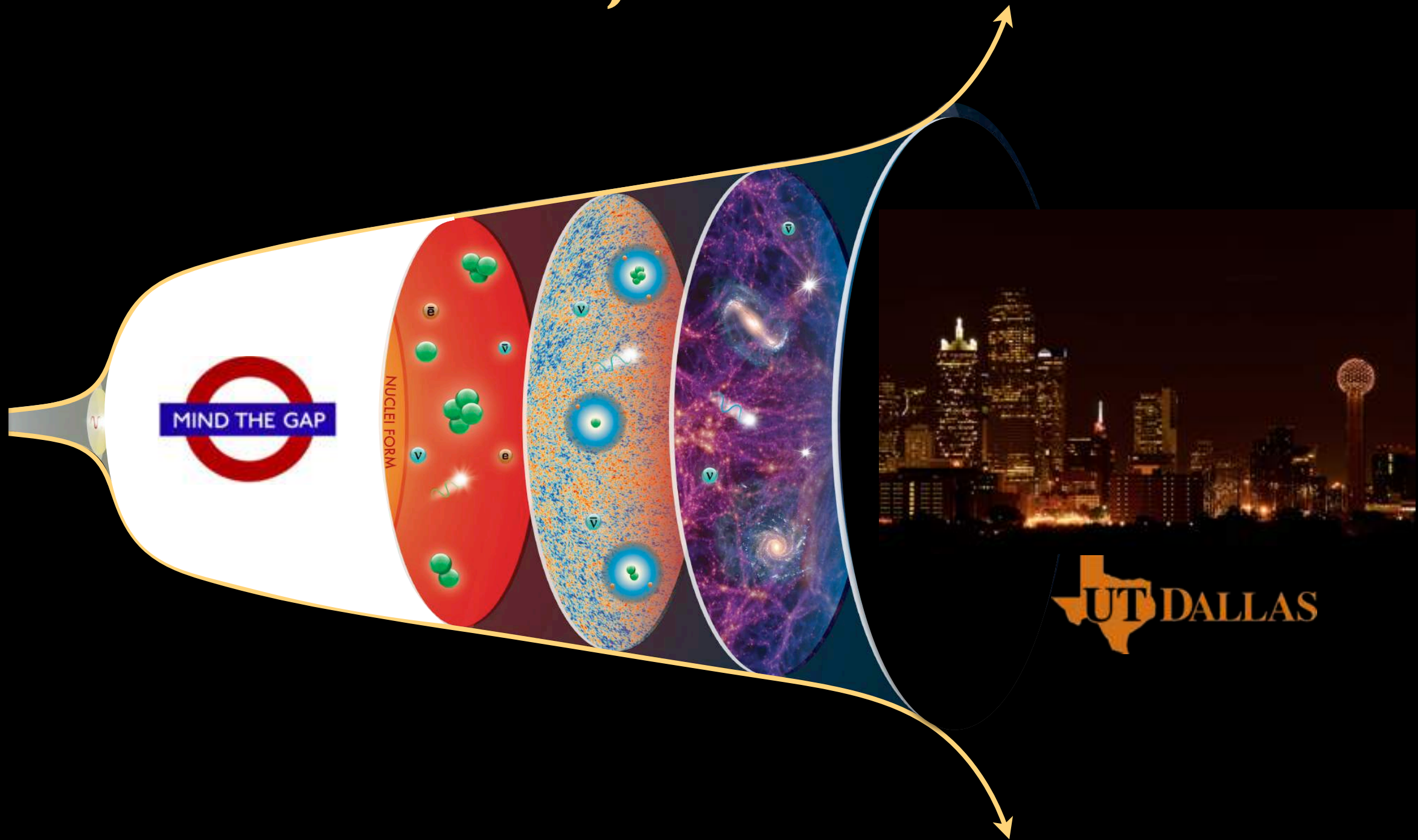
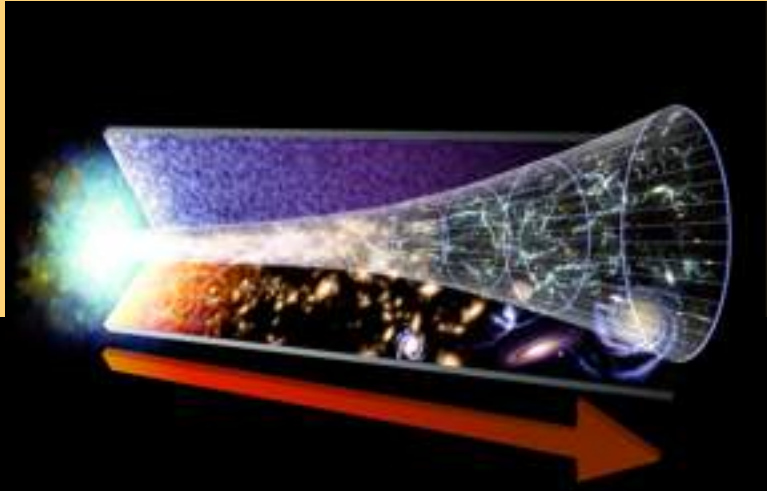


Inflation Ends, What's Next ?



Mustafa A. Amin





Outline

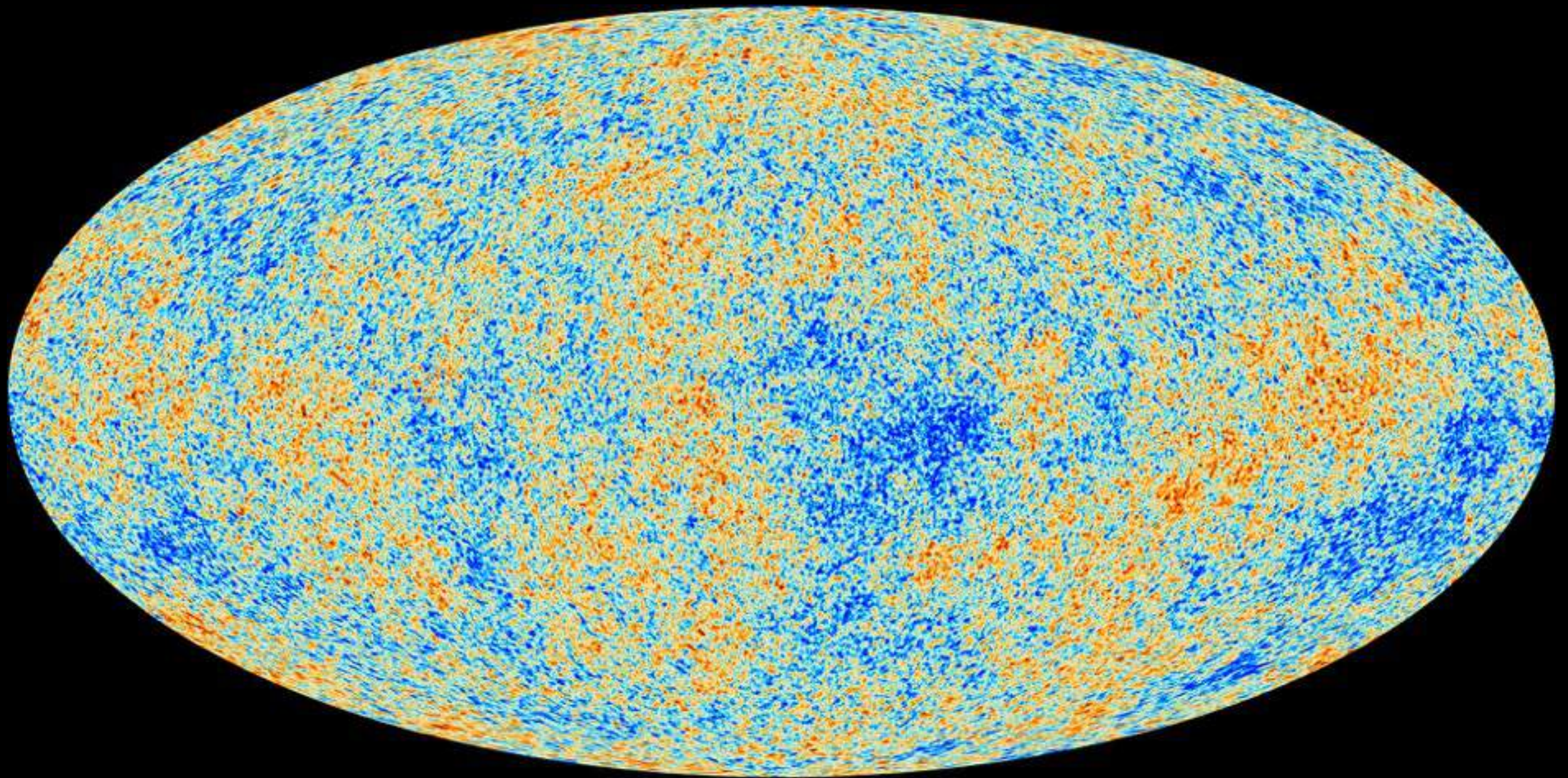
◎ inflation

◎ **cosmic history gap - aftermath of inflation**

- non-perturbative, nonlinear dynamics
- obs. consequences

◎ what's next ?

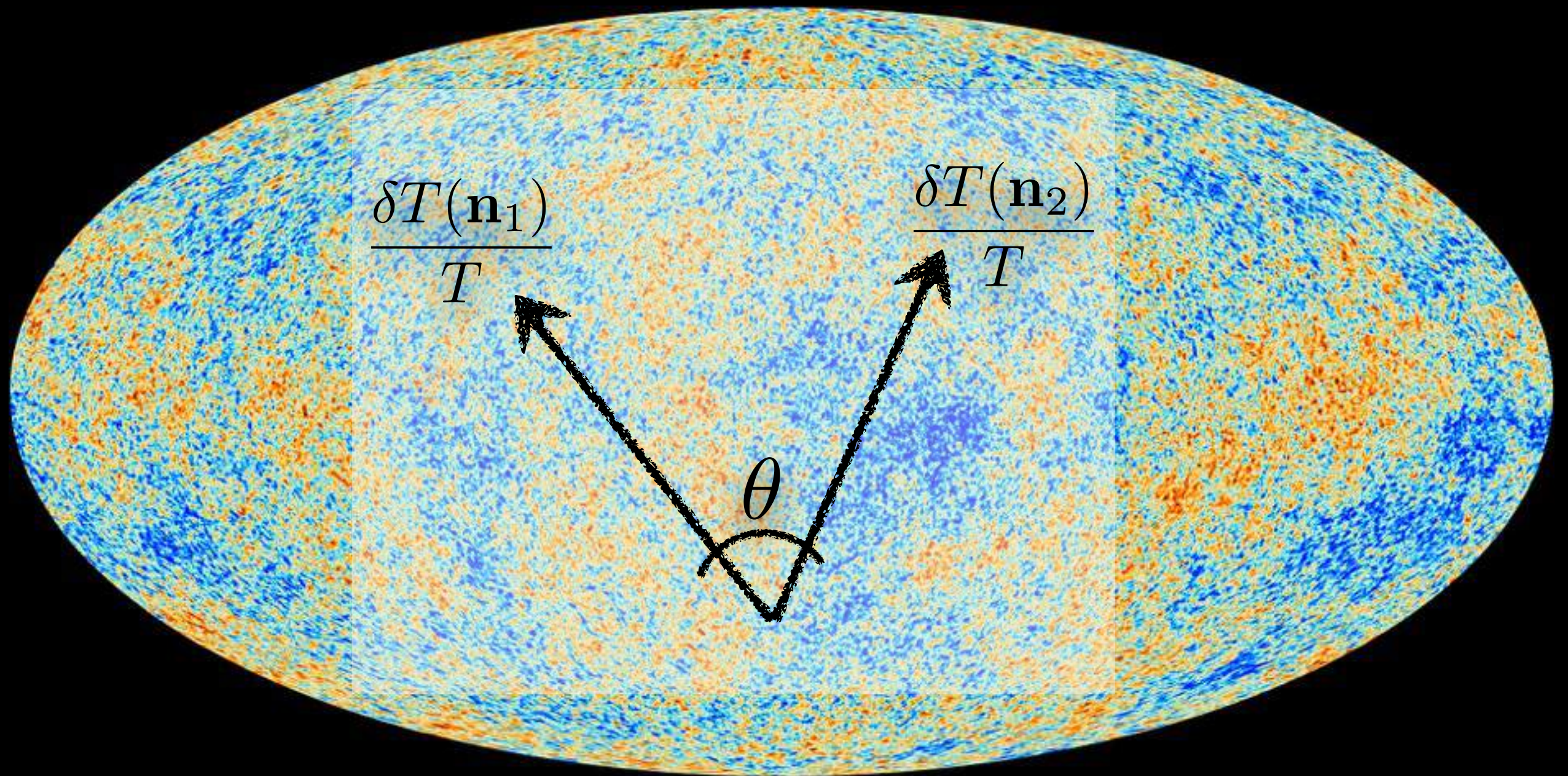
Cosmic Microwave Background (CMB) temperature anisotropies



$$\delta T/T \sim 10^{-5}$$

a conundrum ... and a solution

large angle correlations?



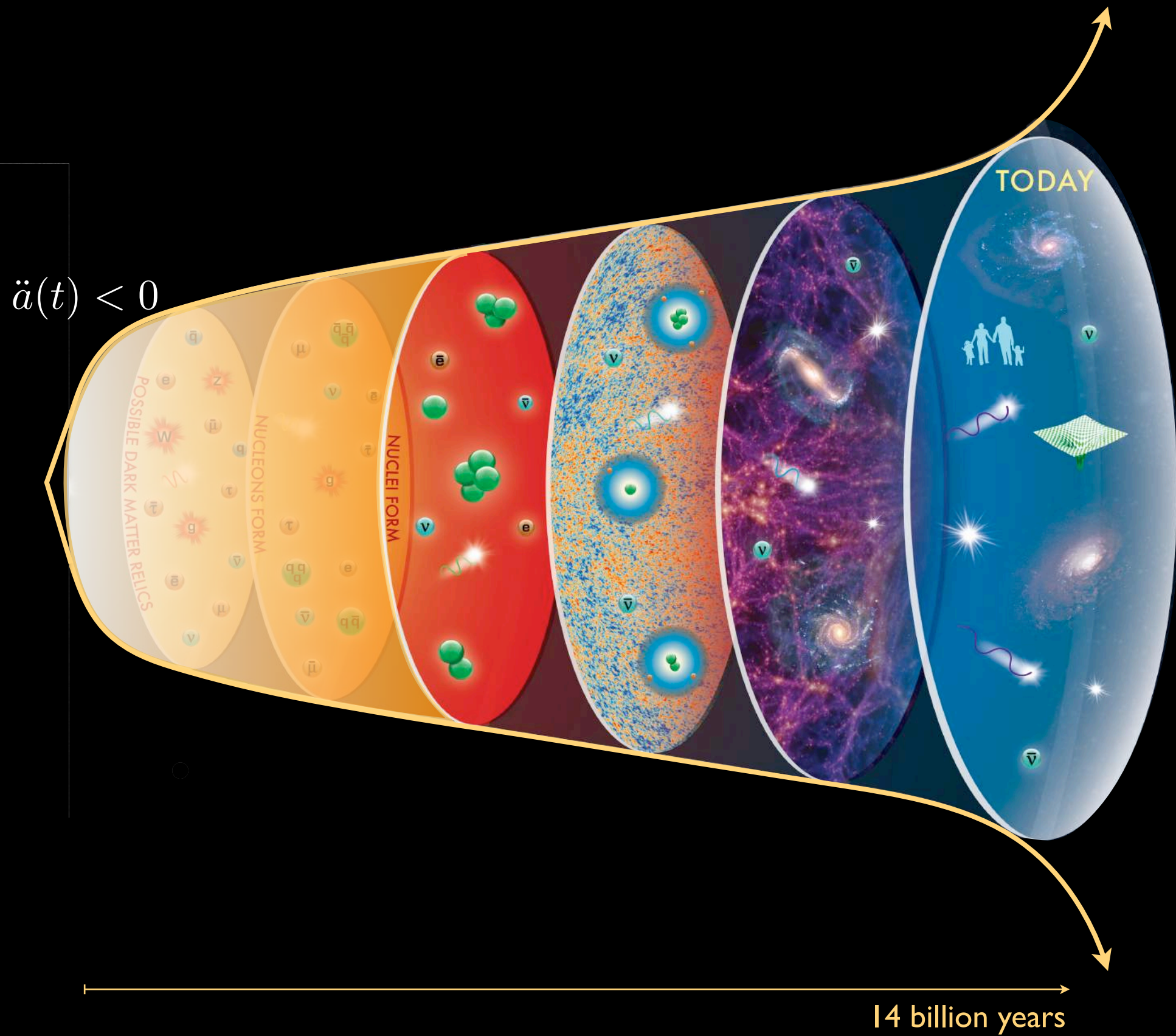
$$C(\theta) = \left\langle \frac{\delta T(\mathbf{n}_1)}{T} \frac{\delta T(\mathbf{n}_2)}{T} \right\rangle$$

standard big bang cosmology

size of causally connected patch



standard cosmology



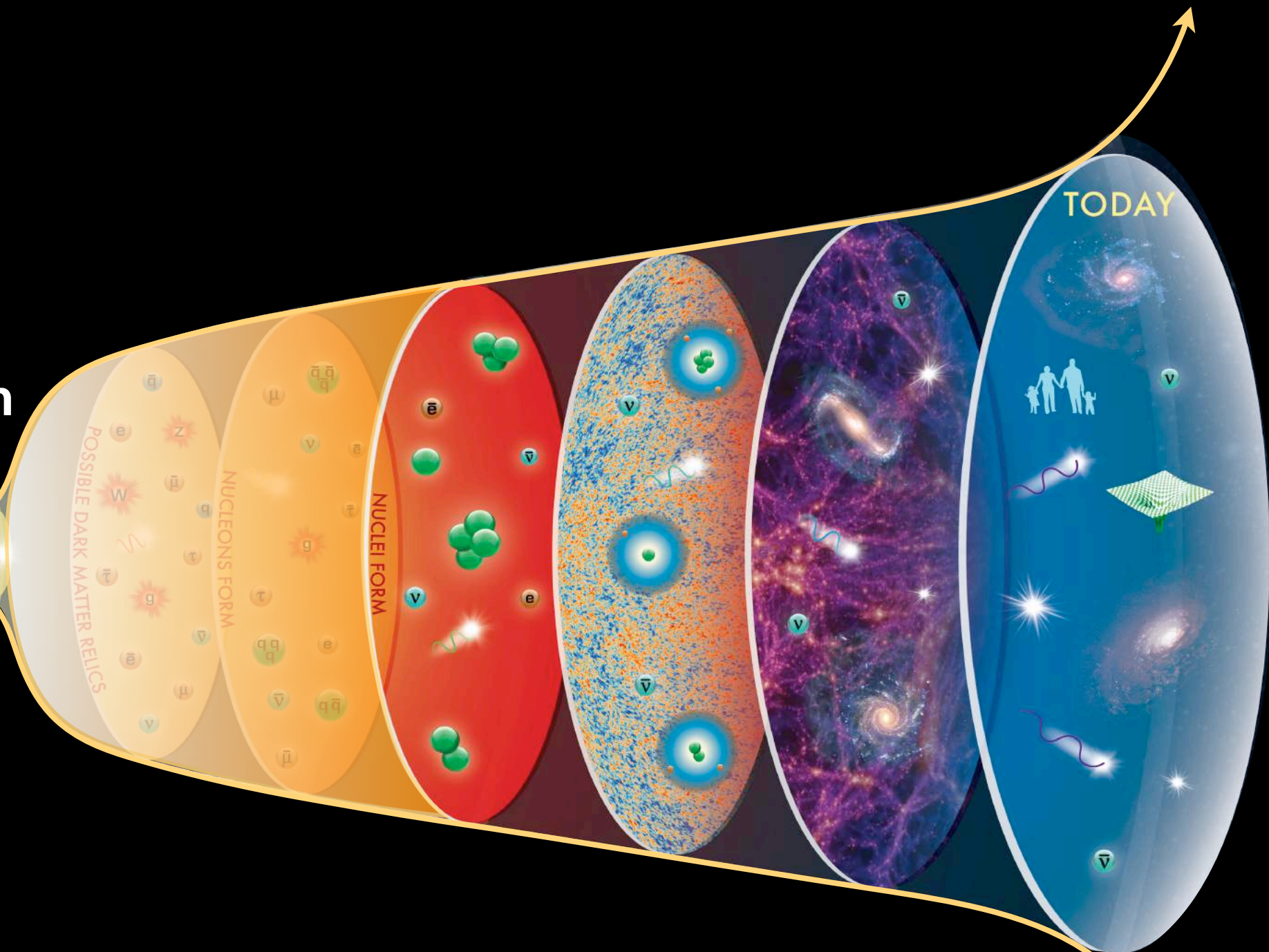
inflationary cosmology

inflation

$$\ddot{a}(t) > 0$$

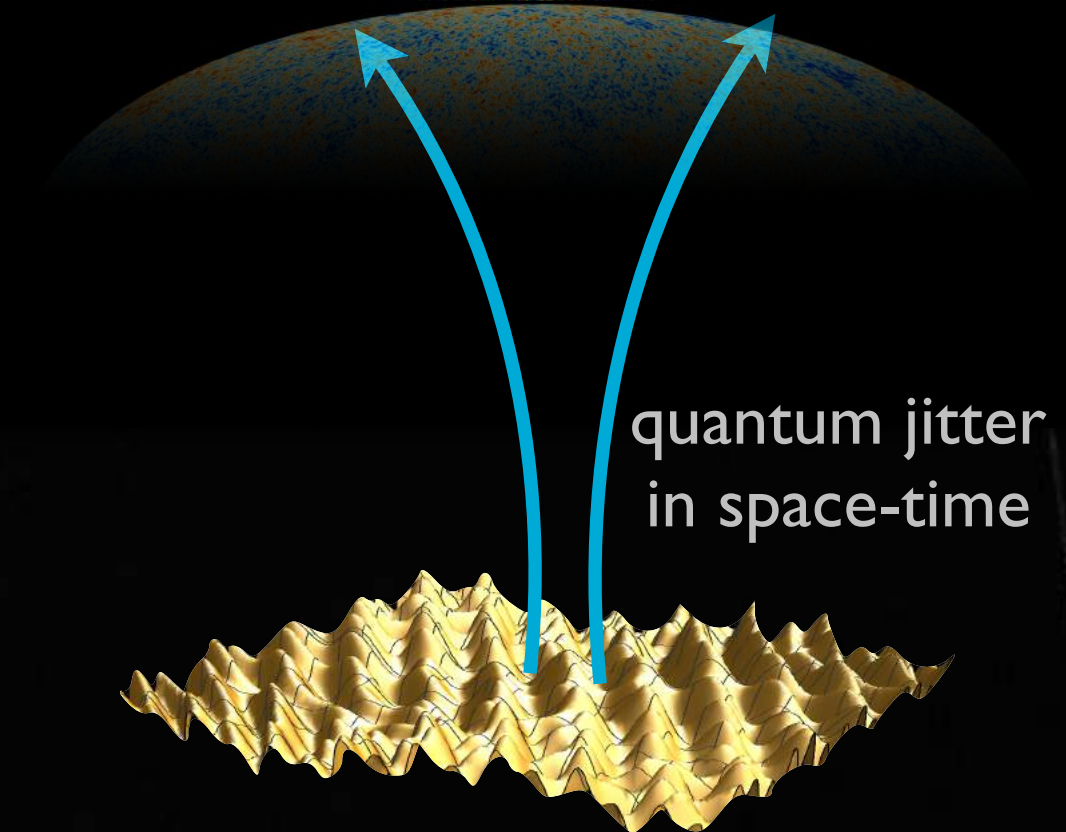
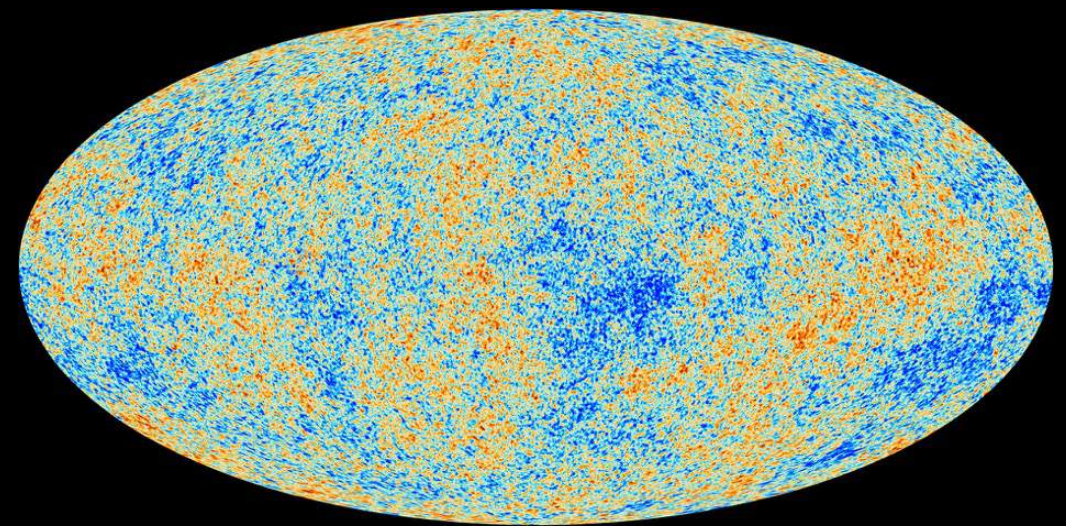
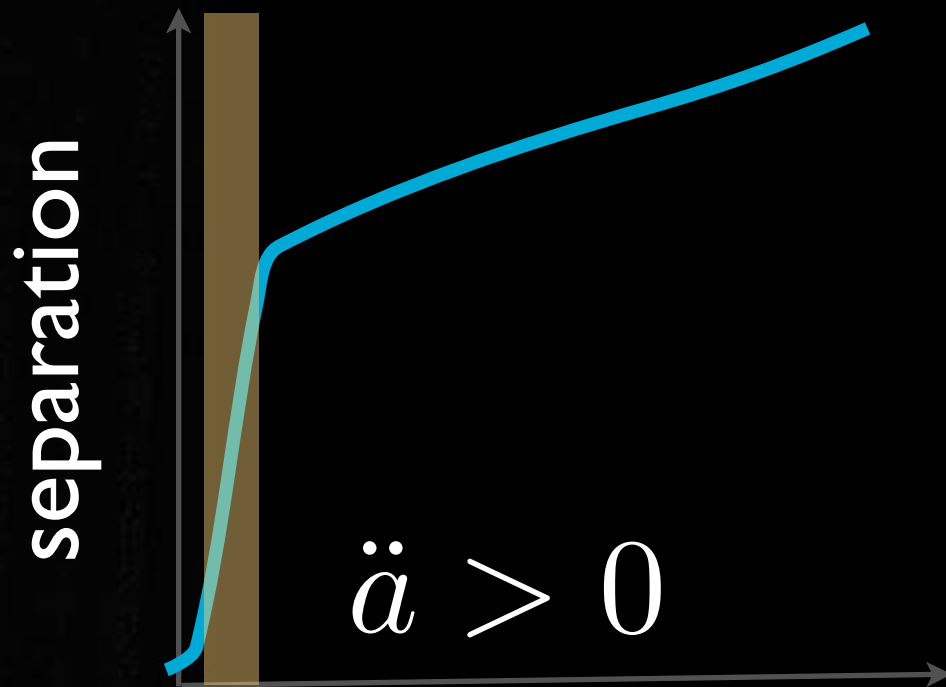
?

