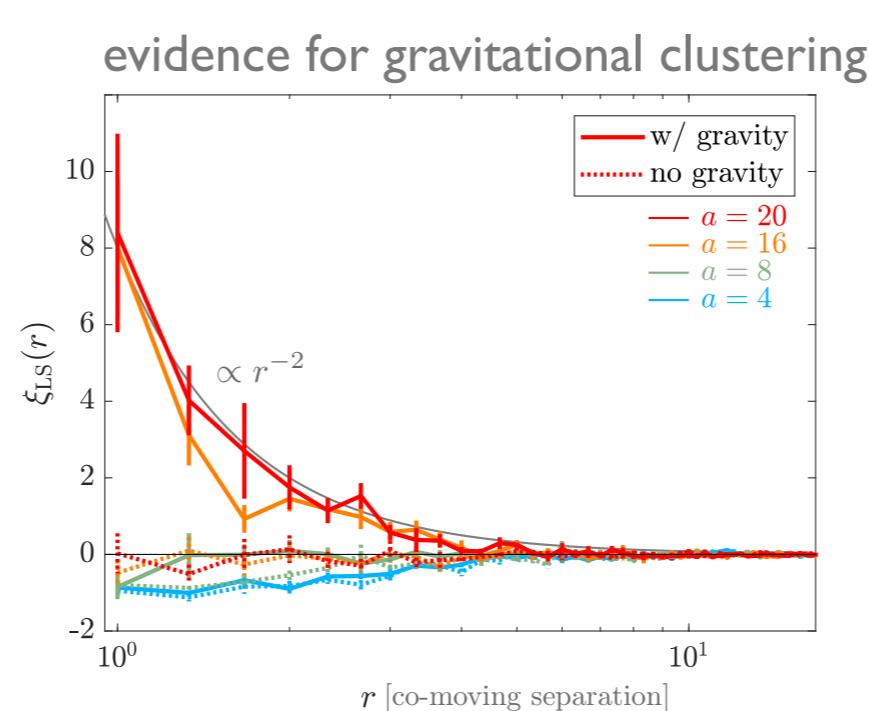
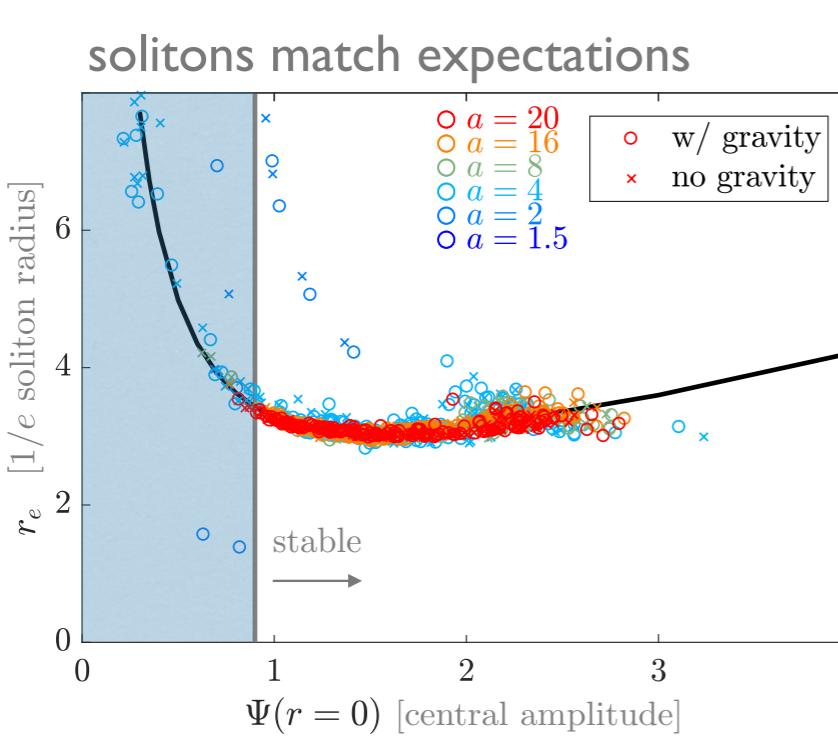
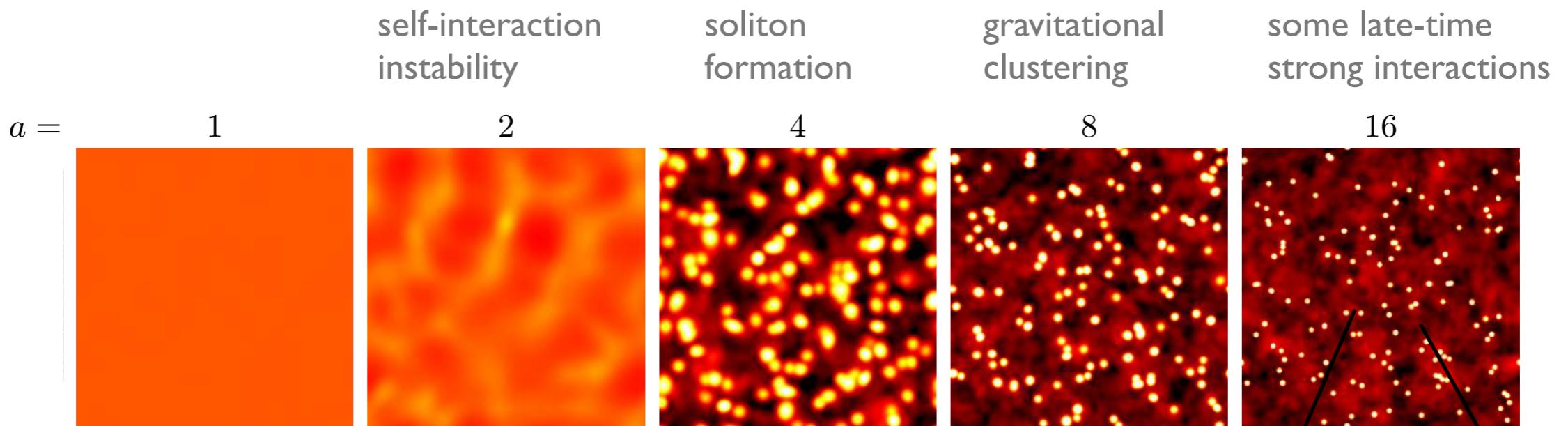
$$\Omega_{\text{GW},0} h_{100}^2 \sim 10^{-6} \left(\frac{M}{m_{\text{pl}}} \right)^2 \lesssim \mathcal{O}[10^{-9}]$$