??

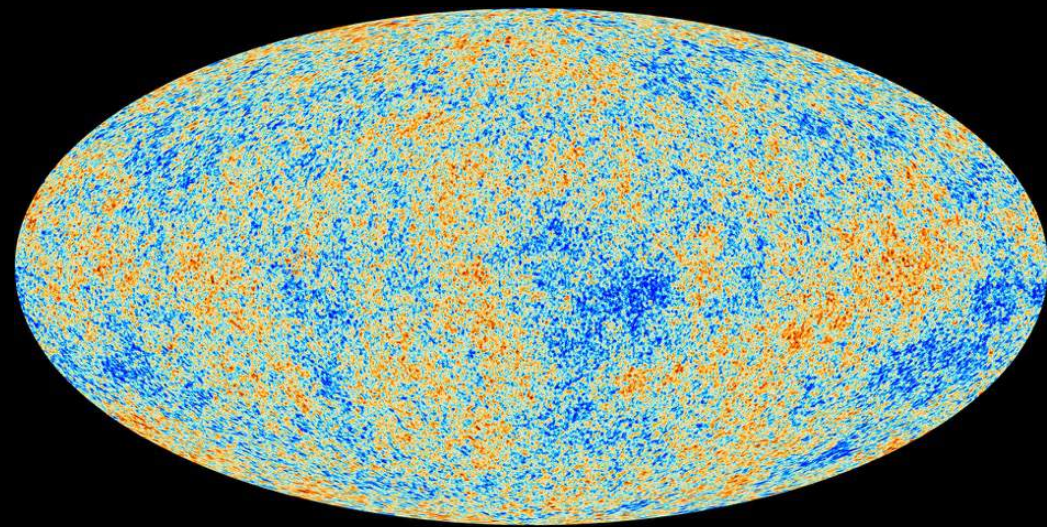


14 billion years

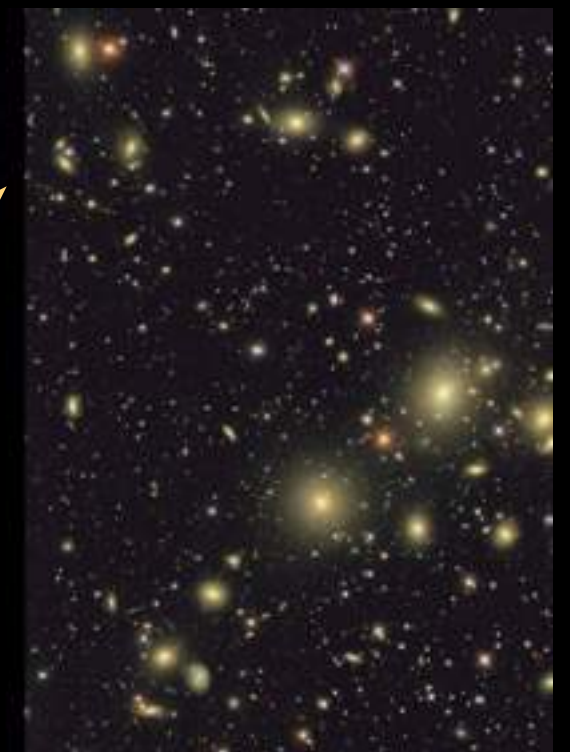
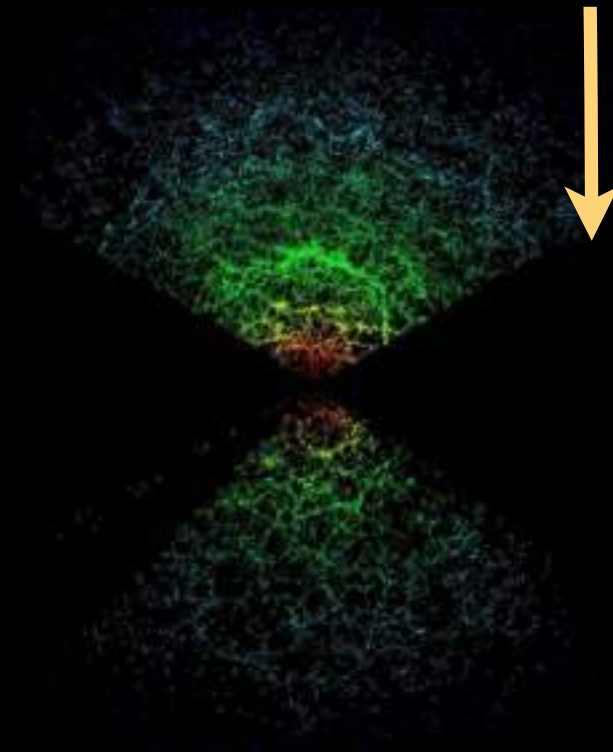
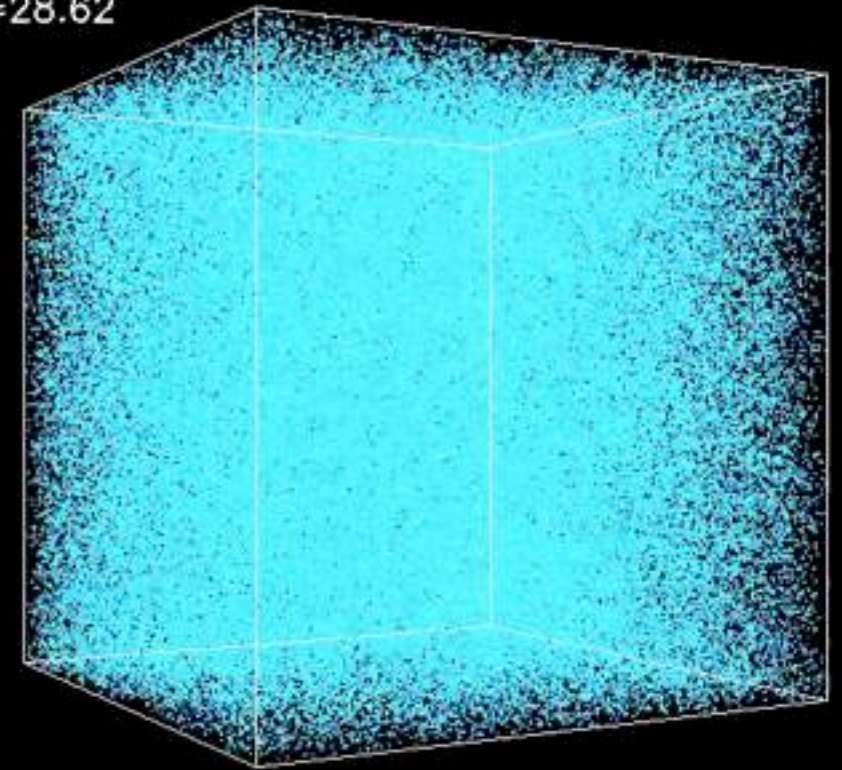
large angle correlations



inflationary cosmology: a calculable framework of initial perturbations*

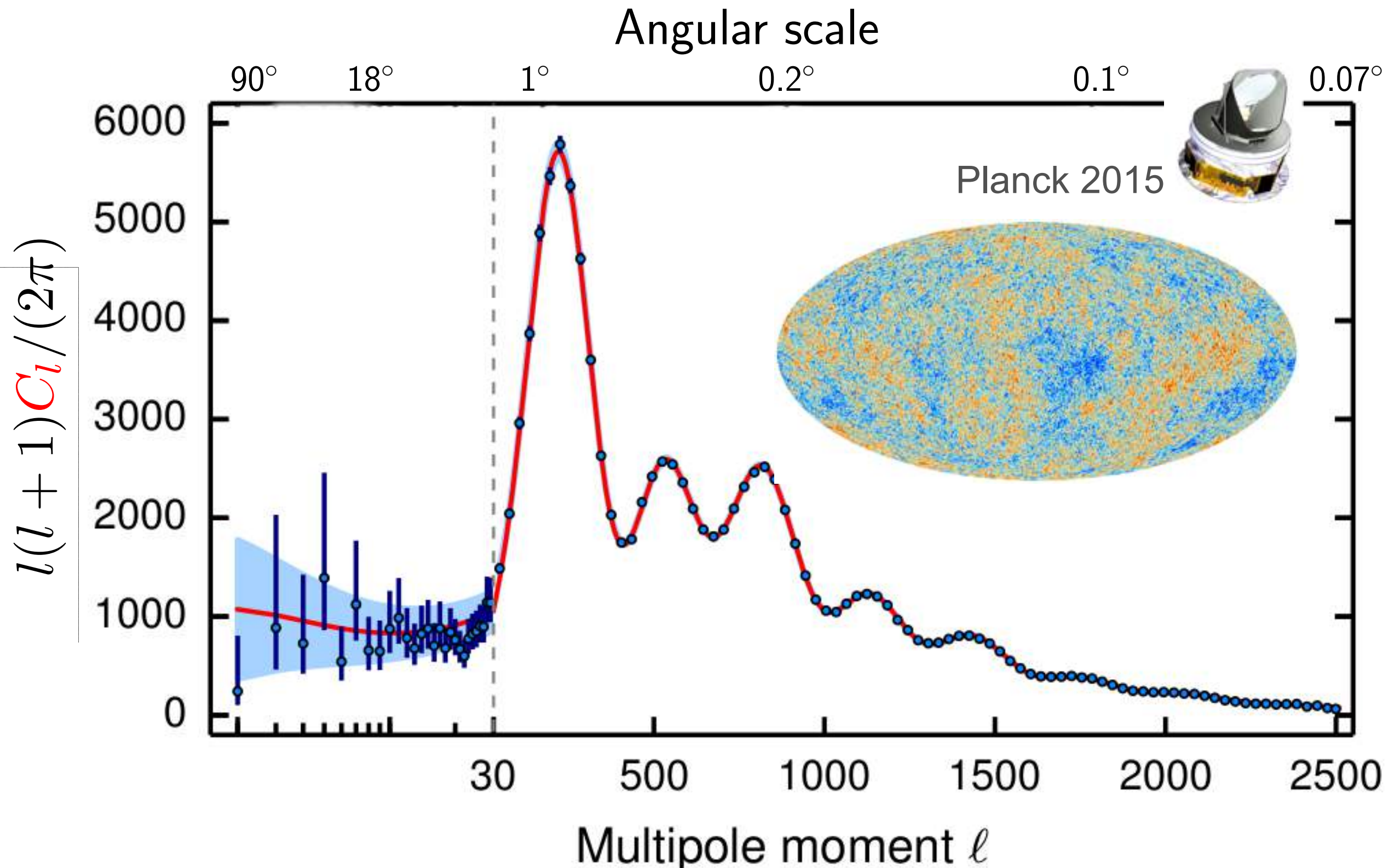


$z=28.62$

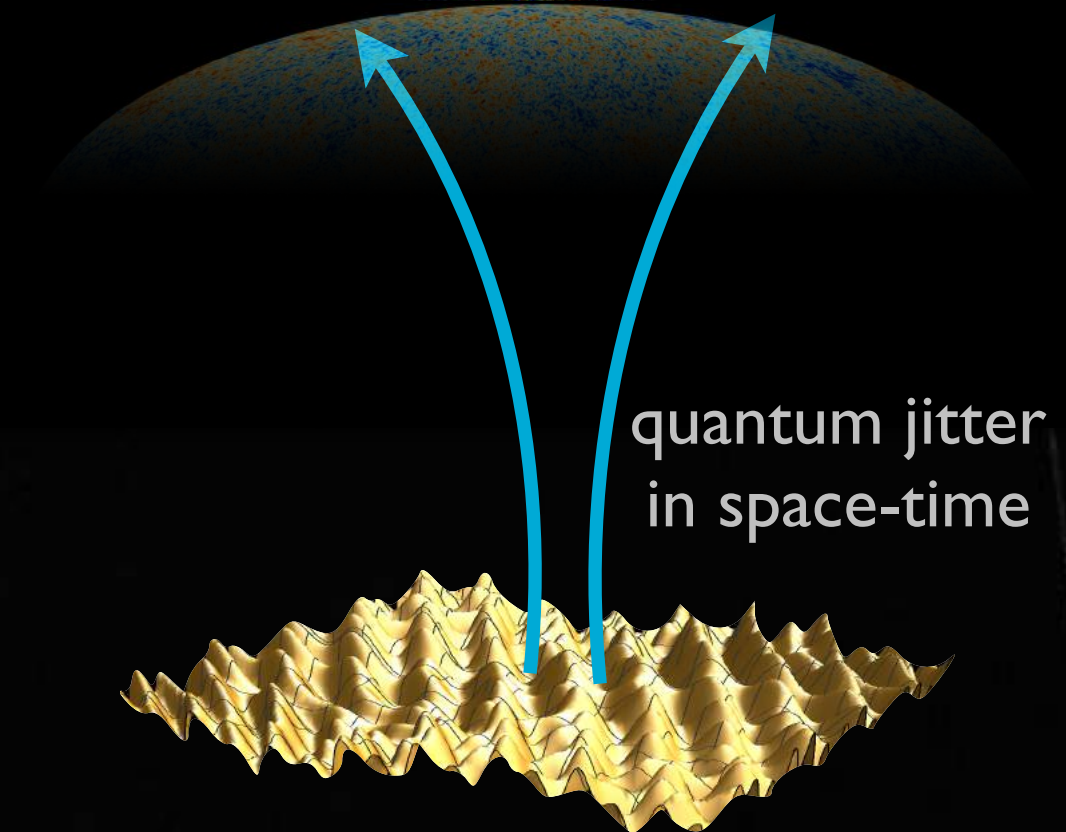
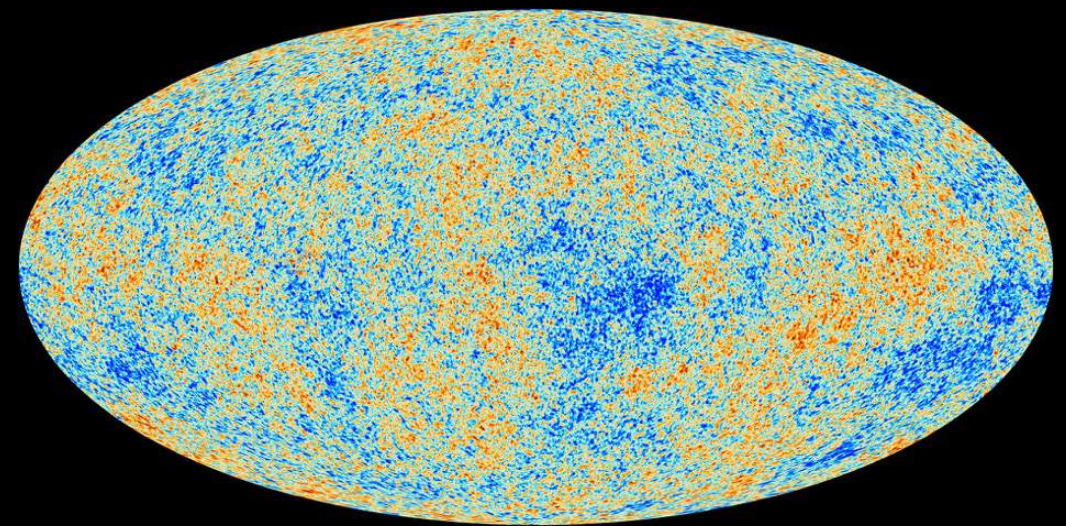
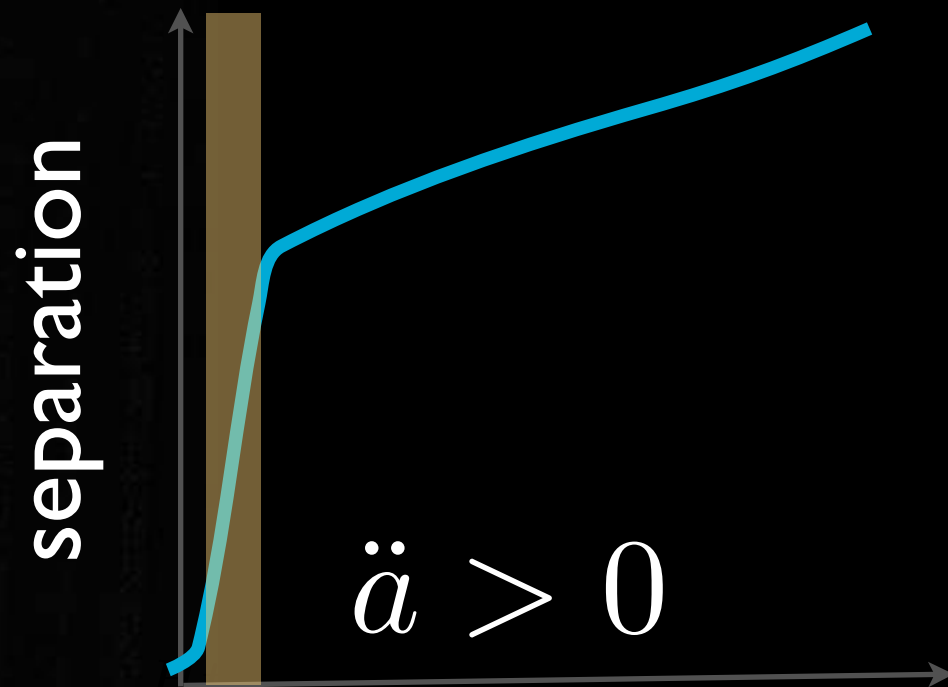


US

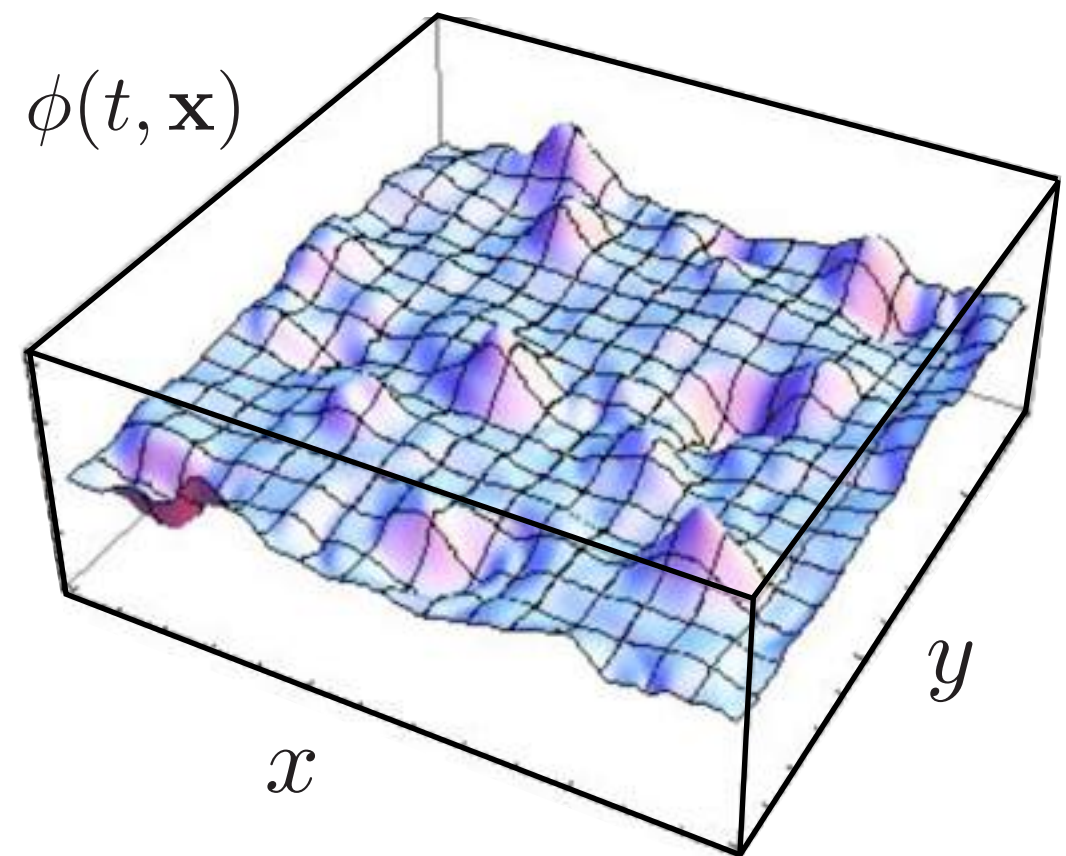
inflation consistent with observations



what drives inflation ?



simplest: single scalar field driven inflation

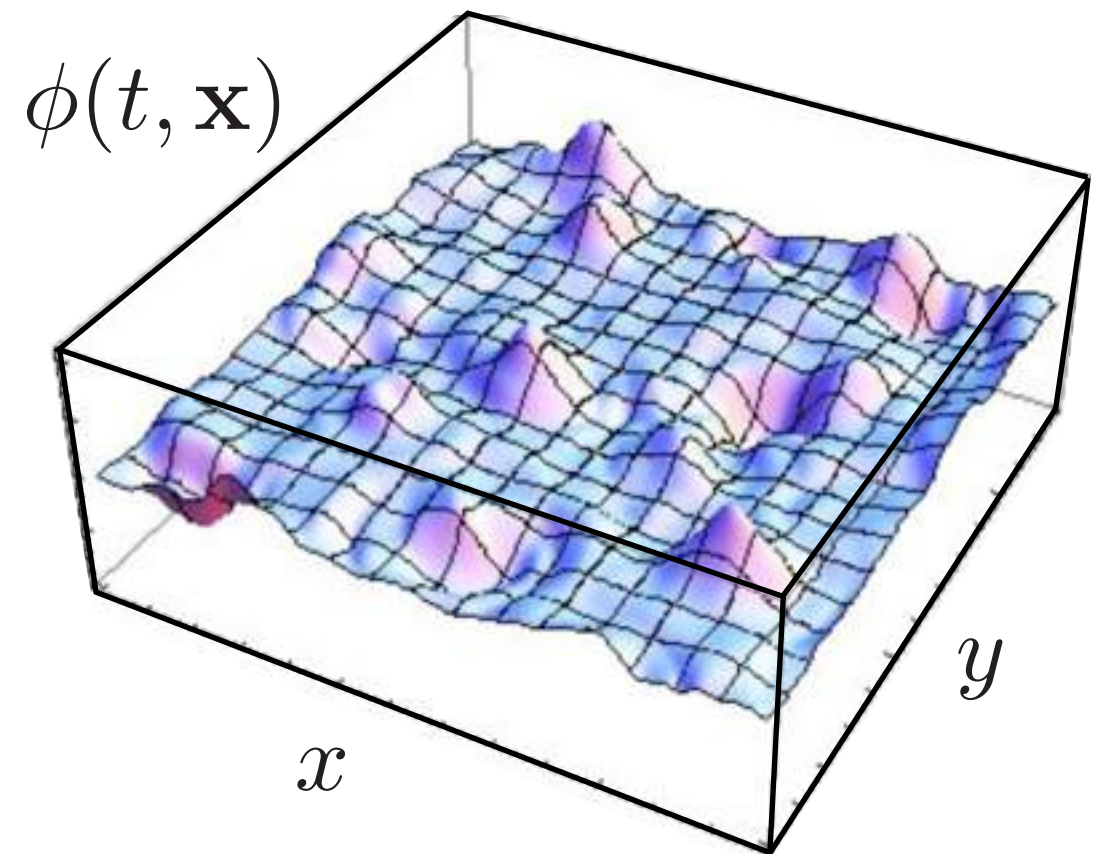


simplest: single scalar field driven inflation

Lagrangian: $\mathcal{L} = \frac{1}{2}(\partial\phi)^2 - V(\phi)$



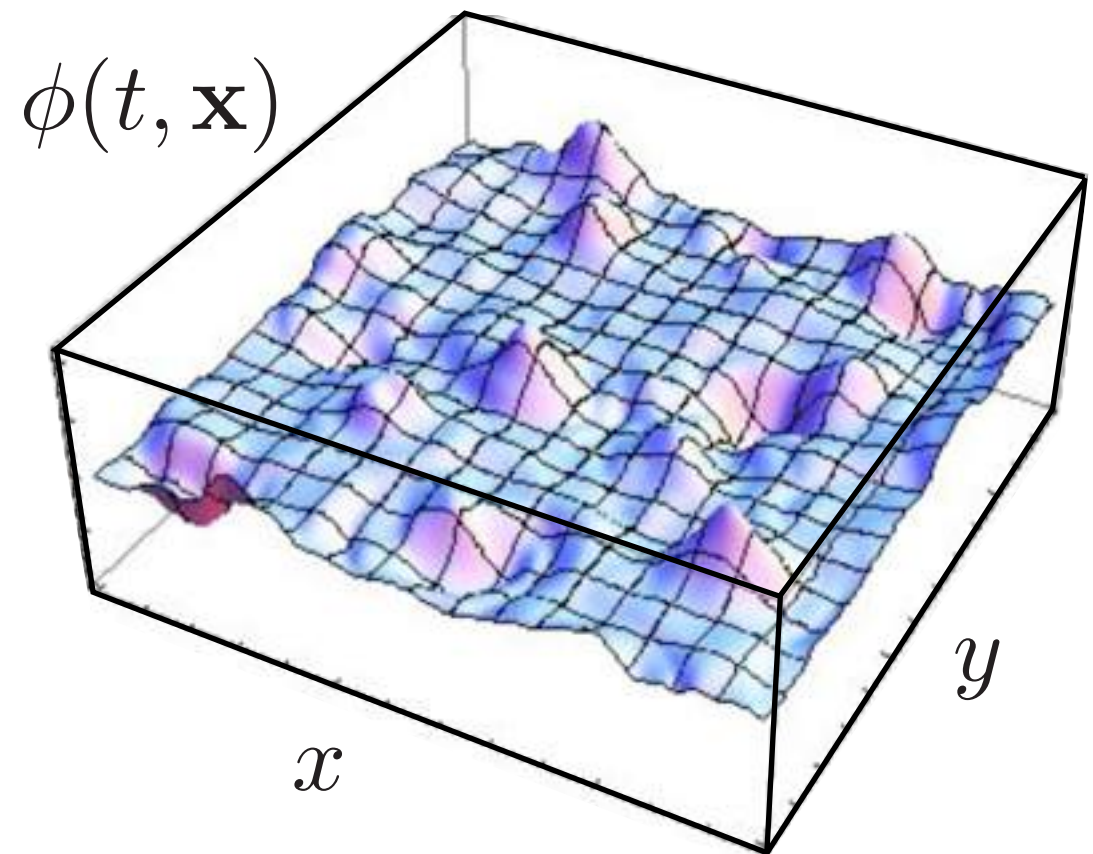
kinetic & gradient energy density



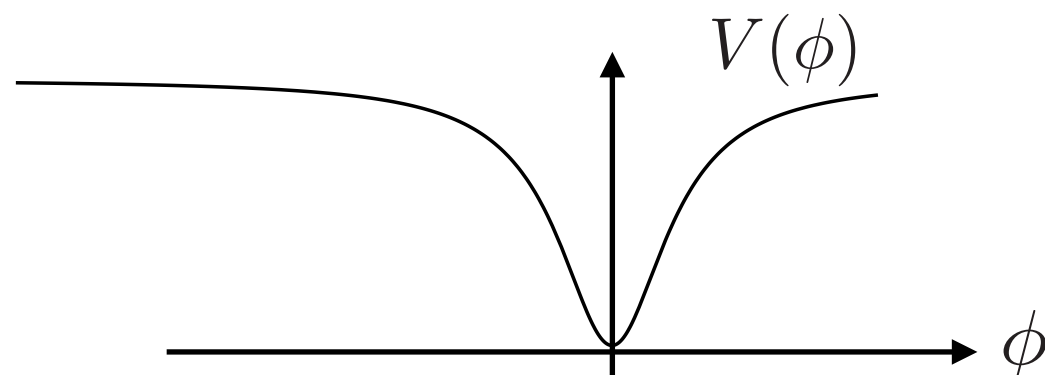
simplest: single scalar field driven inflation

Lagrangian: $\mathcal{L} = \frac{1}{2}(\partial\phi)^2 - V(\phi)$

kinetic & gradient energy density



potential: determines the dynamics

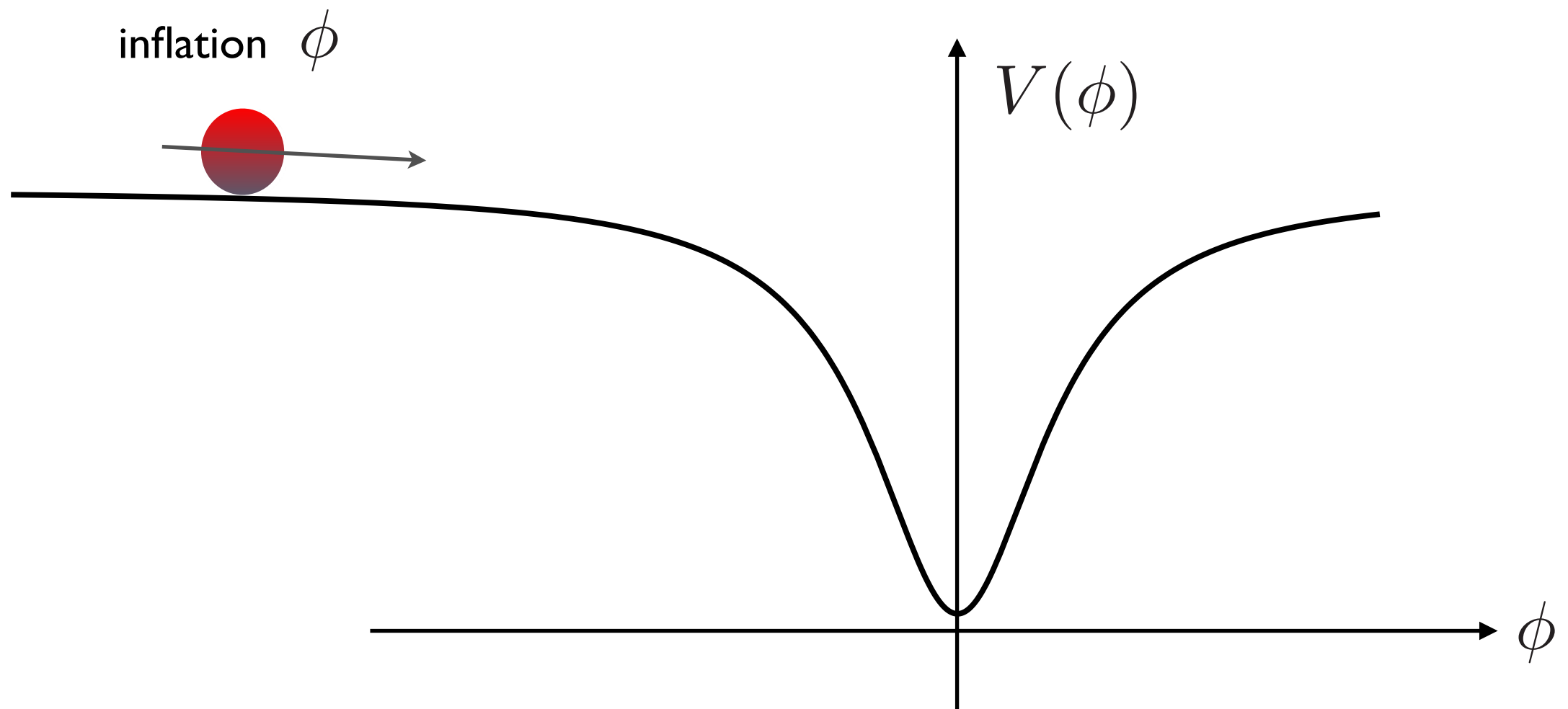
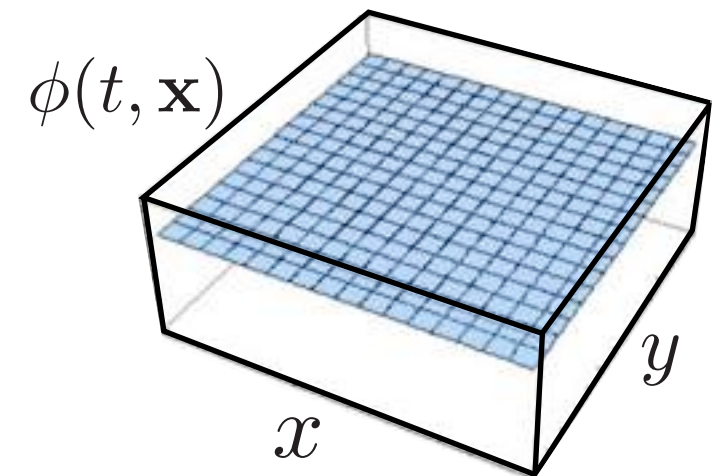


scalar field driven inflation



$$G_{\mu\nu} = 8\pi G T_{\mu\nu}$$

$$\Rightarrow \ddot{a}(t) \propto V(\phi) - \dot{\phi}^2$$



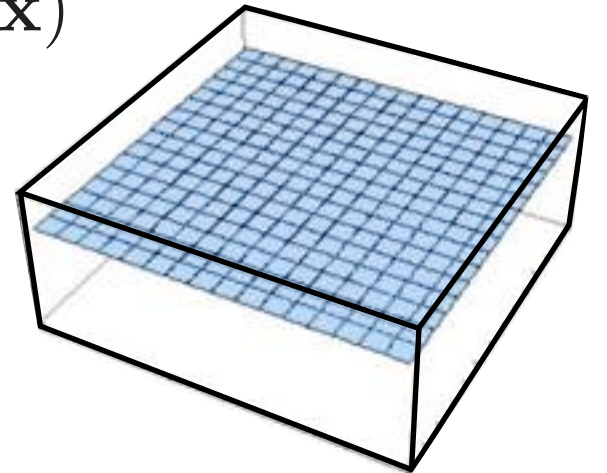
slowly “rolling” field = accelerated expansion



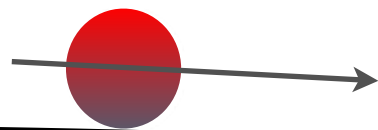
$$G_{\mu\nu} = 8\pi G T_{\mu\nu}$$

$$\Rightarrow \ddot{a} \propto V(\phi) - \dot{\phi}^2 > 0$$

$\phi(t, \mathbf{x})$



inflation ϕ

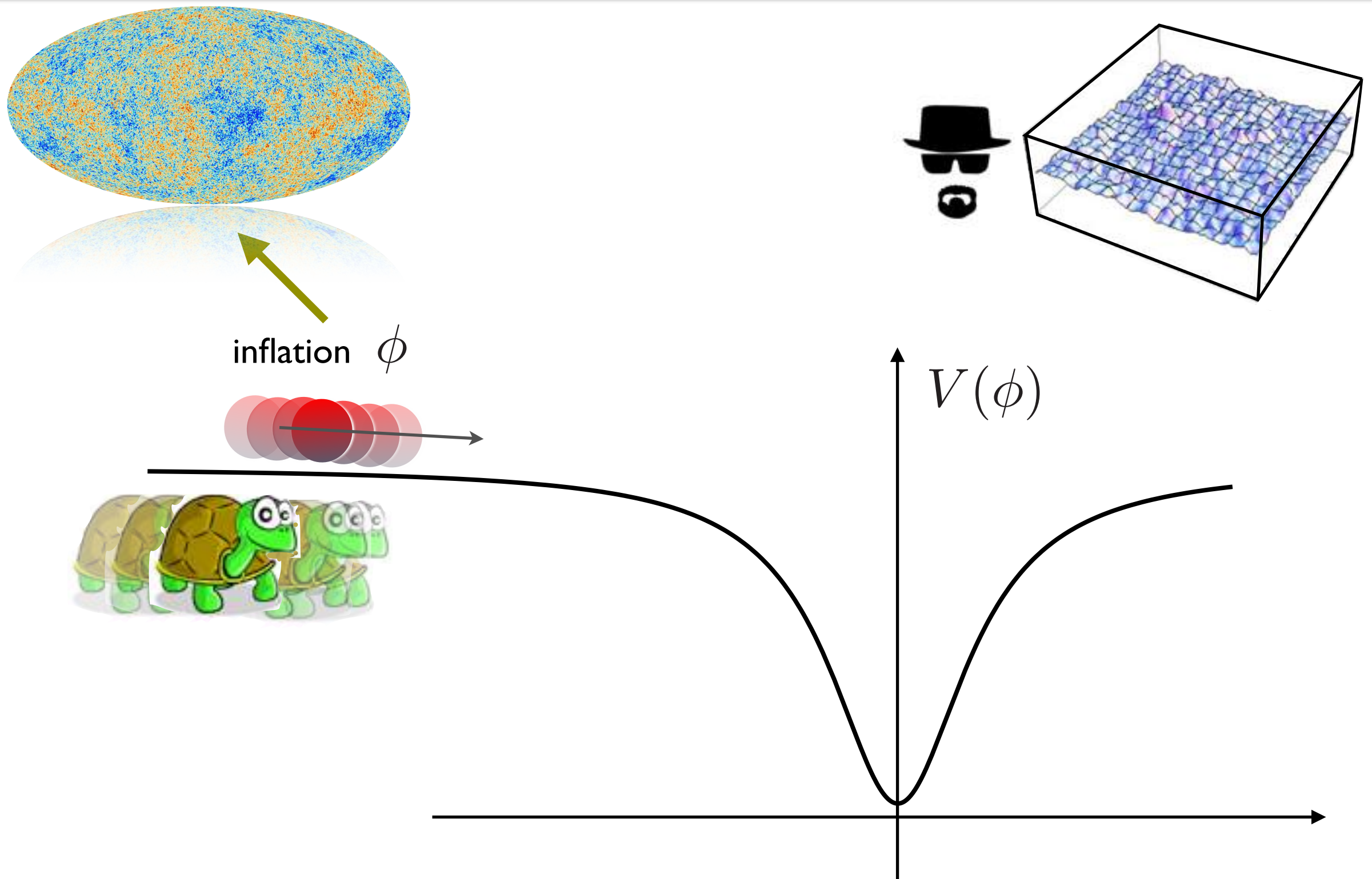


$$\dot{\phi}^2 \ll V(\phi)$$

$V(\phi)$

ϕ

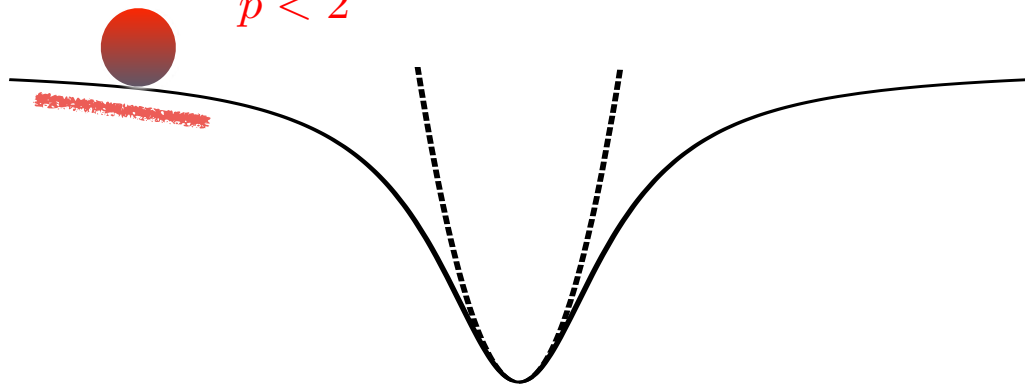
inflationary quantum perturbations



constraints from observations

$$V(\phi) \propto \phi^p$$

$$p < 2$$

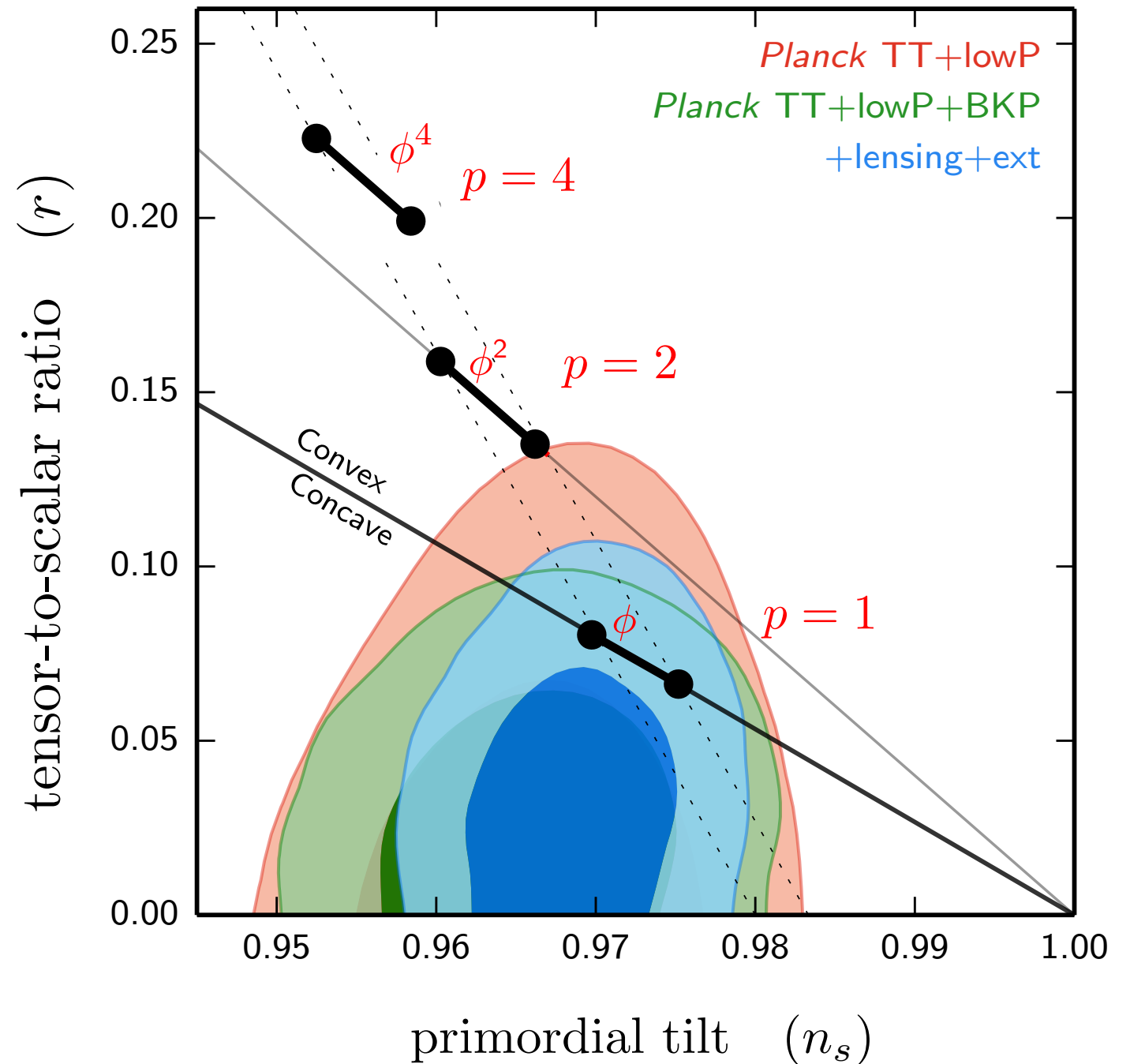


for example:

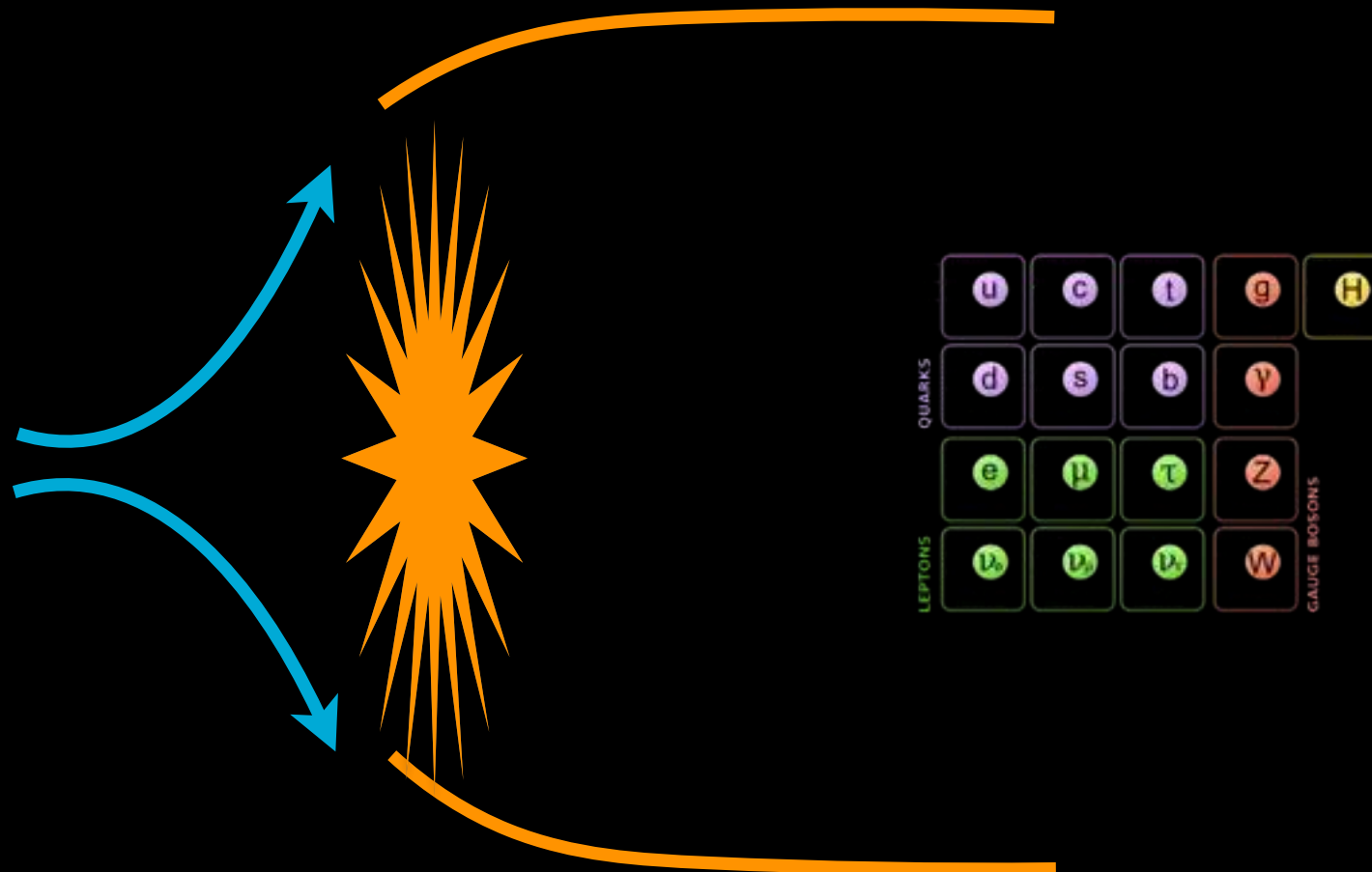
Silverstein & Westphal (2008)

McAllister et. al (2014)

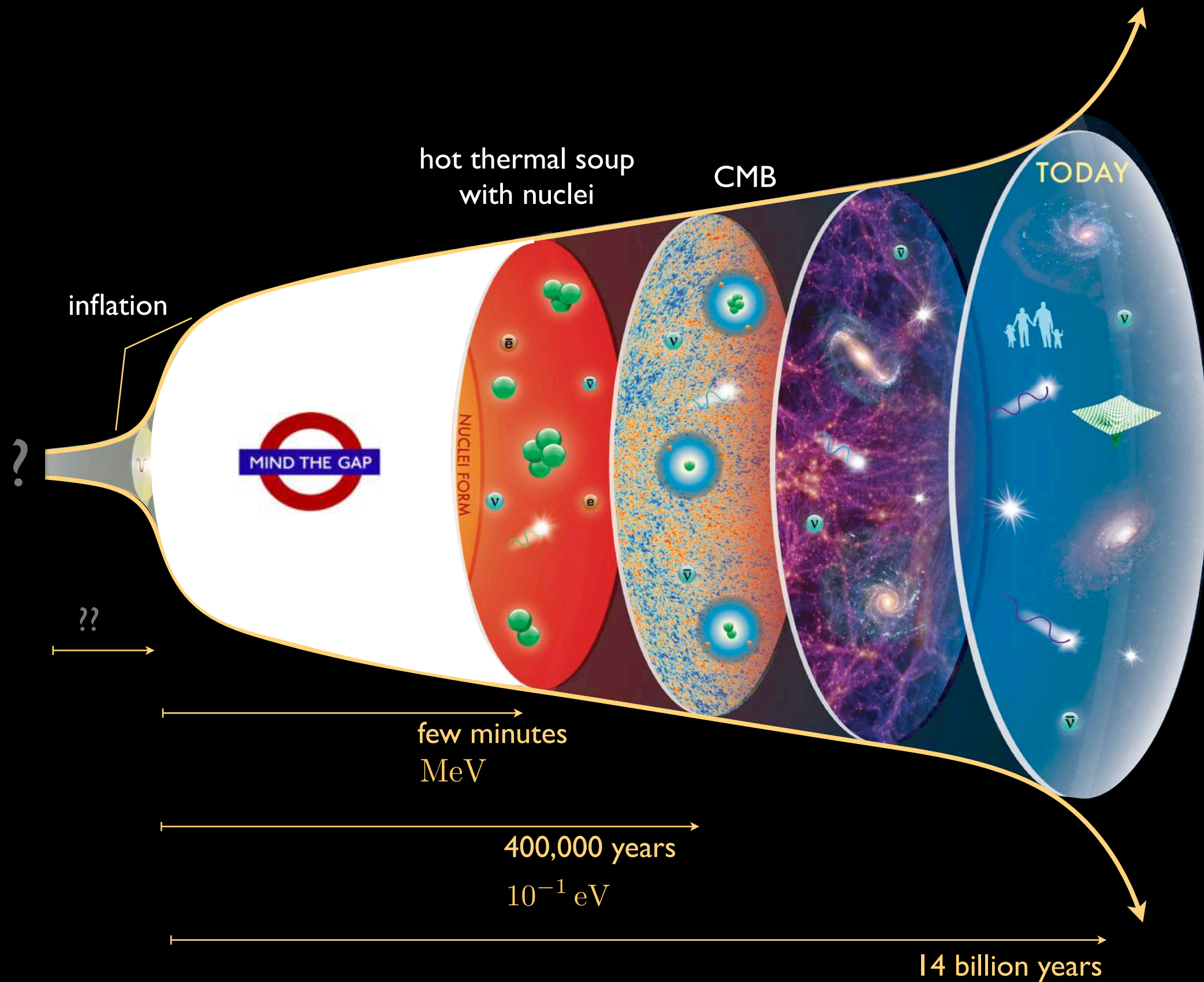
Kallosh & Linde (2014)