relativistic ✗

self-interactions ✓

gravitational int. ✓

summary 2



Numerical GR ✓

self-interactions ✗

formation ✗

summary 3

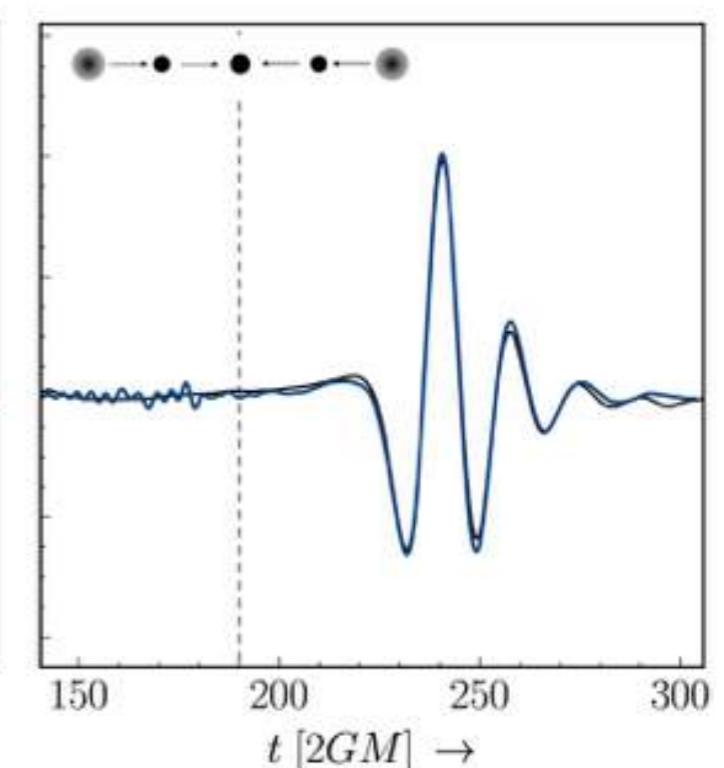