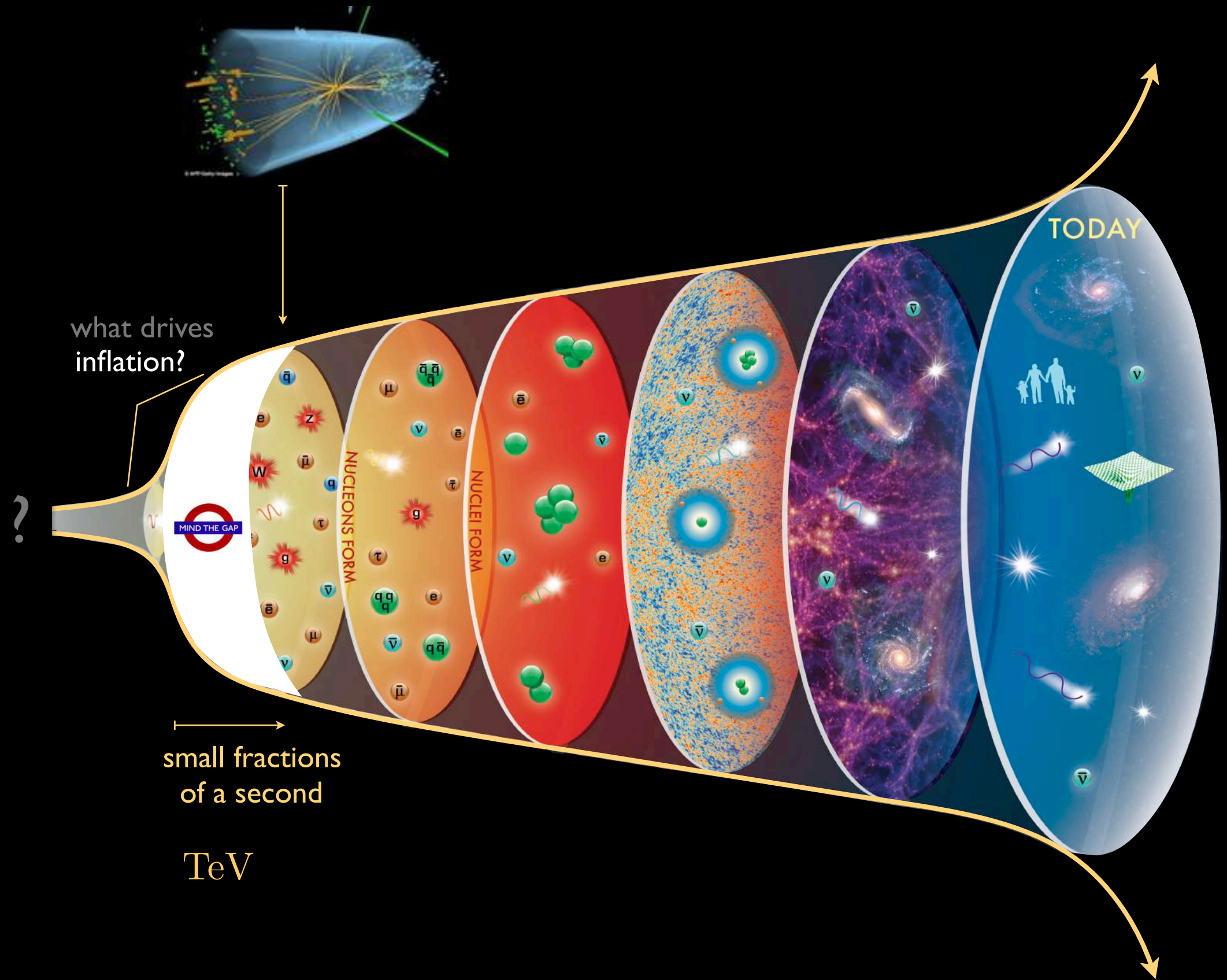


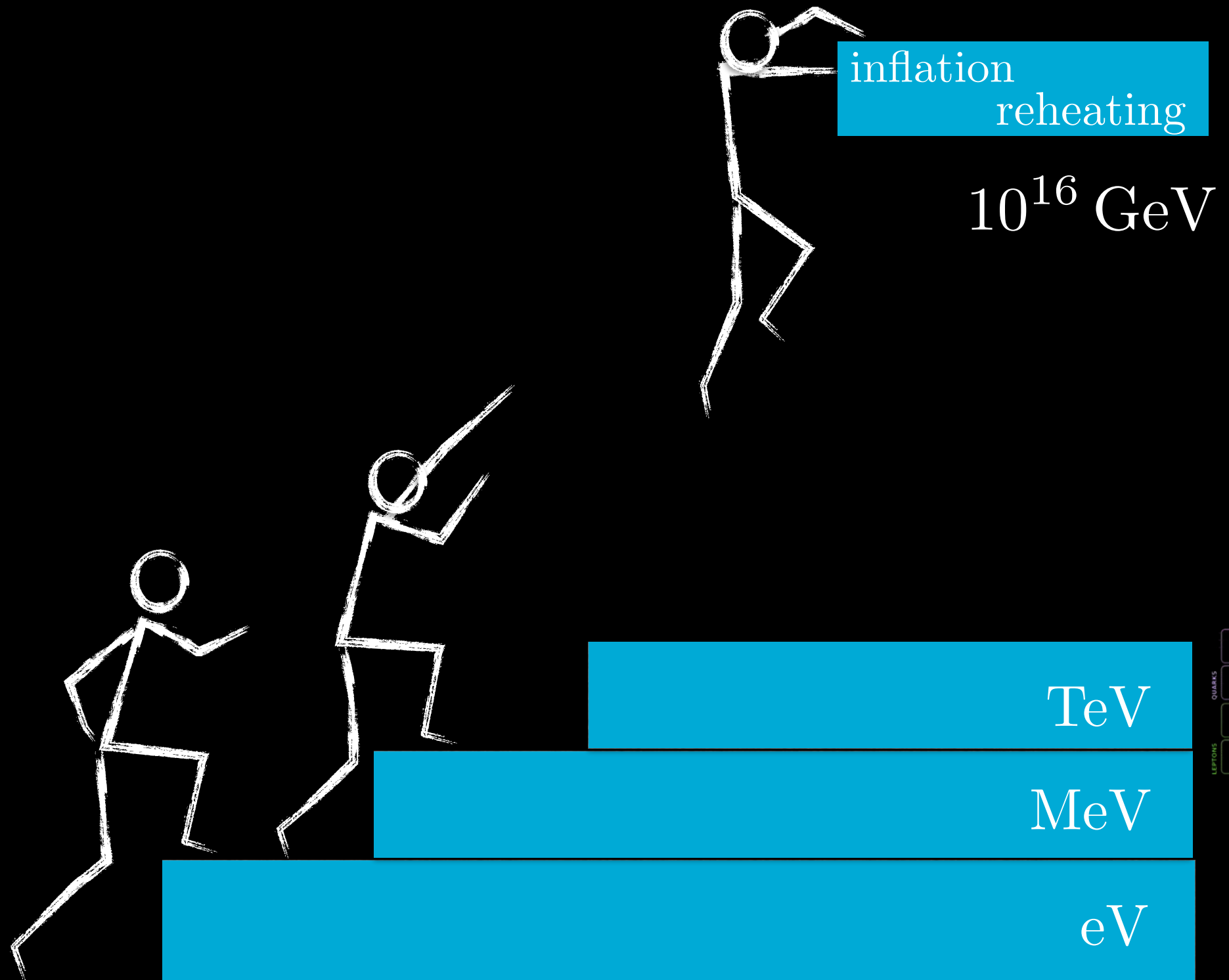
- aftermath of inflation (reheating)?
- Standard Model ?



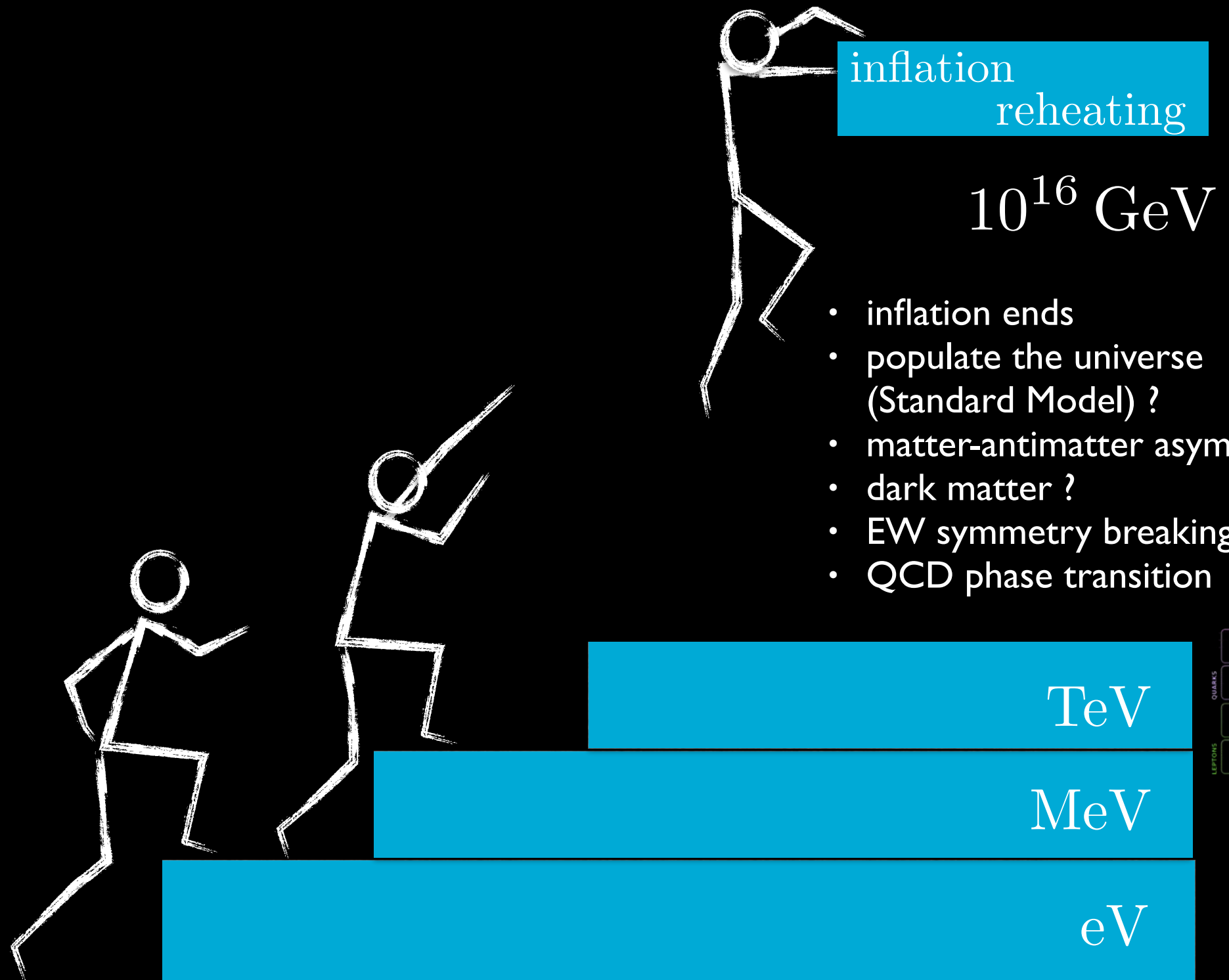
aftermath of inflation: a GAP in our cosmic history







QUARKS	u	c	t	b	H
	d	s	b	y	
	e	μ	τ	Z	
LEPTONS	ν_e	ν_μ	ν_τ	W	
GAUGE BOSONS					



QUARKS	u	c	t	b	H
	d	s	b	y	
	e	μ	τ	Z	
LEPTONS	ν_e	ν_μ	ν_τ	W	
GAUGE BOSONS					

modeling inflation and its aftermath

SIMPLE

problem oriented

COMPLEX

modeling inflation and its aftermath

SIMPLE

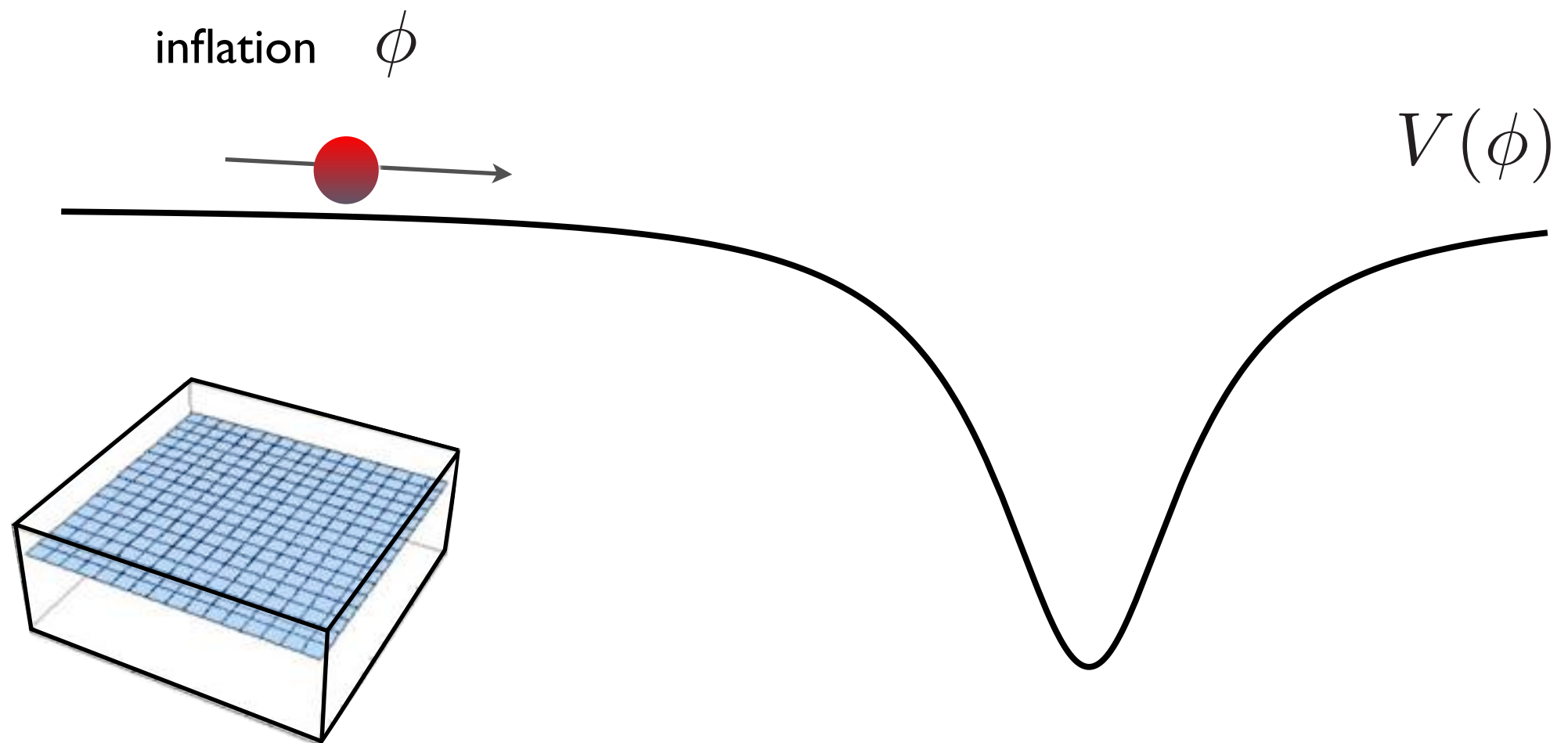
problem oriented

COMPLEX

during inflation

$$\ddot{a}(t) \propto V(\phi) - \dot{\phi}^2 > 0$$

$$\dot{\phi}^2 \ll V(\phi)$$



ending inflation

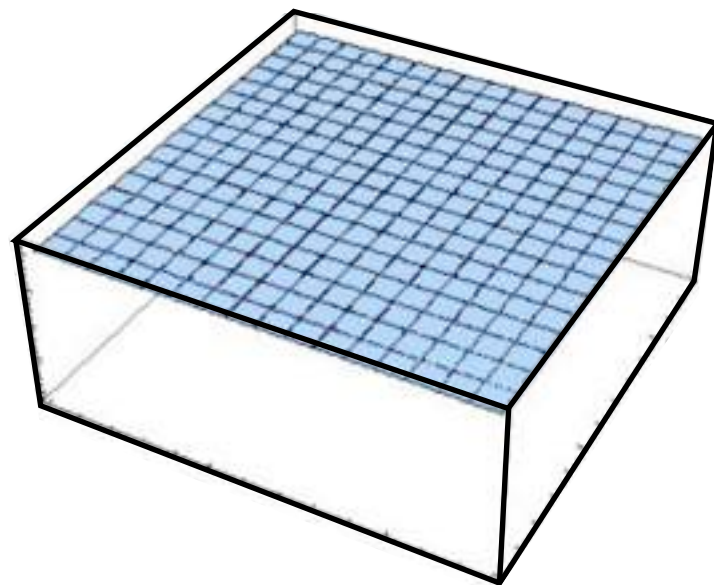
$$\ddot{a}(t) \propto V(\phi) - \dot{\phi}^2 < 0$$

$$\frac{1}{2}\dot{\phi}^2 \sim V(\phi)$$

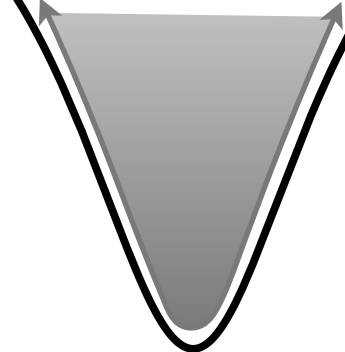
inflation

end: oscillatory regime

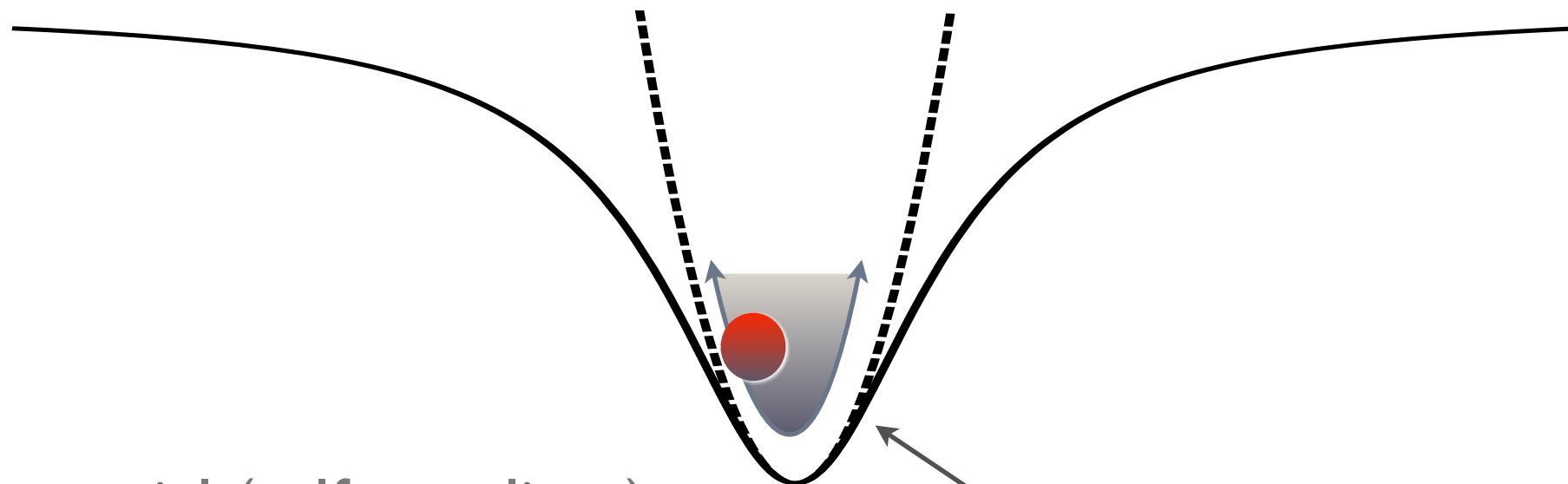
$V(\phi)$



ϕ



detailed dynamics ?



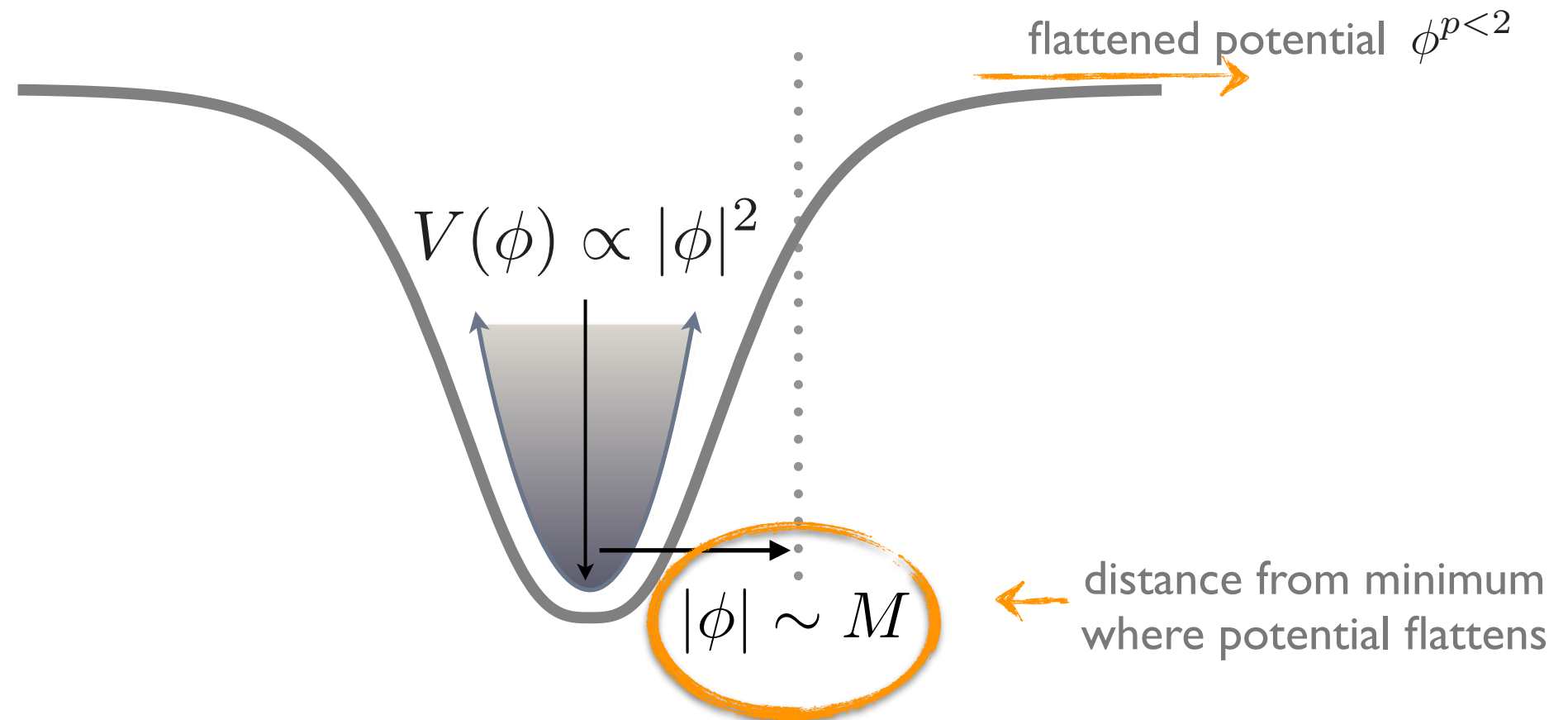
- shape of the potential (self couplings)
- couplings to other fields

 χ, ψ

mass	~ 0.3 MeV/c ²	~ 1.25 GeV/c ²	~ 173.2 GeV/c ²	0	~ 125 GeV/c ²
charge	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0	0
spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	0	0
	u up	c charm	t top	g gluon	H Higgs boson
	d down	s strange	b bottom	γ photon	
	e electron	μ muon	τ tau	Z Z boson	
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W W boson	

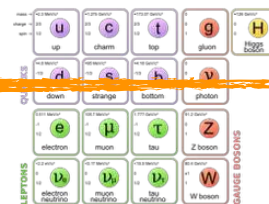
Traschen & Brandenburger (1990)
Kofman, Linde & Starobinsky (1994)
Shtanov, Traschen & Brandenberger (1995)
Kofman, Linde & Starobinsky (1997)
review: MA, Kaiser, Karouby & Hertzberg (2014)

end of inflation in “simple” models



- shape of the potential (self couplings)

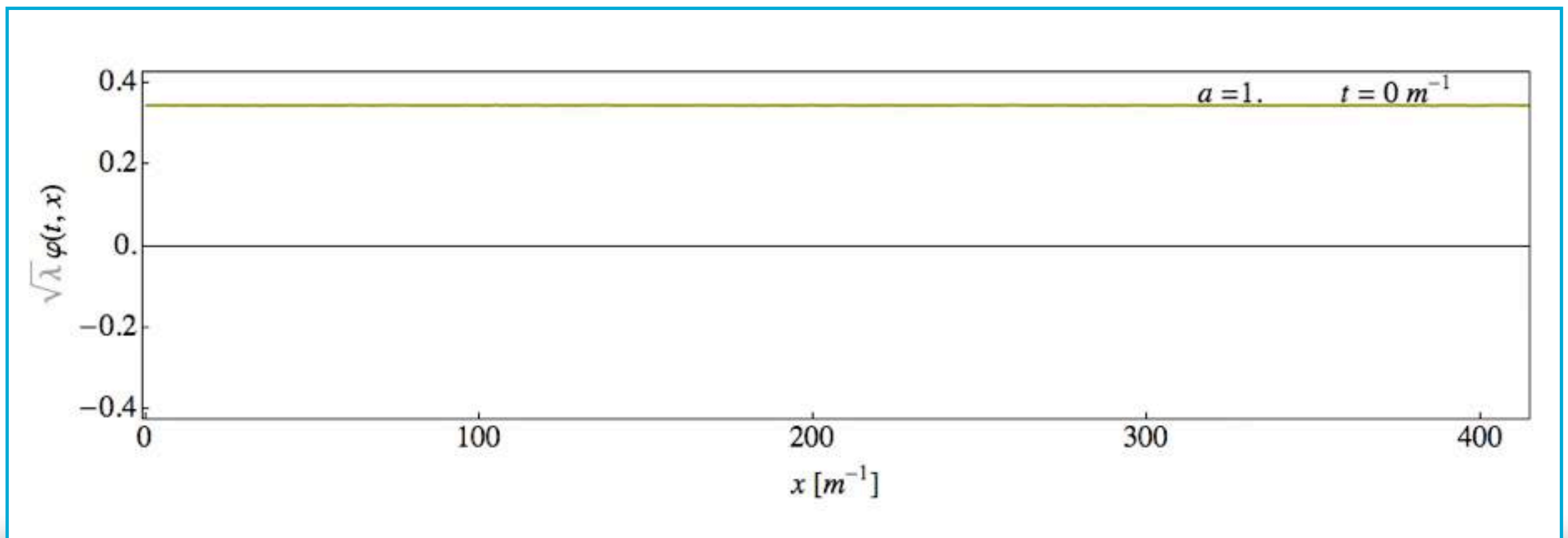
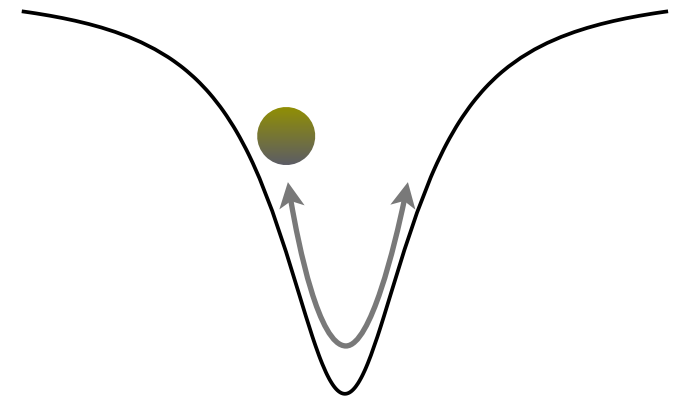
- ~~couplings to other fields~~



χ, ψ

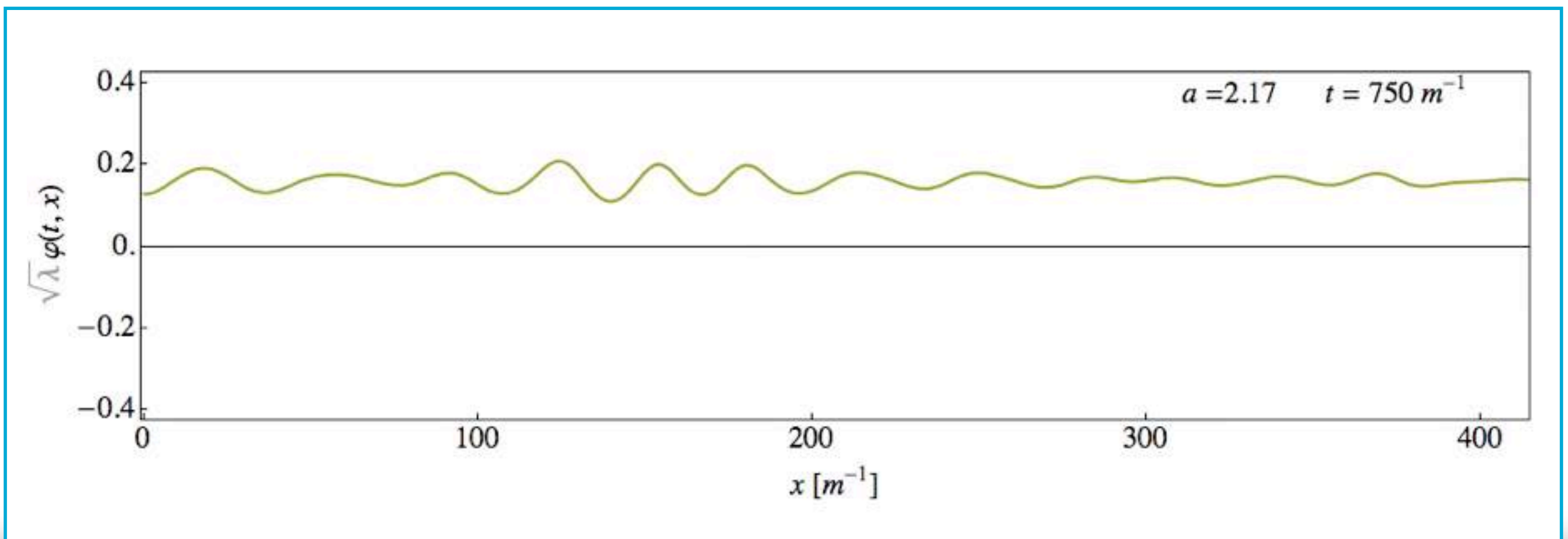
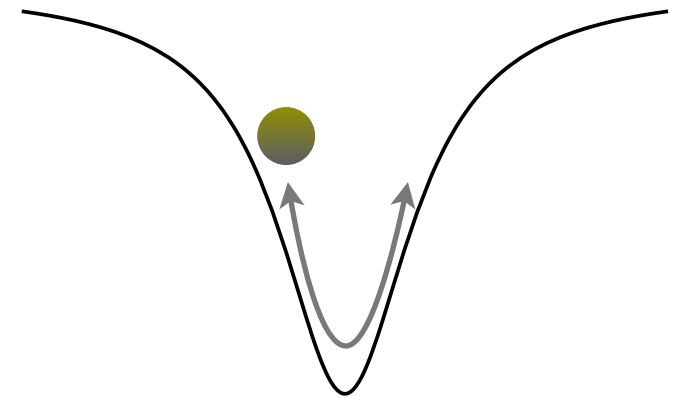
dynamics after inflation

$$\square\phi = V'(\phi)$$



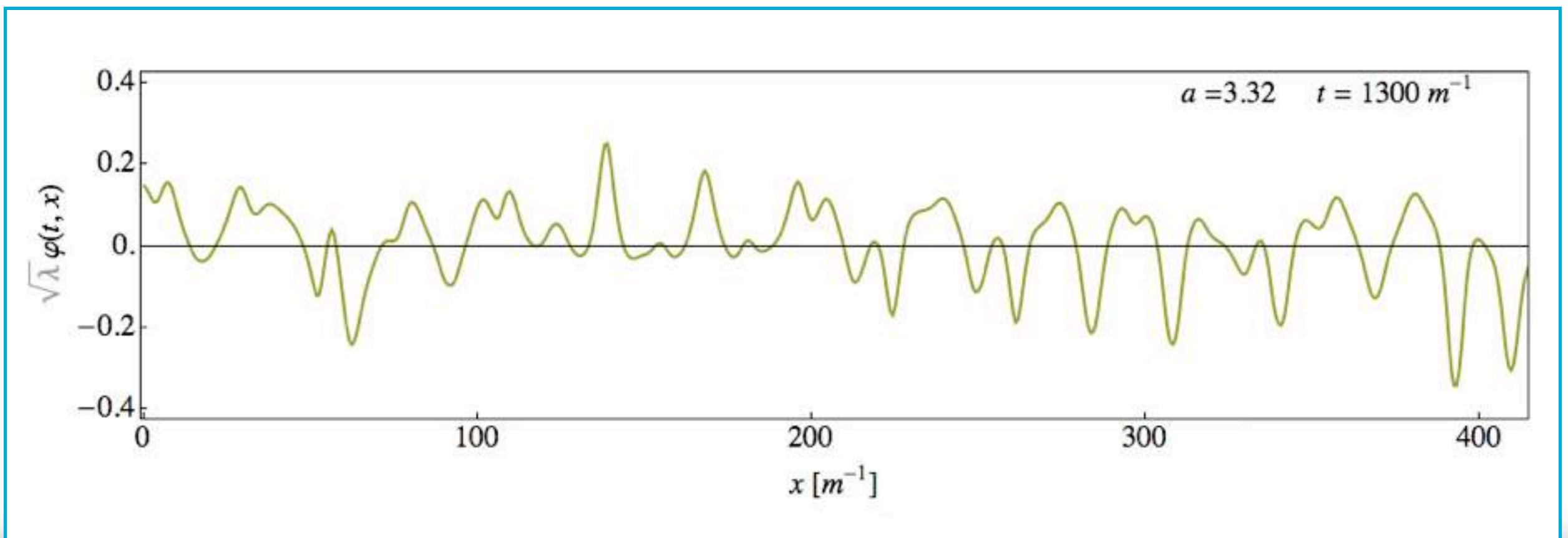
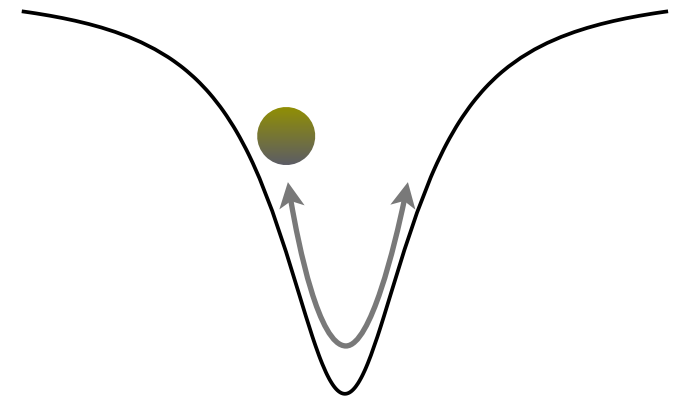
dynamics after inflation

$$\square\phi = V'(\phi)$$



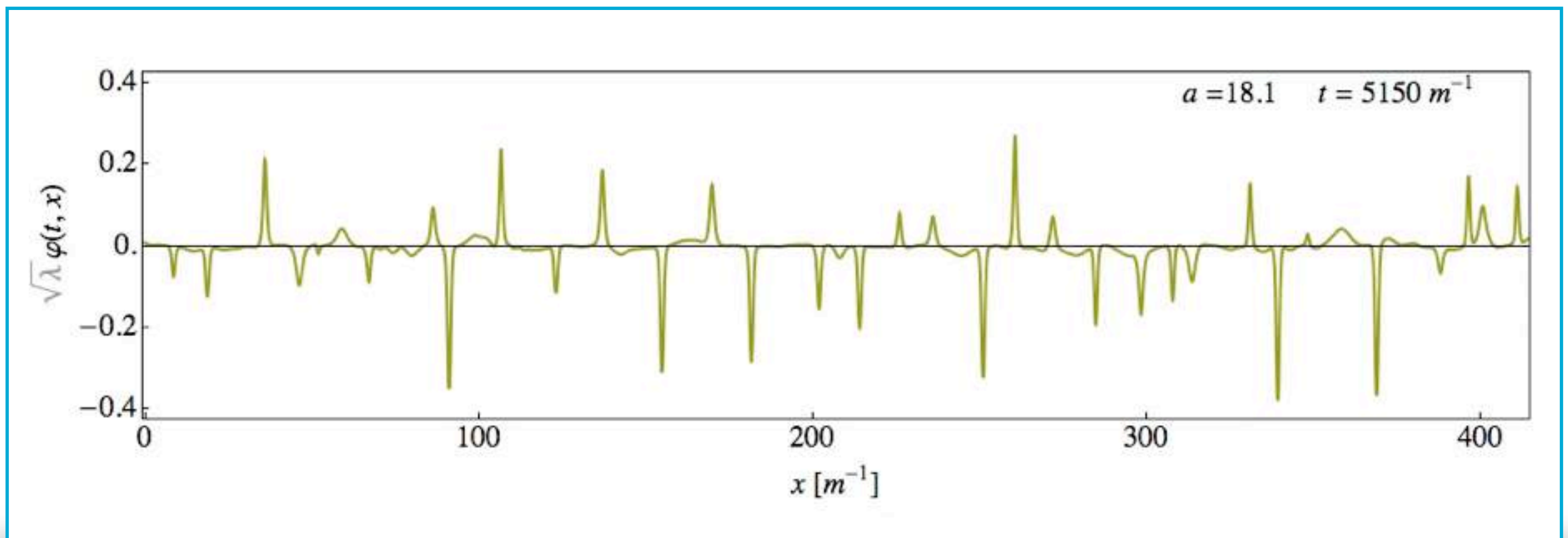
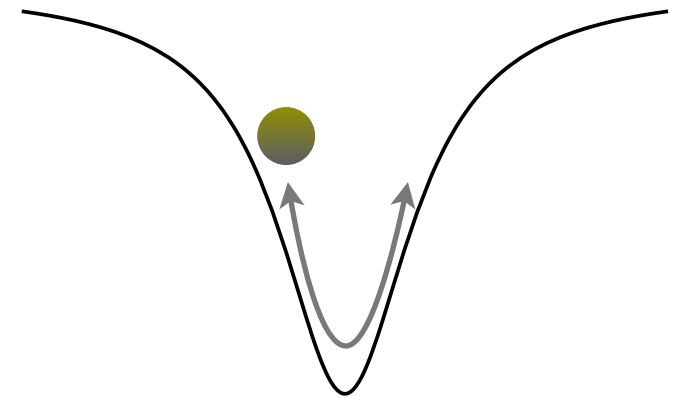
dynamics after inflation

$$\square\phi = V'(\phi)$$



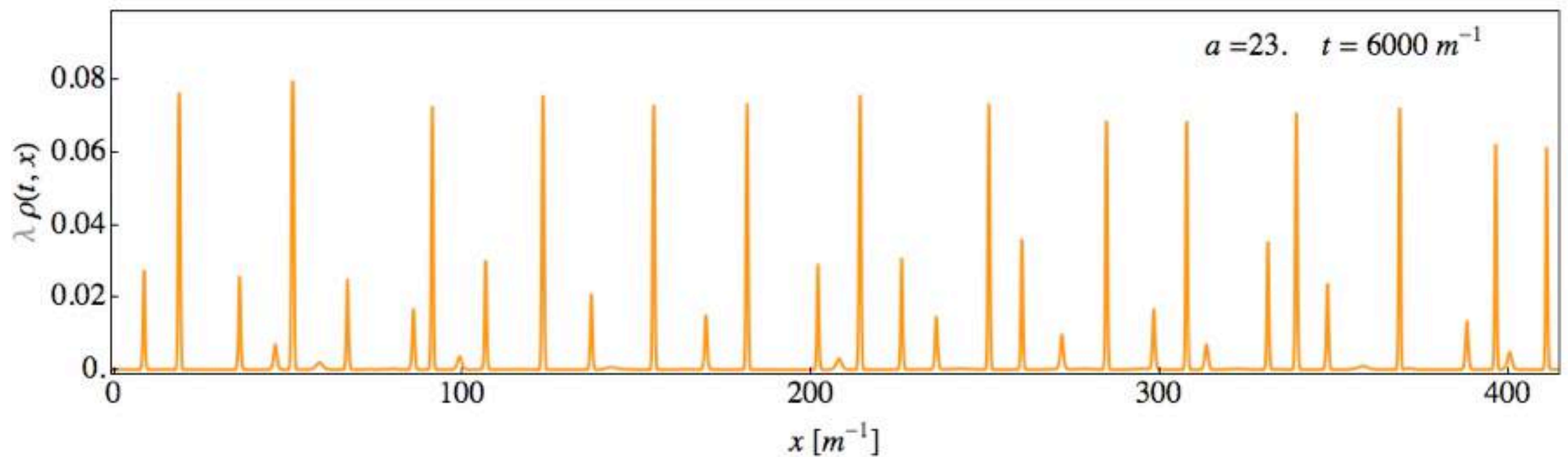
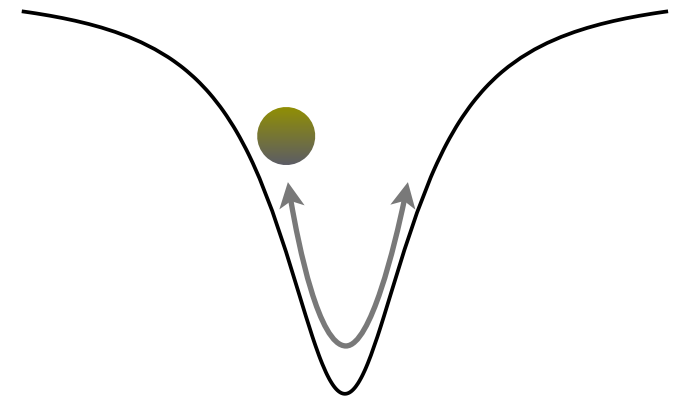
dynamics after inflation

$$\square\phi = V'(\phi)$$

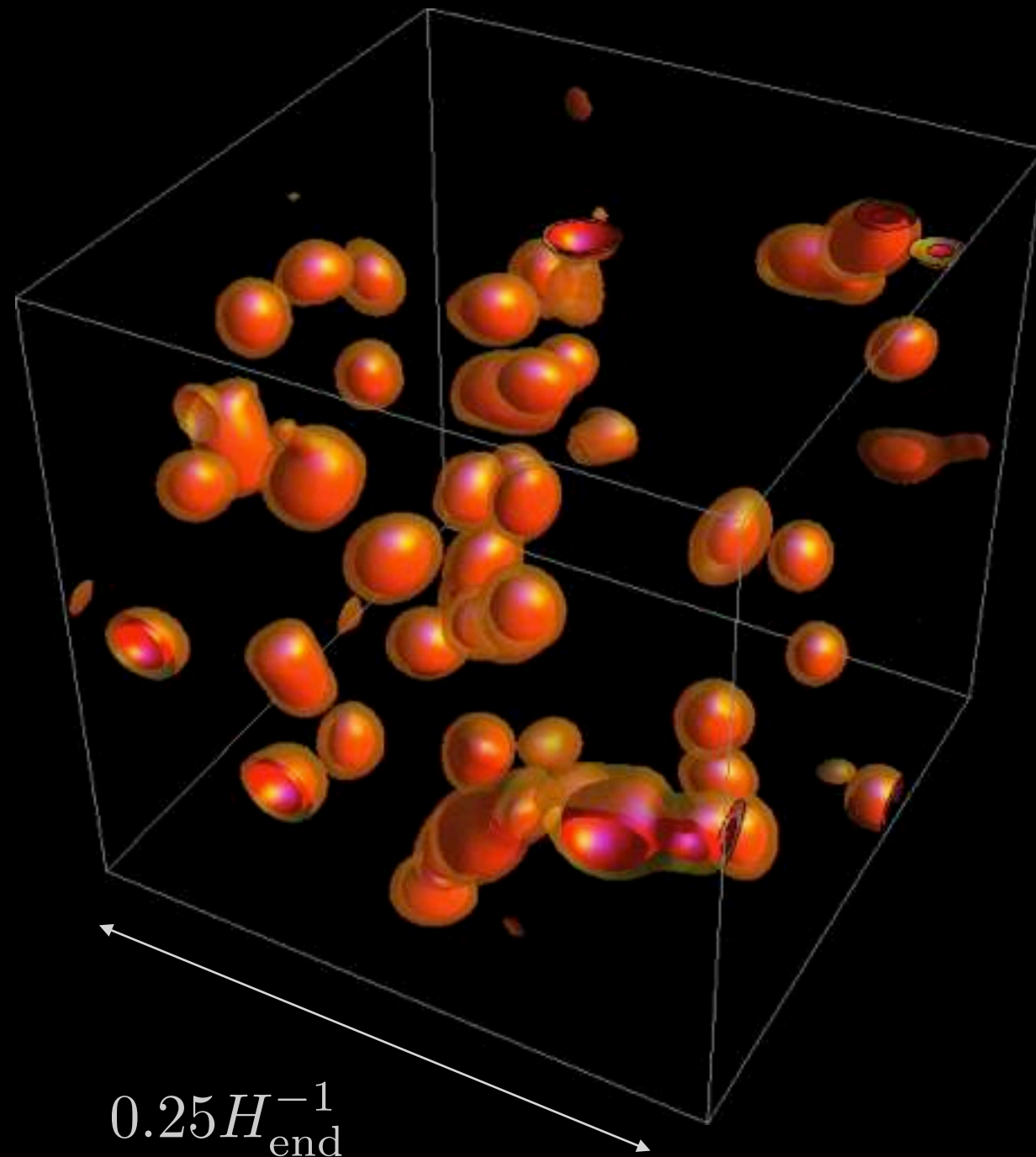


dynamics after inflation

$$\square \varphi = V'(\varphi)$$



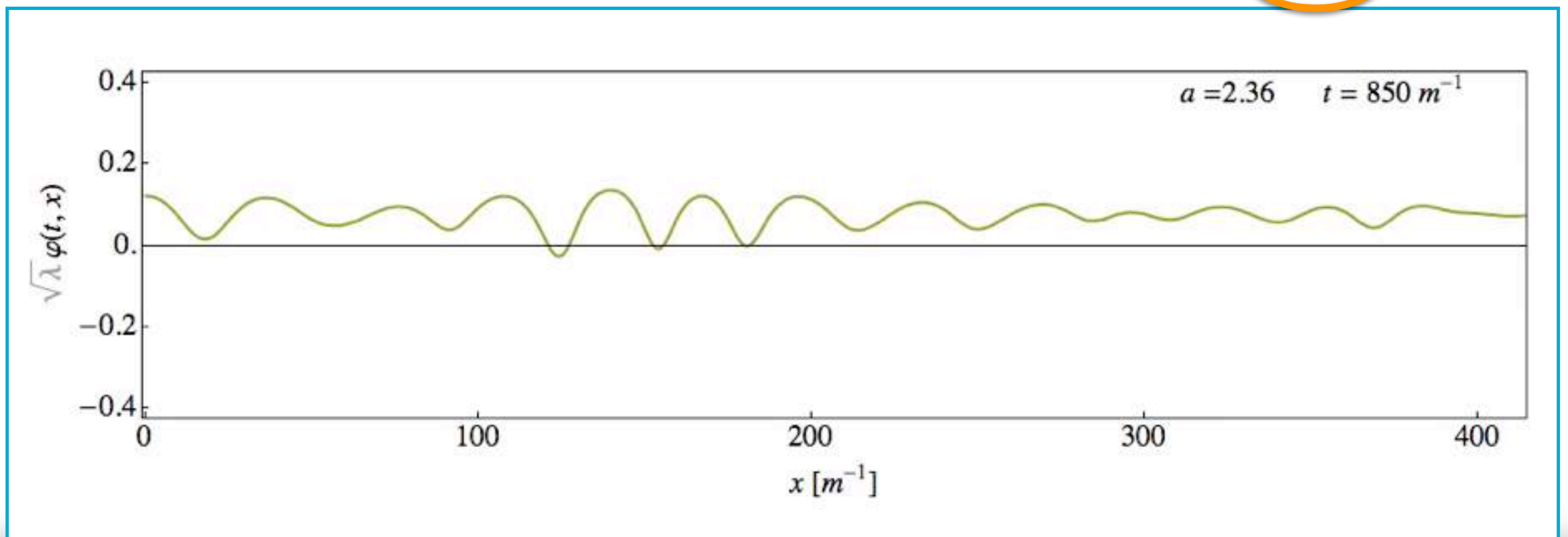
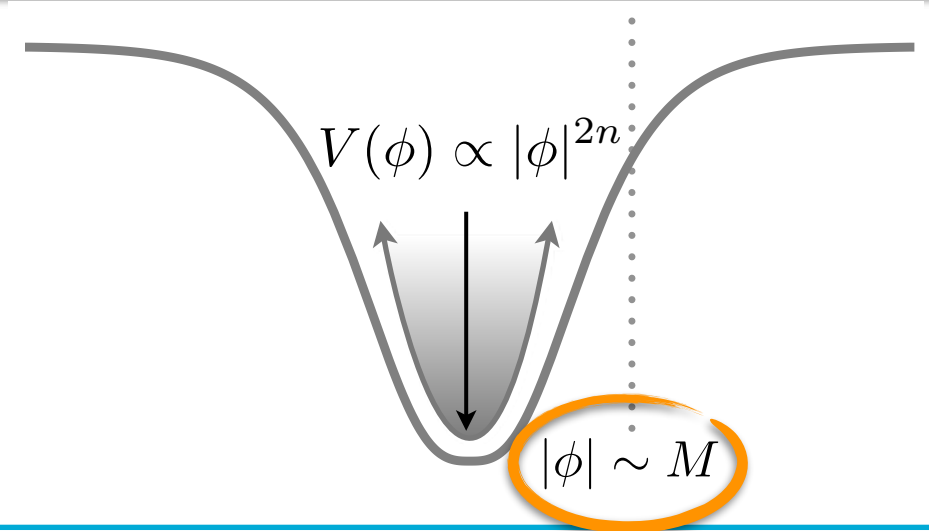
now in 3D: (iso-density surfaces)



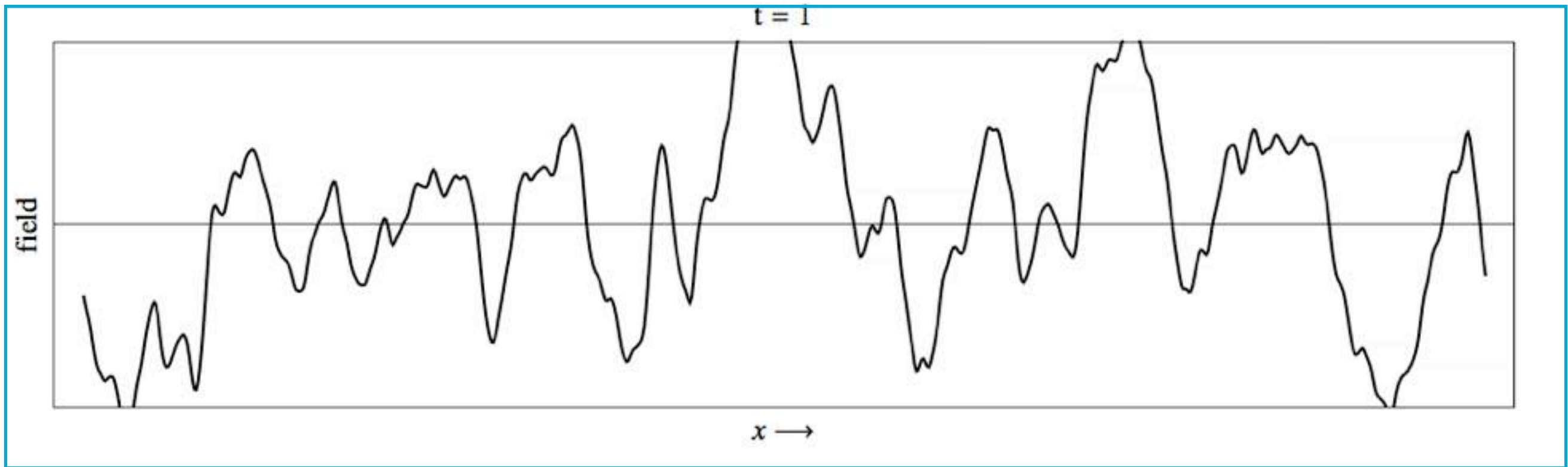
MA, Easter, Finkel, Flaugher & Hertzberg (2011)

condition for rapid fragmentation ?

$$\frac{\text{growth-rate of fluctuations}}{\text{expansion rate}} \sim \frac{m_{\text{pl}}}{M} \gg 1$$

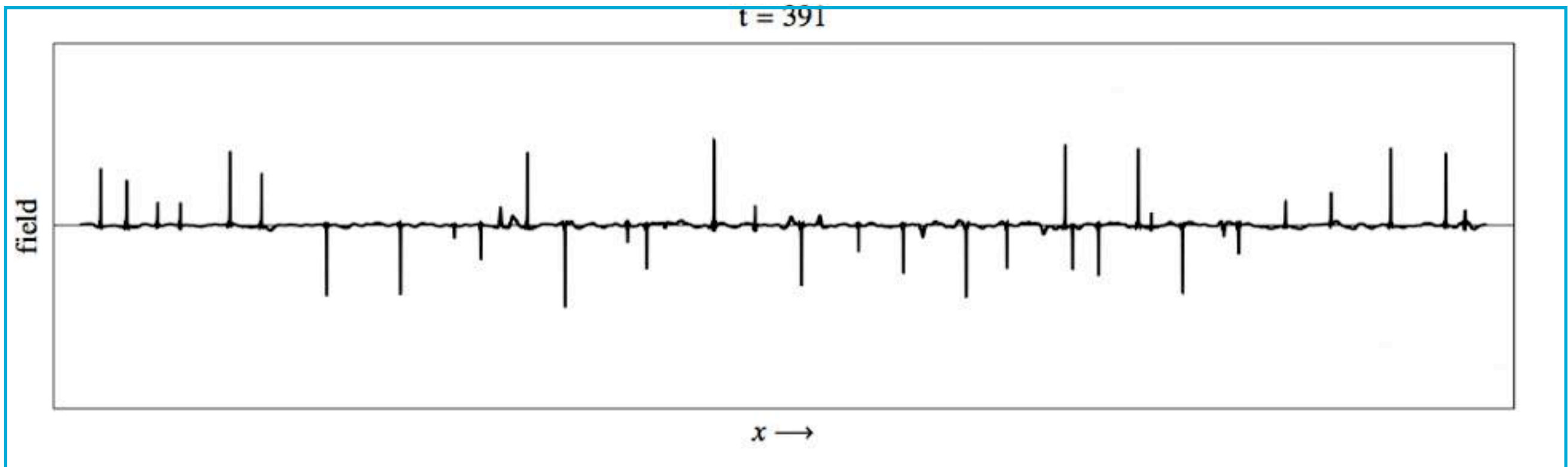


insensitive to initial conditions



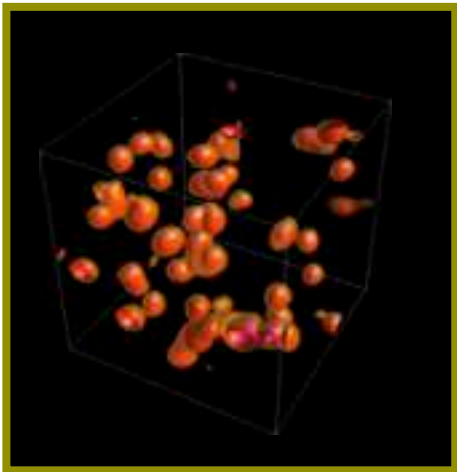
simulation of “quasi-thermal” example in Farhi et. al 2008

insensitive to initial conditions



simulation of “quasi-thermal” example in Farhi et. al 2008

lumps ?