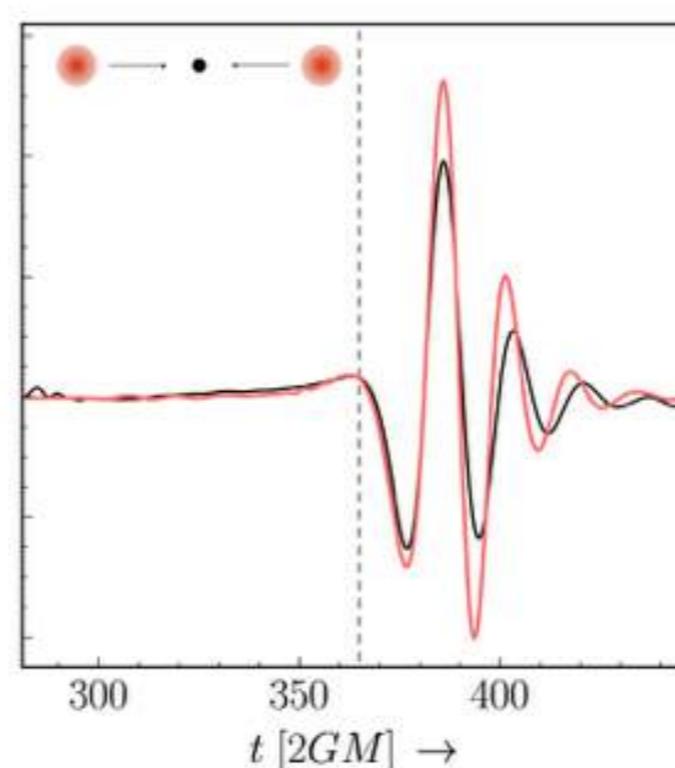
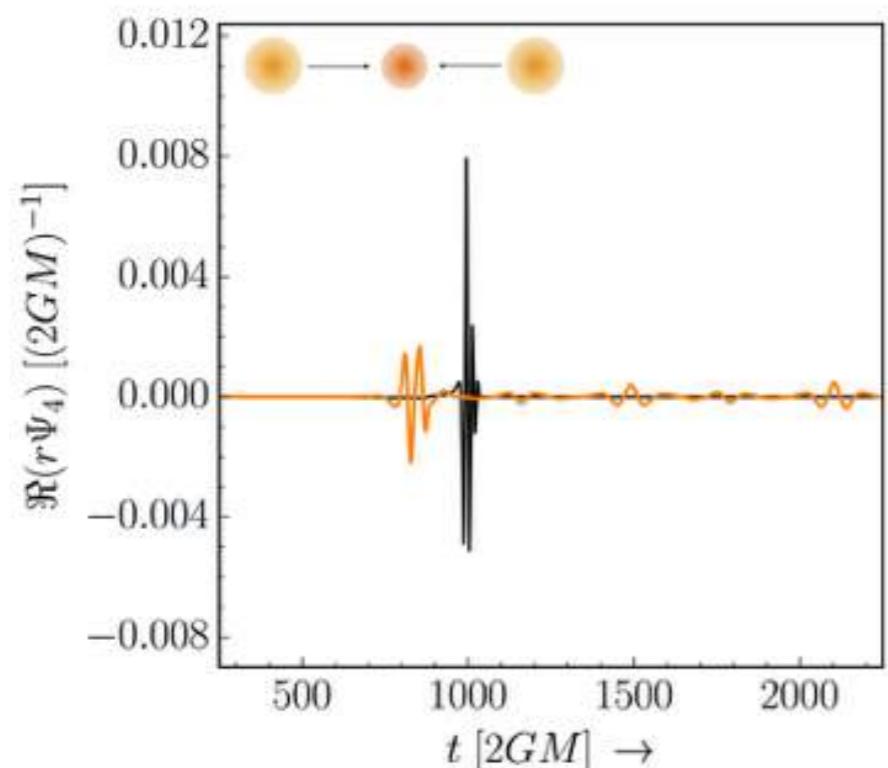
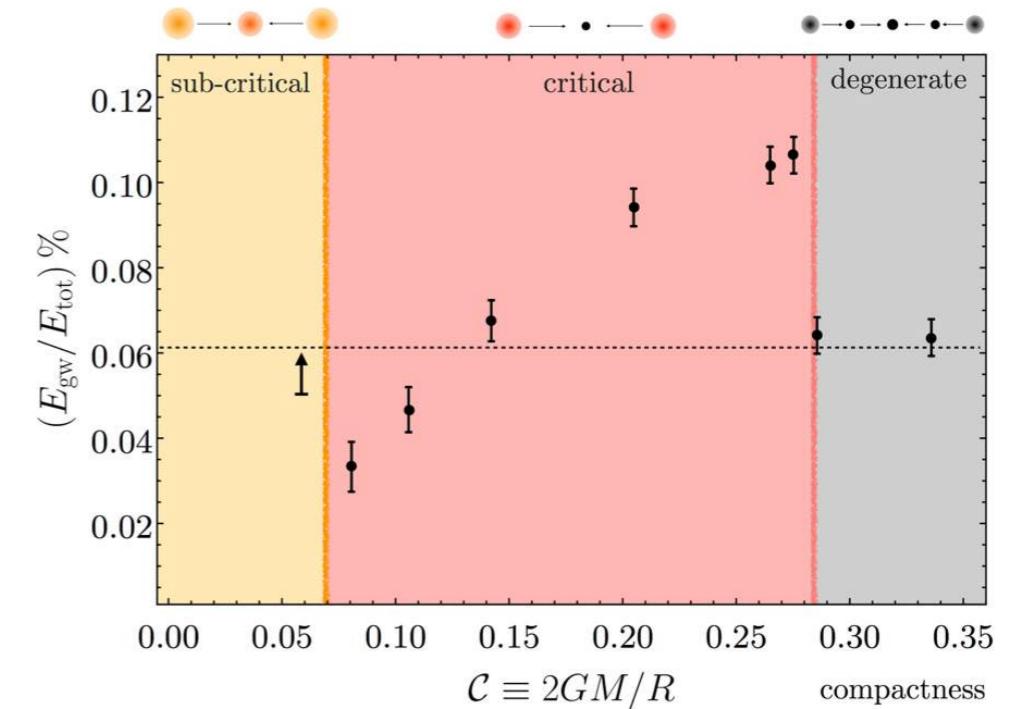
sub-critical



critical



degenerate



thanks