(1) oscillatory (2) spatially localized (3) **very long lived**

$$\mathcal{L} = T(X, \varphi) - V(\varphi)$$

$$T(X, \varphi) = X + \xi_2 X^2 + \xi_3 \varphi X^2 + \dots$$

$$V(\varphi) = \frac{1}{2}\varphi^2 + \frac{\lambda_3}{3}\varphi^3 + \frac{\lambda_4}{4}\varphi^4 + \frac{\lambda_5}{5}\varphi^5 + \dots$$

$$\Delta = \xi_2 - \lambda_4 + \frac{10}{9}\lambda_3^2 > 0.$$

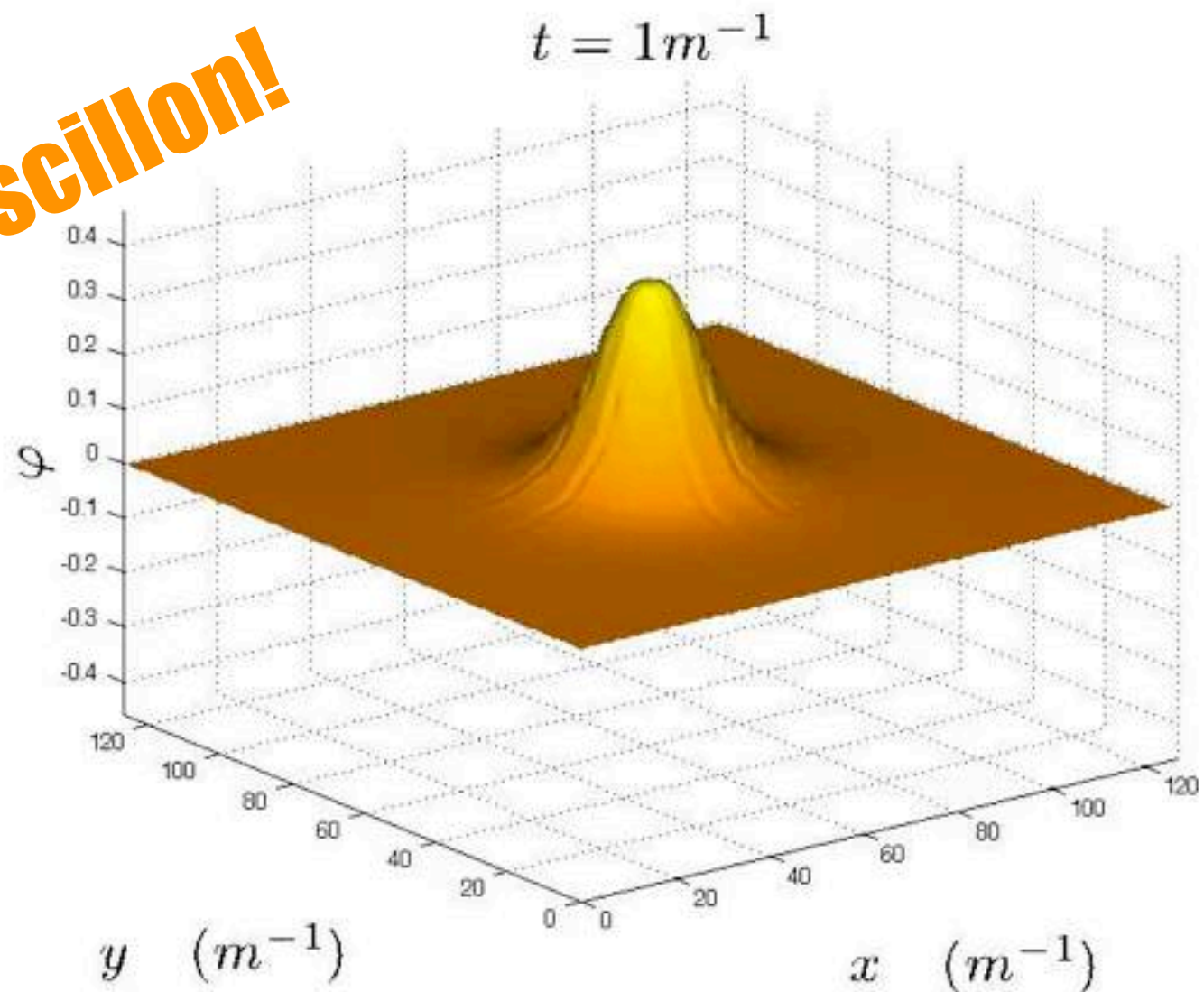
existence and stability:

MA (2013)

MA & Shirokoff (2010)

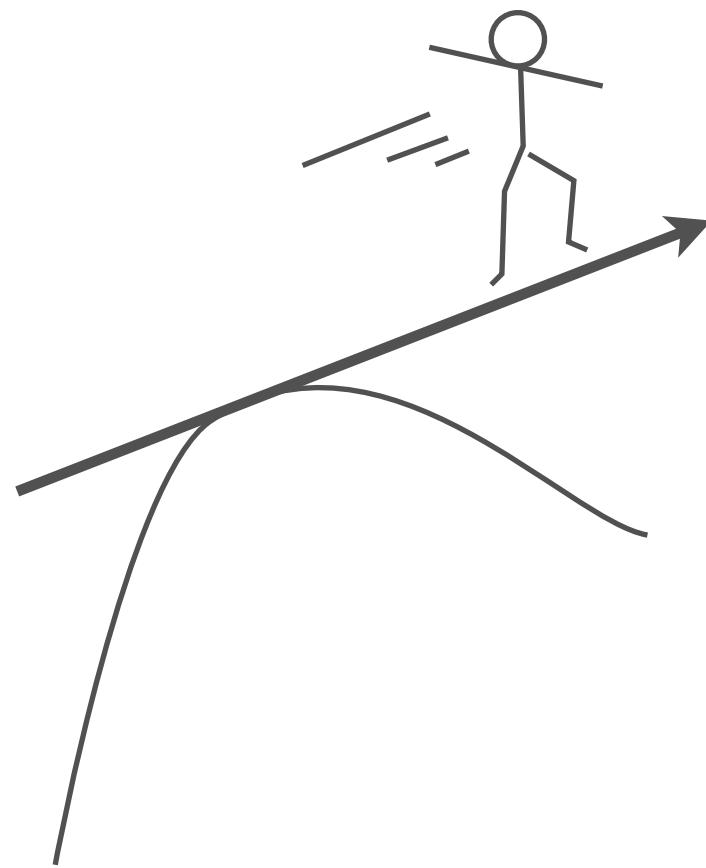
Hertzberg (2011)

oscillon!



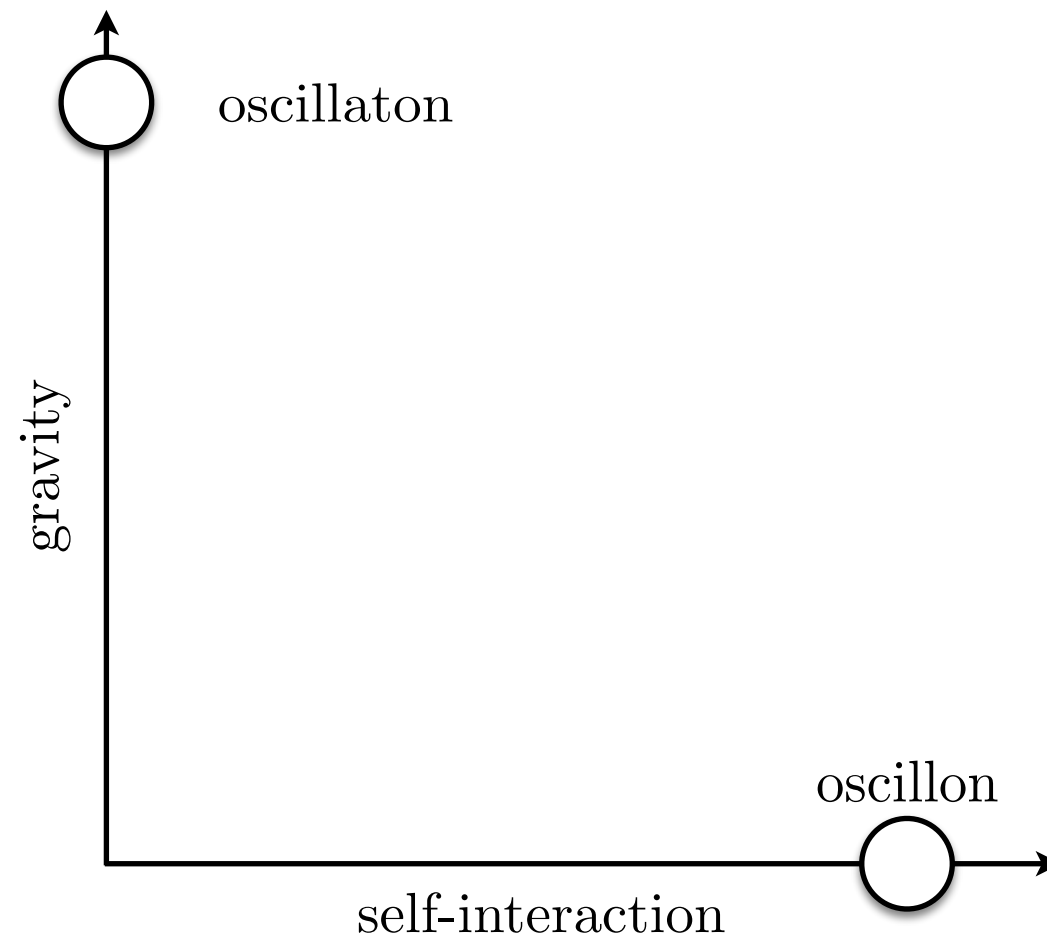
Bogolubsky & Makhankov (1976), Gleiser (1994), Copeland, Gleiser and Mueller et al. (1995) ...

tangential digression



oscillons and friends

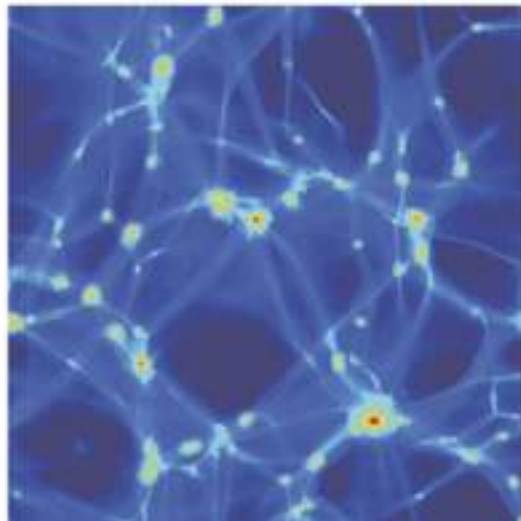
non-topological solitons



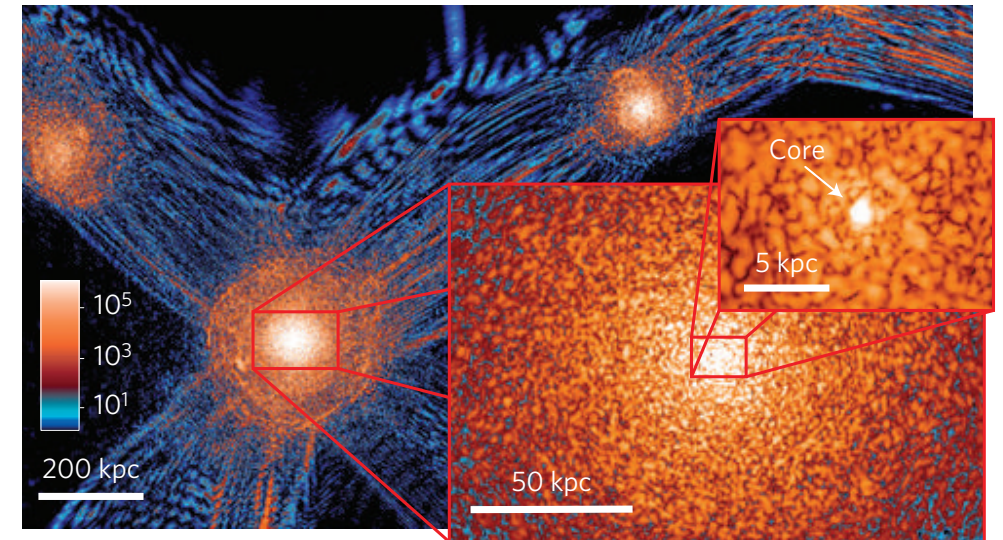
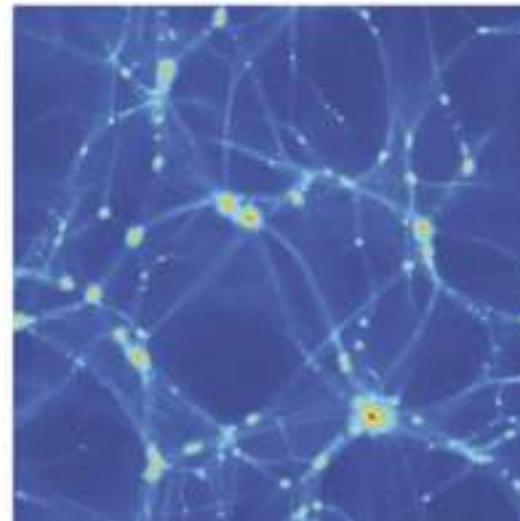
oscillons and friends

- late universe

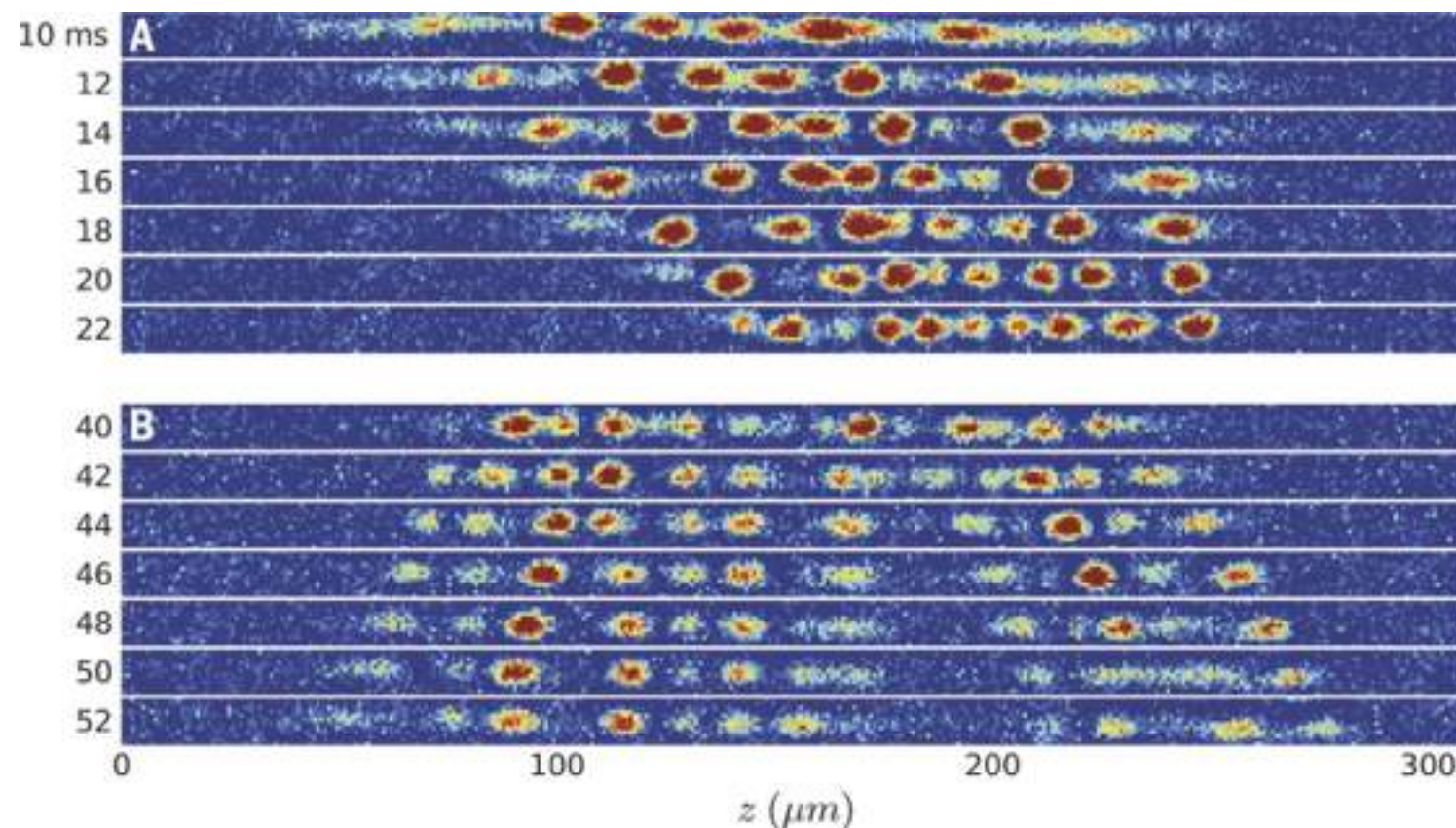
(a) ψ DM



(b) CDM



Schive et. al (2014)

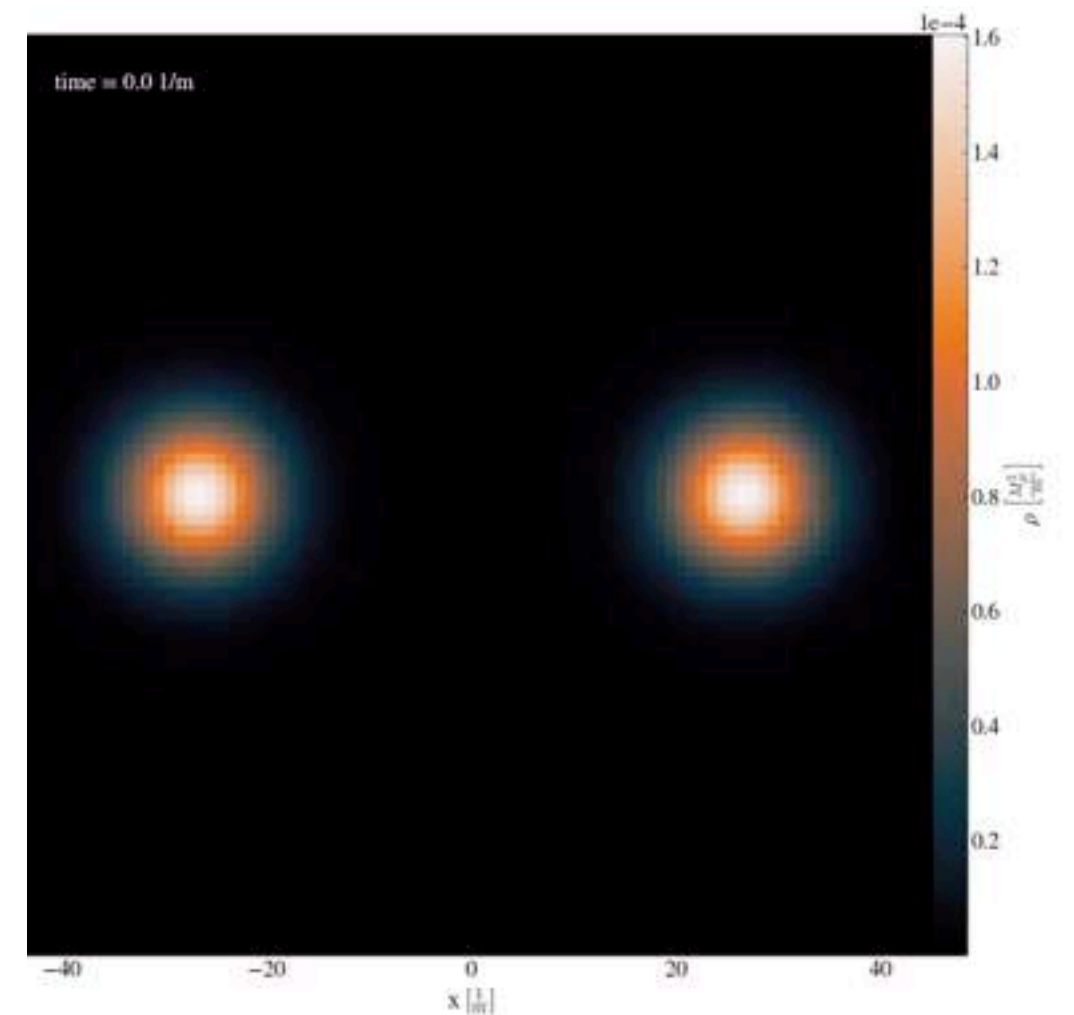
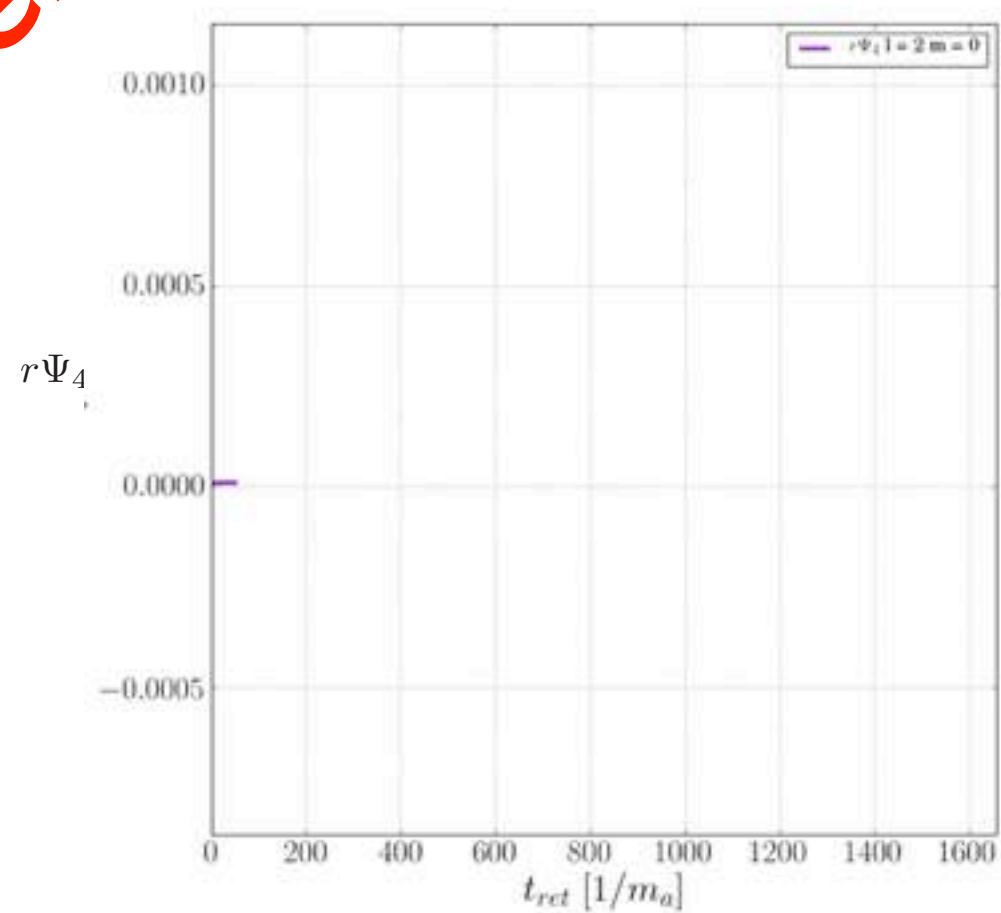


Nguyen, Luo & Hulet (2017)

gravitational waves from dense oscillaton collisions

relevance for late universe processes — axion stars

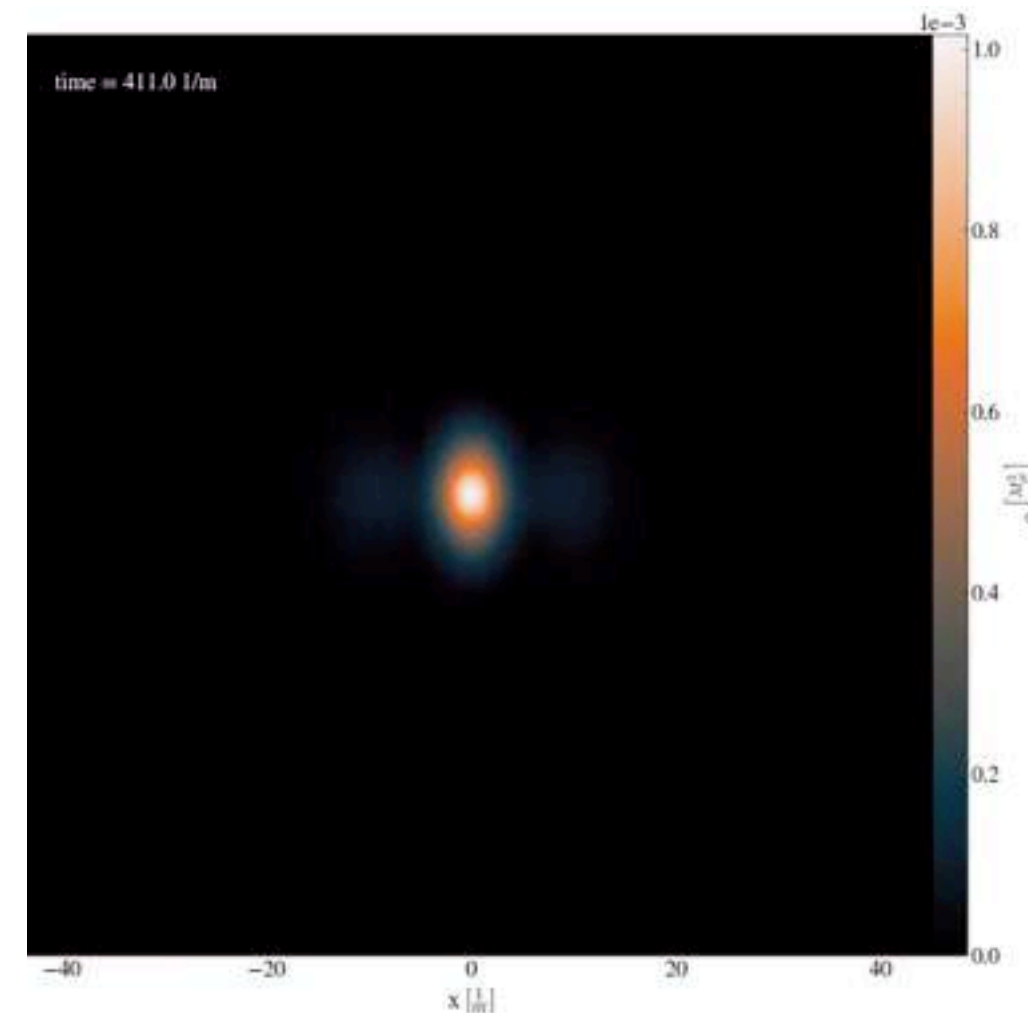
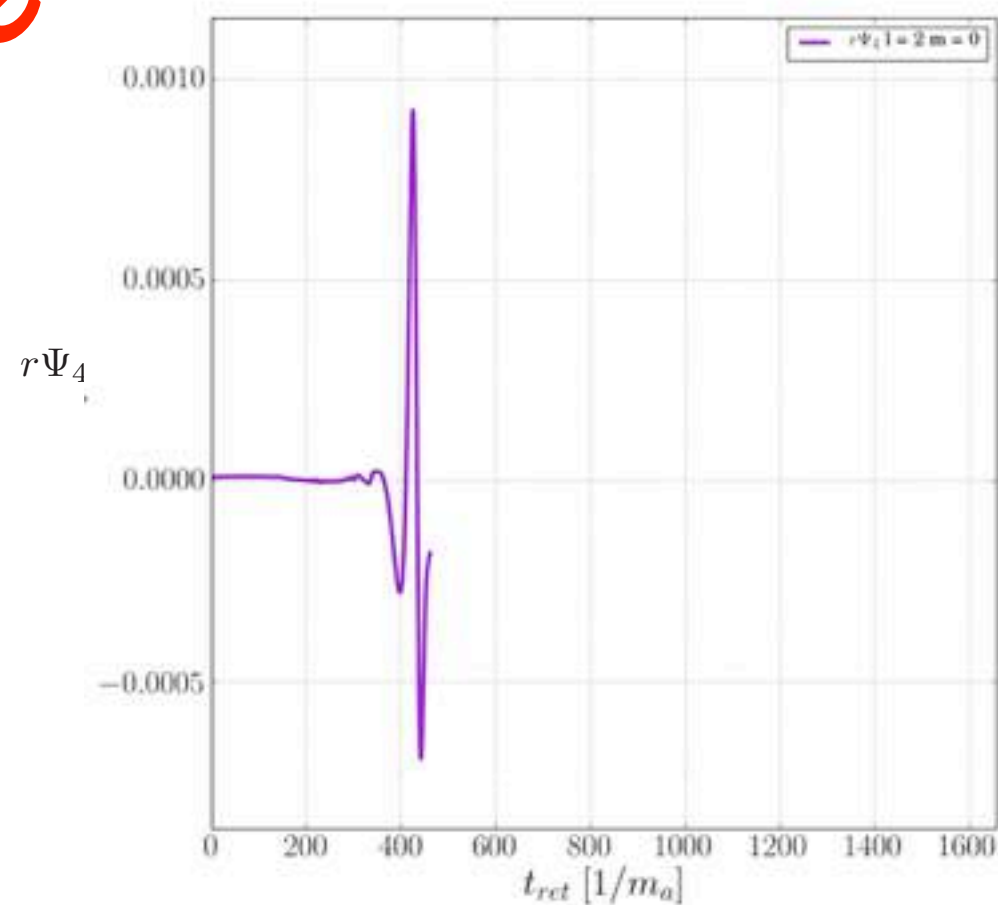
Helfer, Lim, Garcia & Amin (in prep)



gravitational waves from dense oscillaton collisions

relevance for late universe processes — axion stars

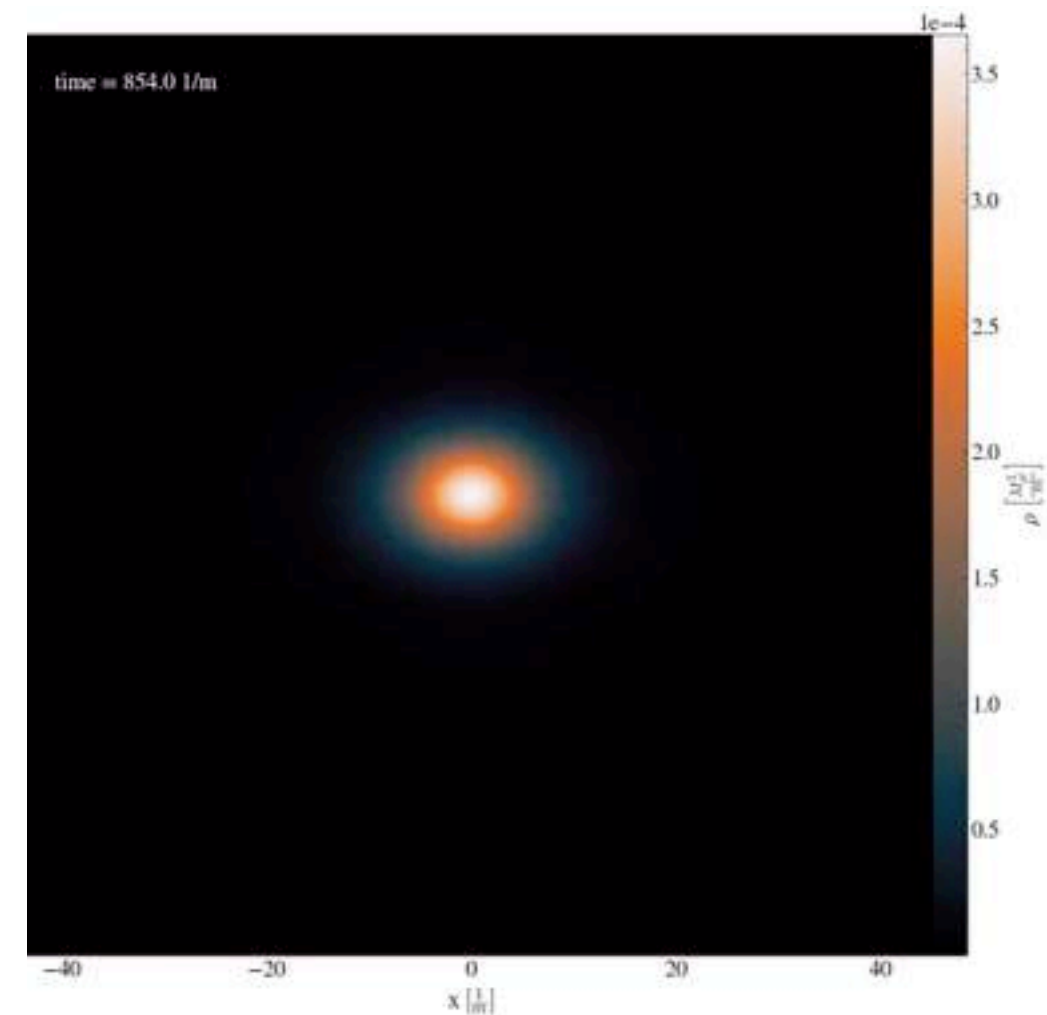
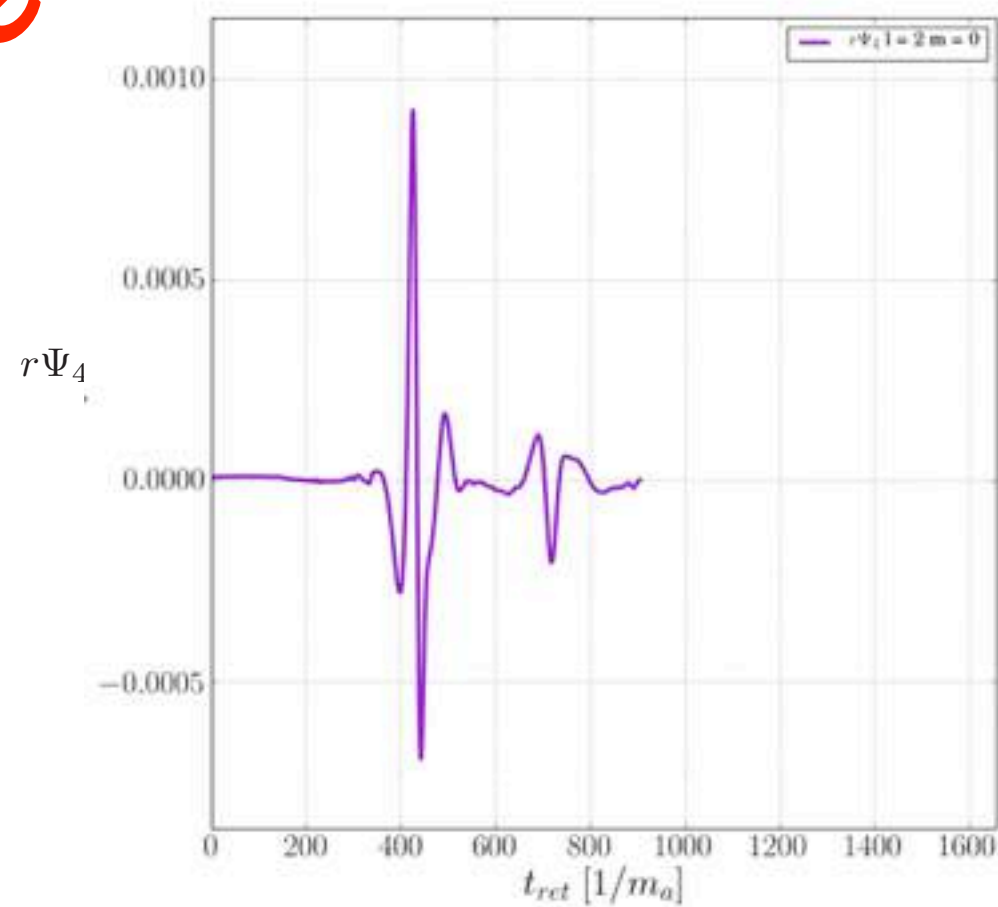
Helfer, Lim, Garcia & Amin (in prep)



gravitational waves from dense oscillaton collisions

relevance for late universe processes — axion stars

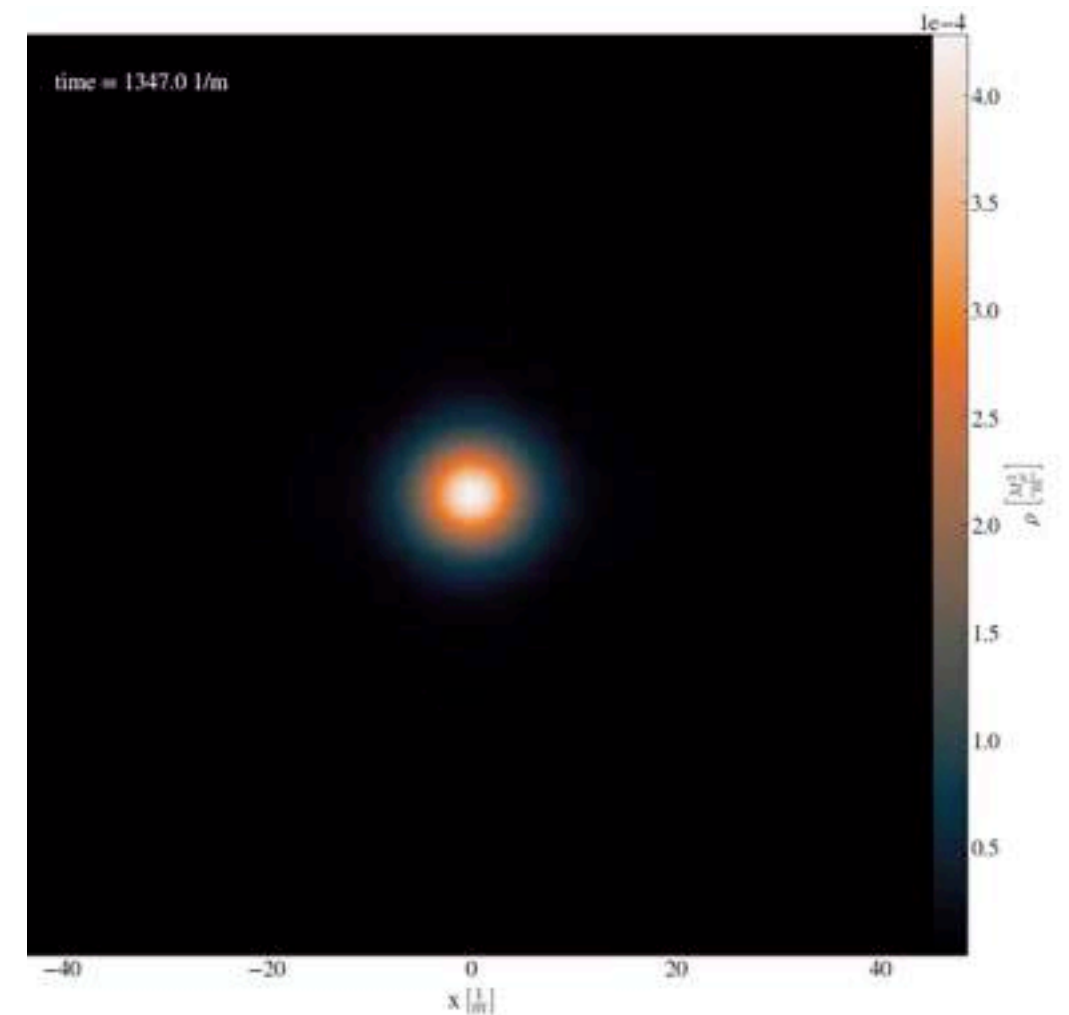
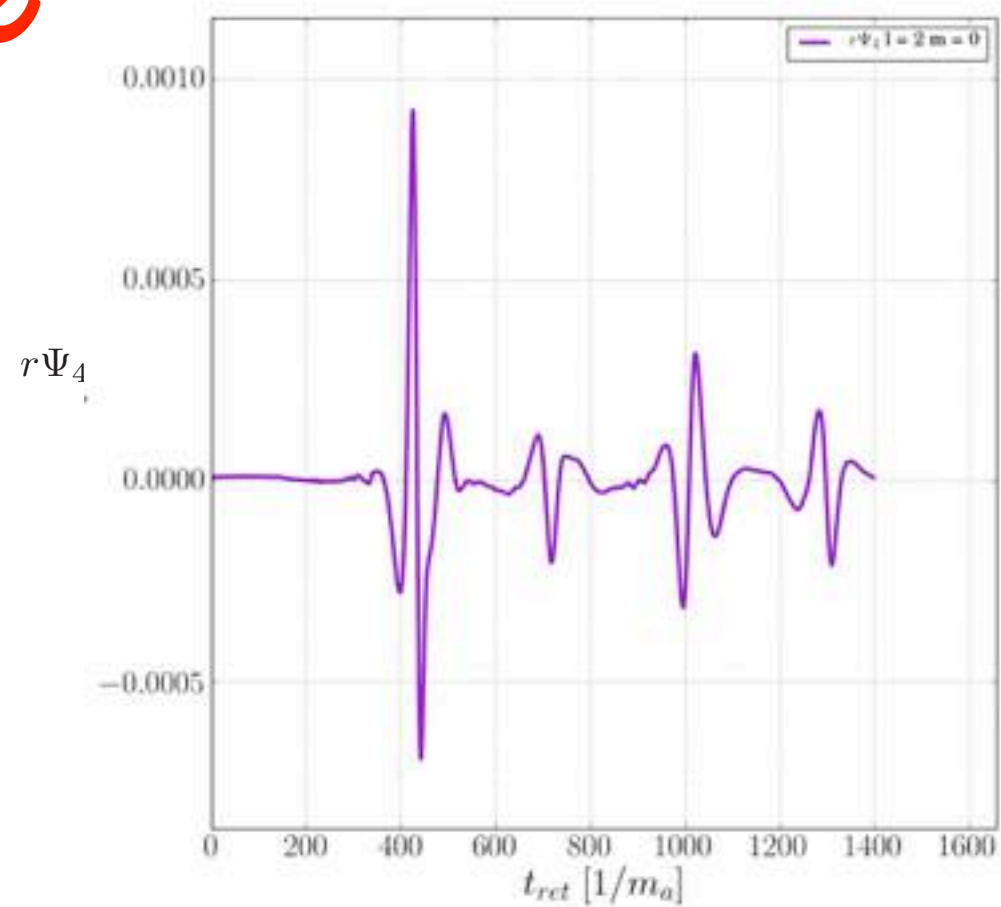
Helfer, Lim, Garcia & Amin (in prep)



gravitational waves from dense oscillaton collisions

relevance for late universe processes — axion stars

Helfer, Lim, Garcia & Amin (in prep)

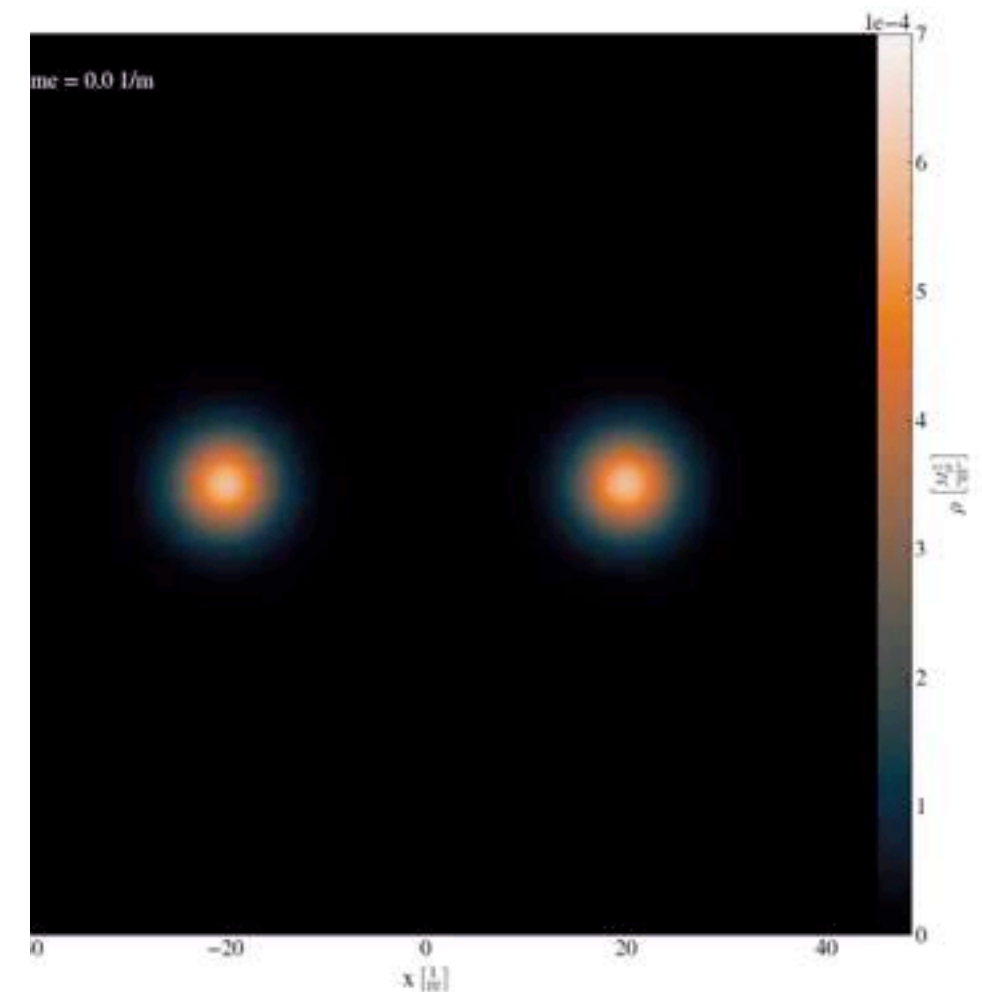
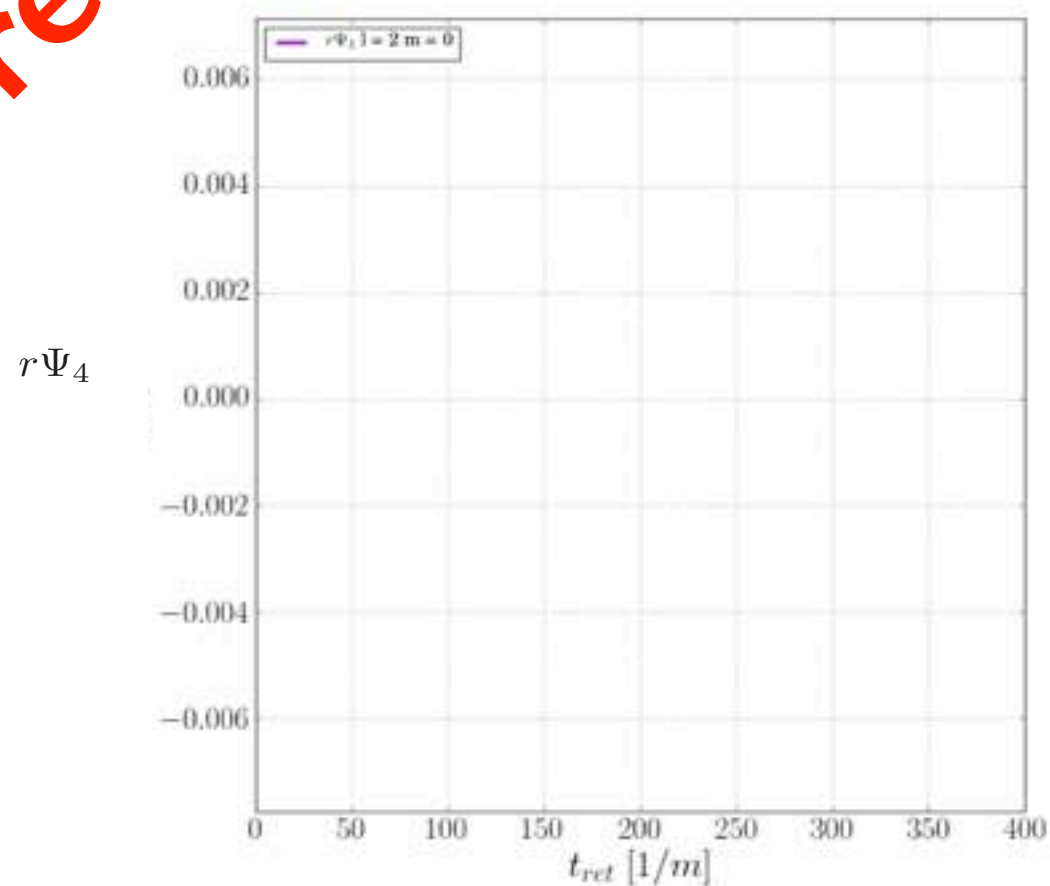


gravitational waves from oscillaton collisions

* critical case

relevance for late universe processes — axion stars

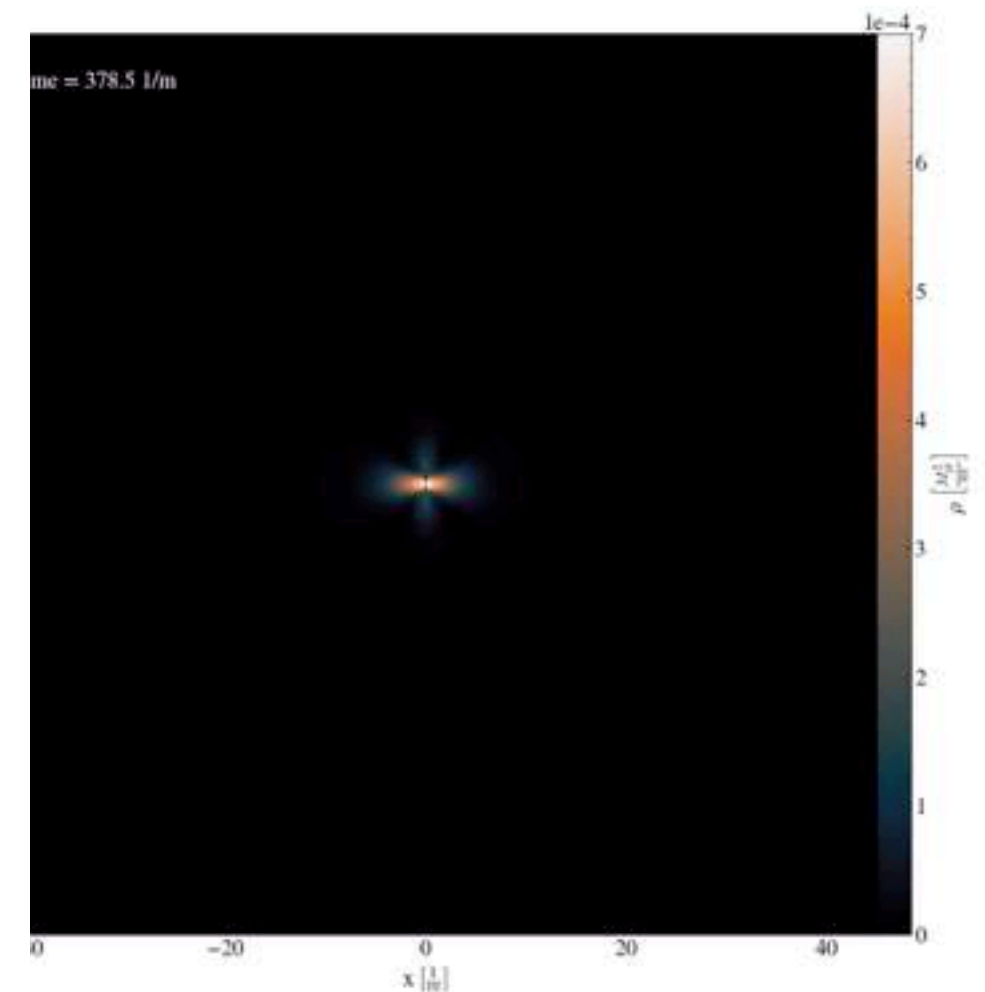
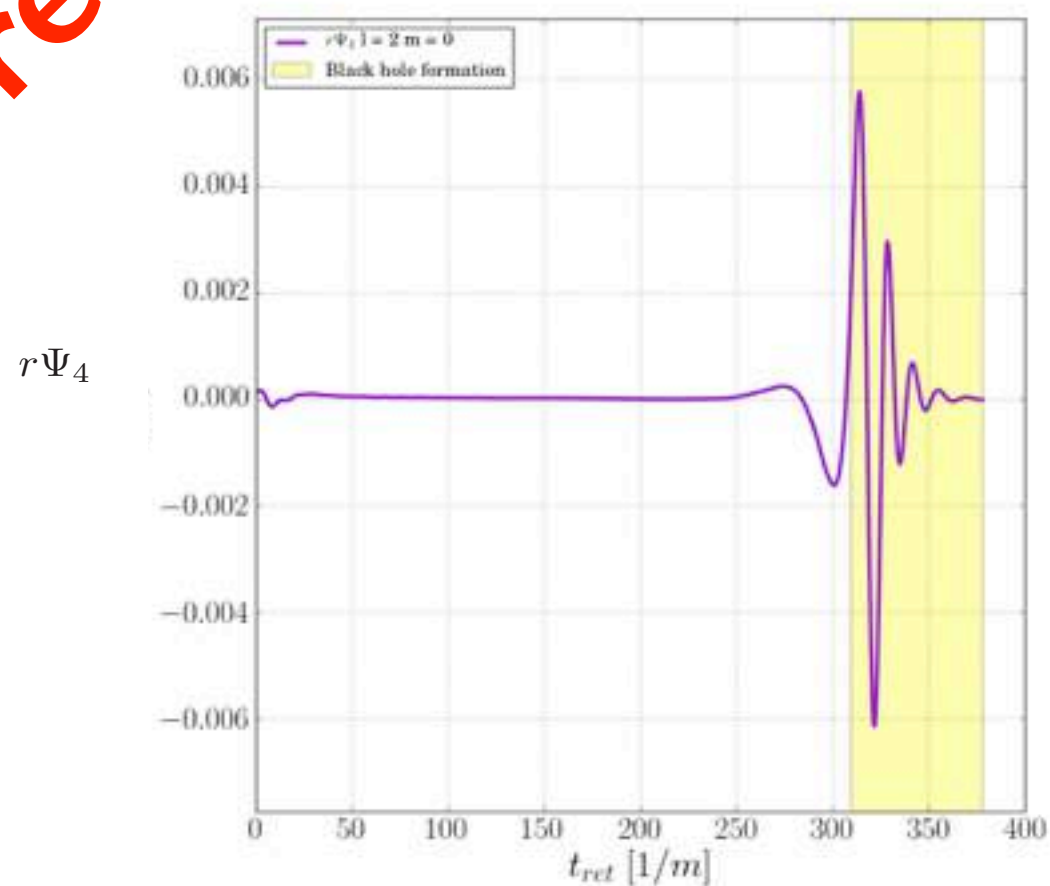
Helfer, Lim, Garcia & Amin (in prep)



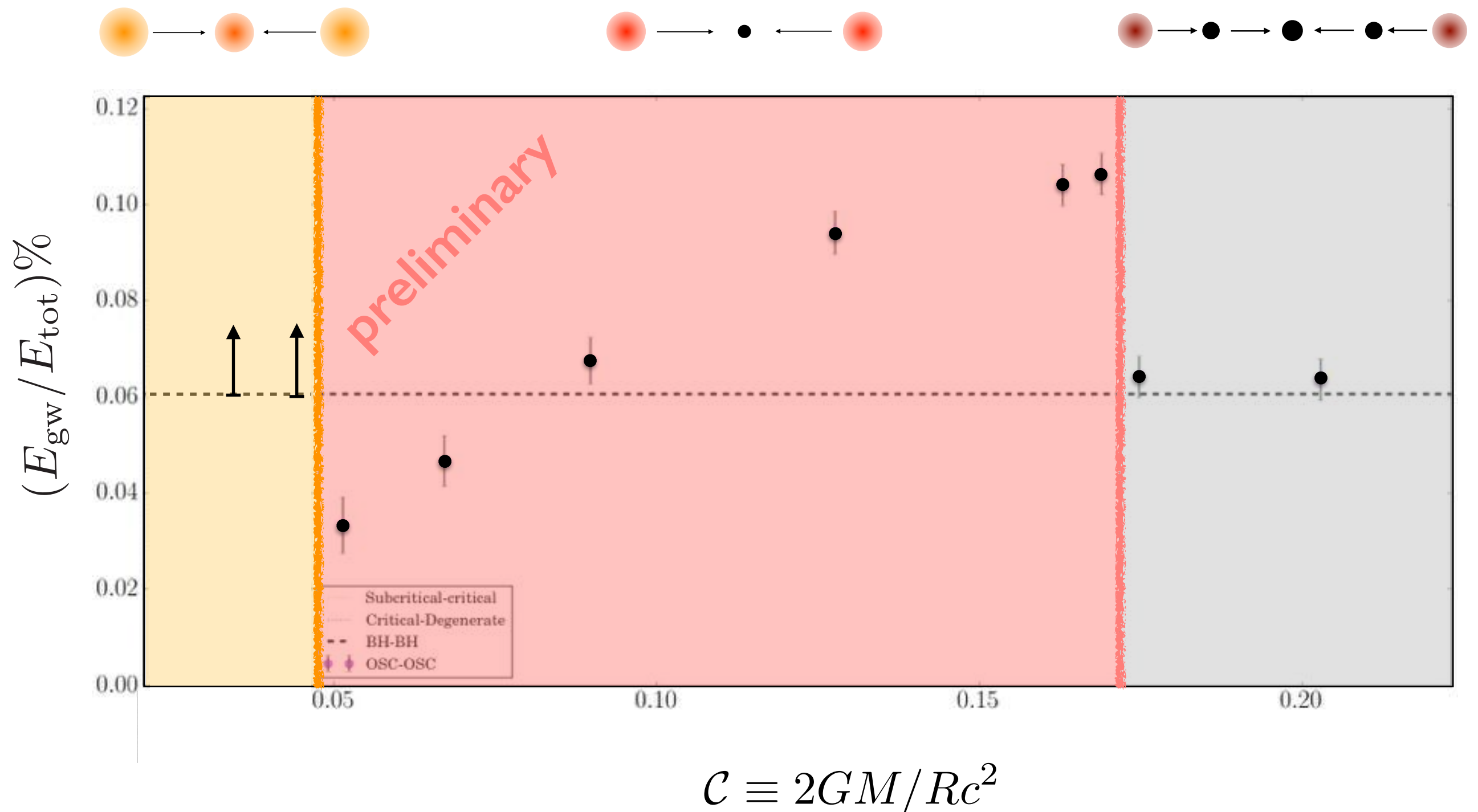
gravitational waves from oscillaton collisions

relevance for late universe processes — axion stars

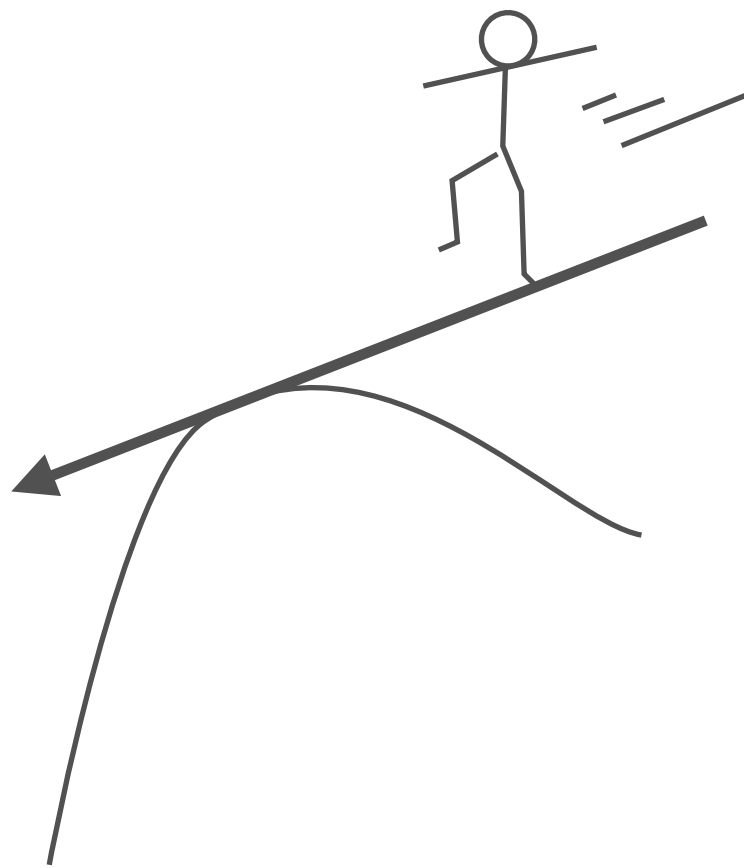
Helfer, Lim, Garcia & Amin (in prep)



super emitters + “sharp” transitions in gravitational wave output



end tangential digression



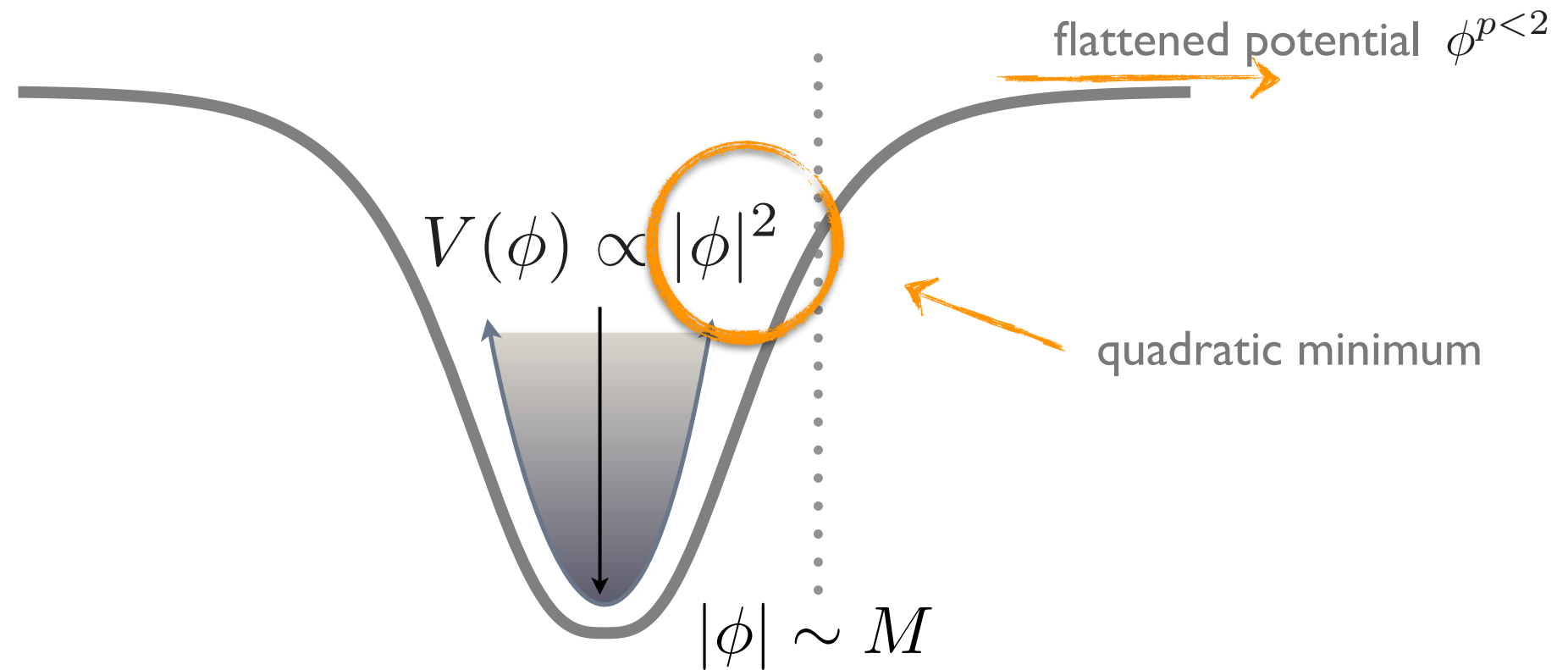
modeling end of inflation

SIMPLE

problem oriented

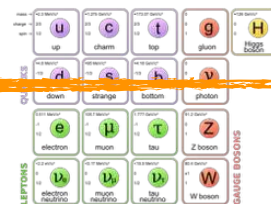
COMPLEX

end of inflation in “simple” models



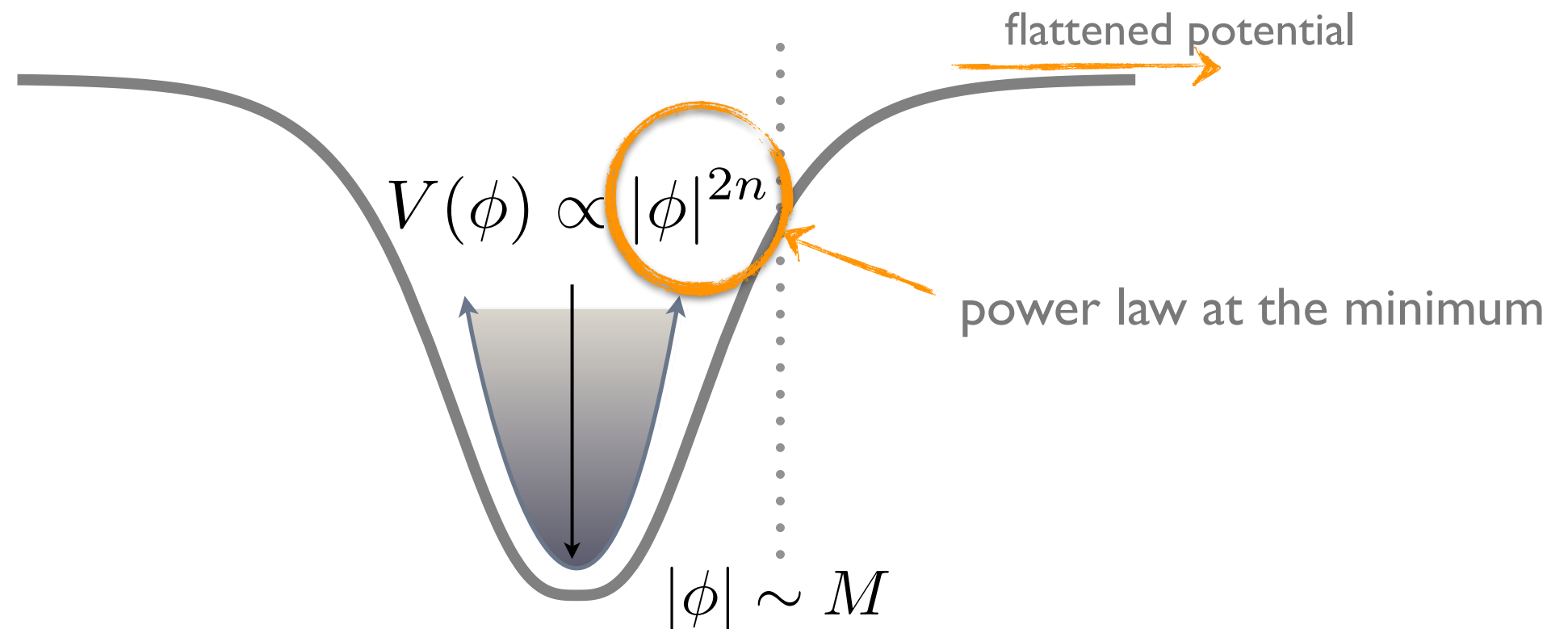
- shape of the potential (self couplings)

- ~~couplings to other fields~~



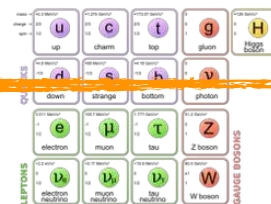
χ, ψ

end of inflation in “simple” models



- shape of the potential (self couplings)

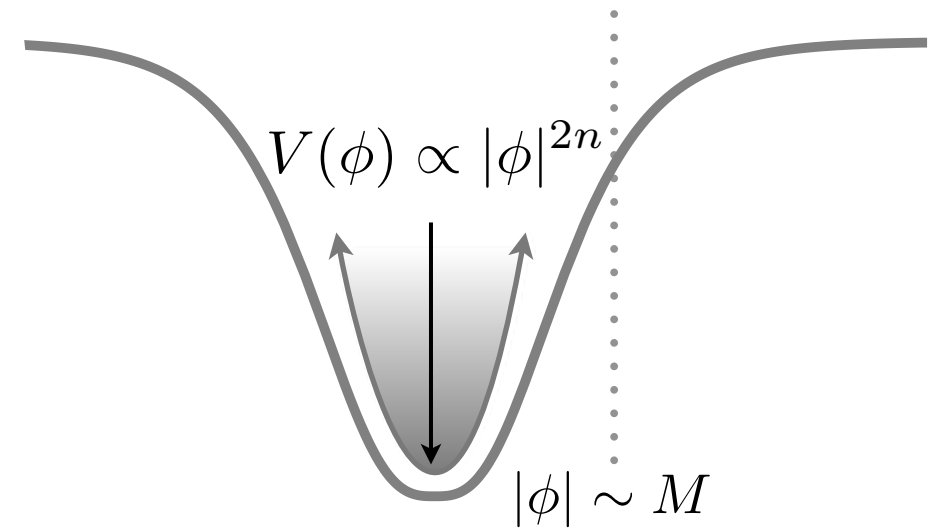
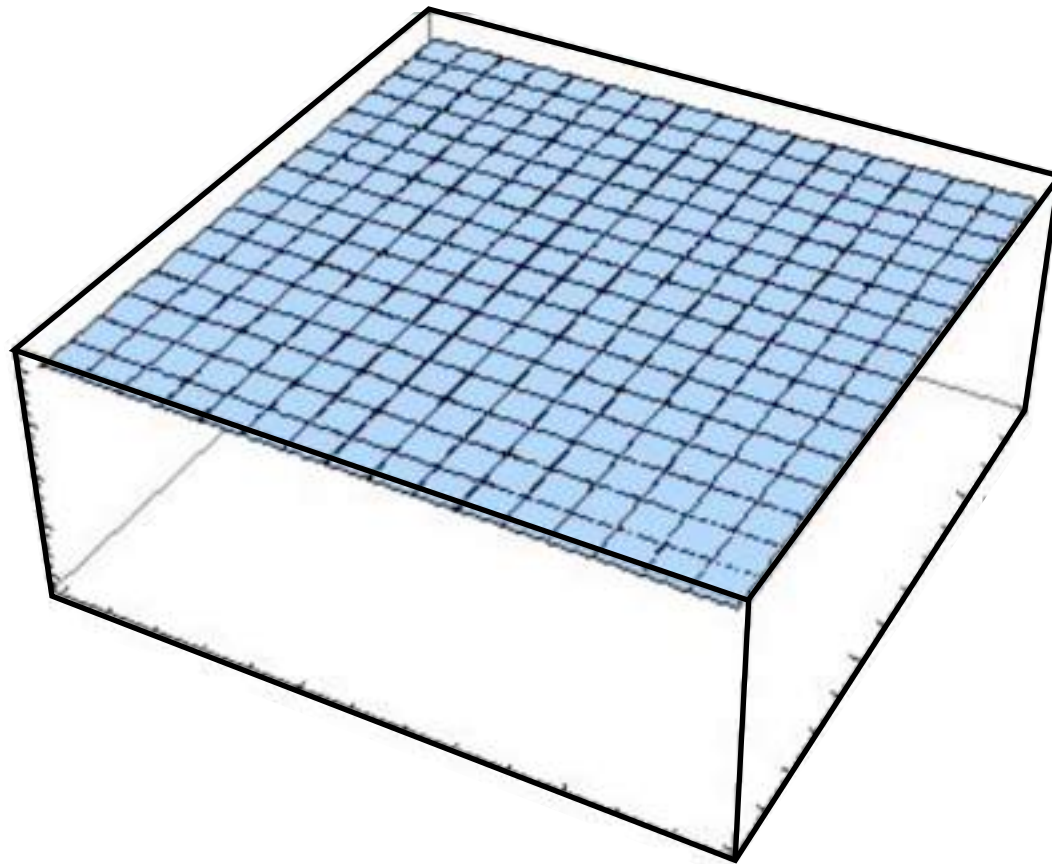
- ~~couplings to other fields~~



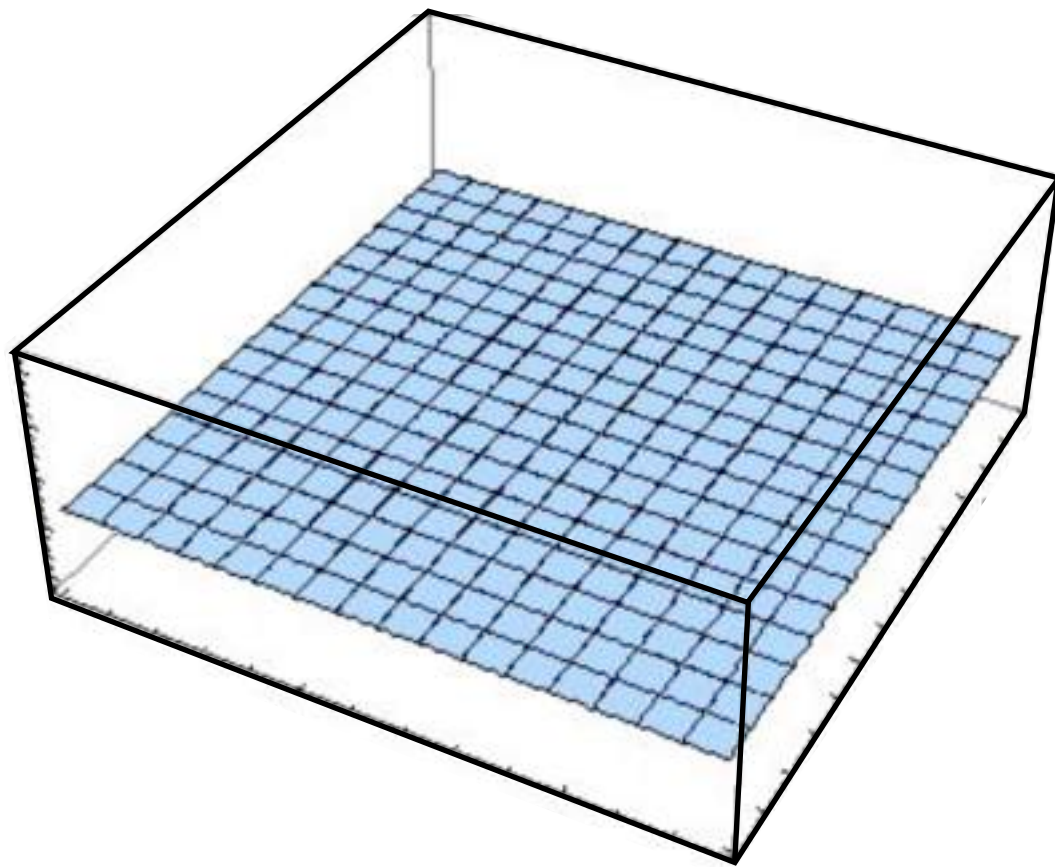
χ, ψ



homogeneous dynamics

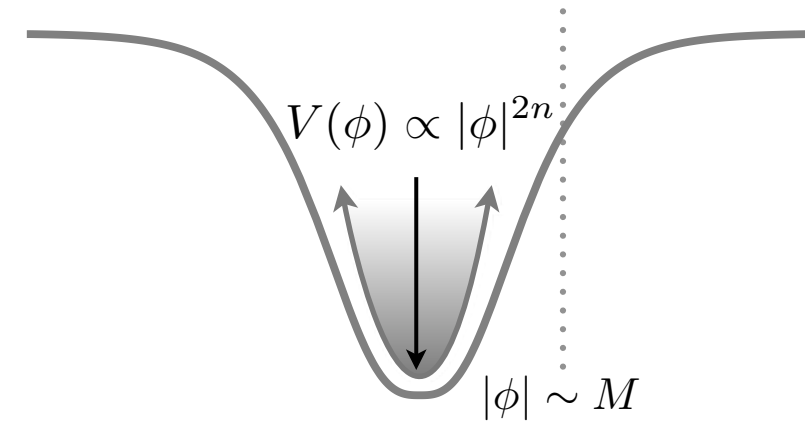


homogeneous eq. of state



eq. of state $w = \frac{\text{pressure}}{\text{density}}$

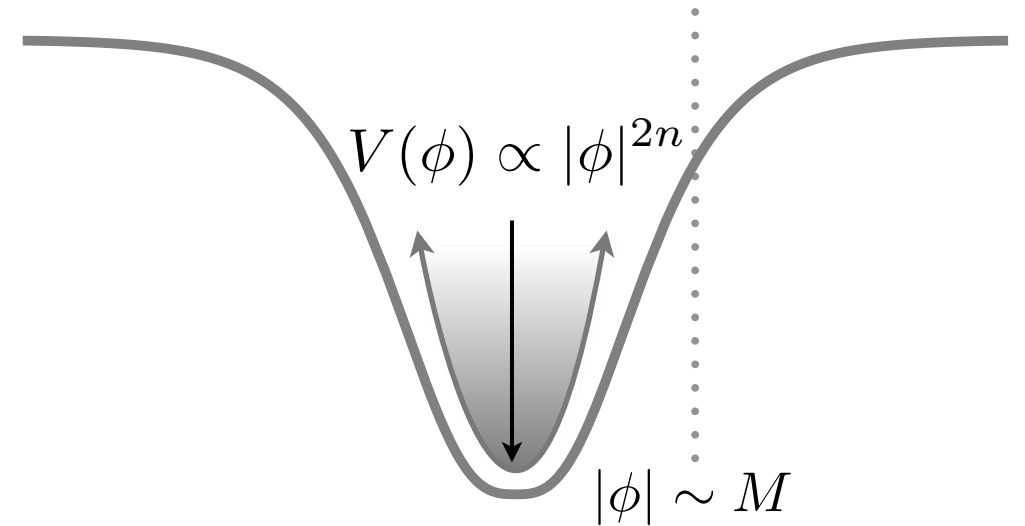
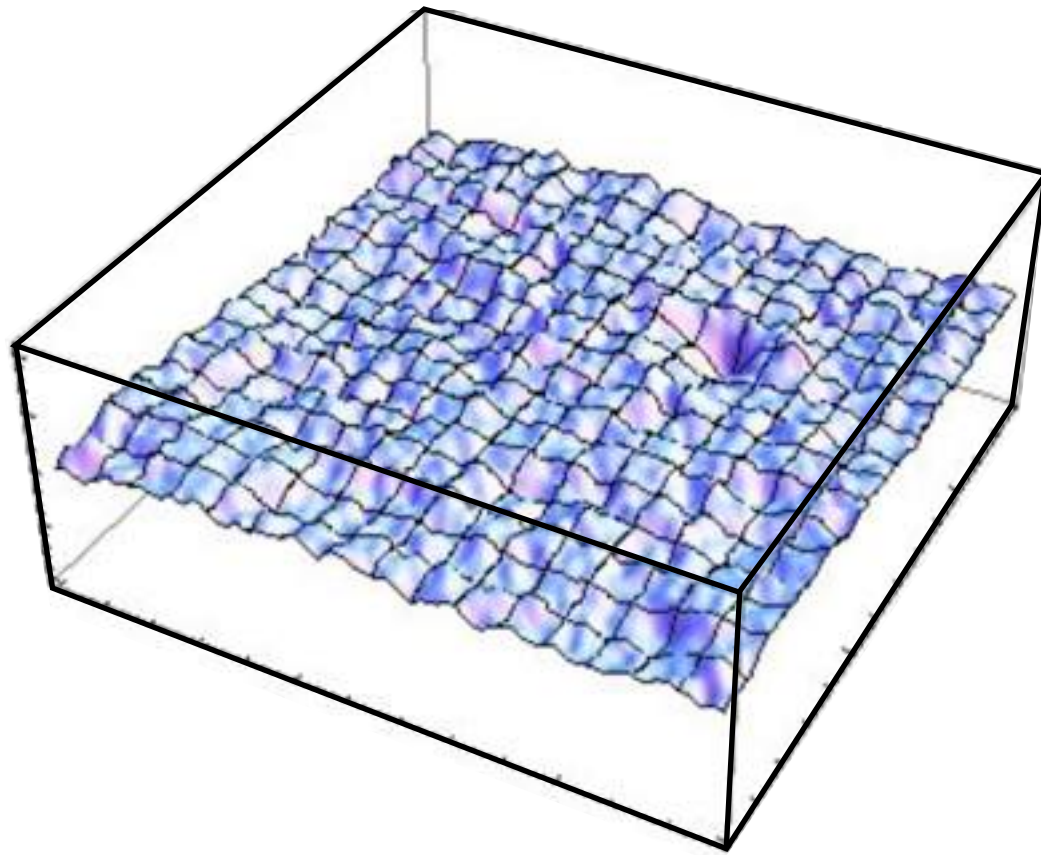
$$w \equiv \frac{\langle p \rangle_s}{\langle \rho \rangle_s} = \frac{\langle \dot{\phi}^2/2 - (\nabla \phi)^2/6a^2 - V \rangle_s}{\langle \dot{\phi}^2/2 + (\nabla \phi)^2/2a^2 + V \rangle_s} \approx \frac{n-1}{n+1}$$



Turner (1983)

* can be obtained from a viral theorem

fragmentation is (almost) inevitable



(i) existence of wings (self-couplings) $M \lesssim m_{\text{pl}}$

and/ or

(ii) non-quadratic minimum $n > 1$

* directly related to competition between growth rate and expansion, but duration depends on parameters

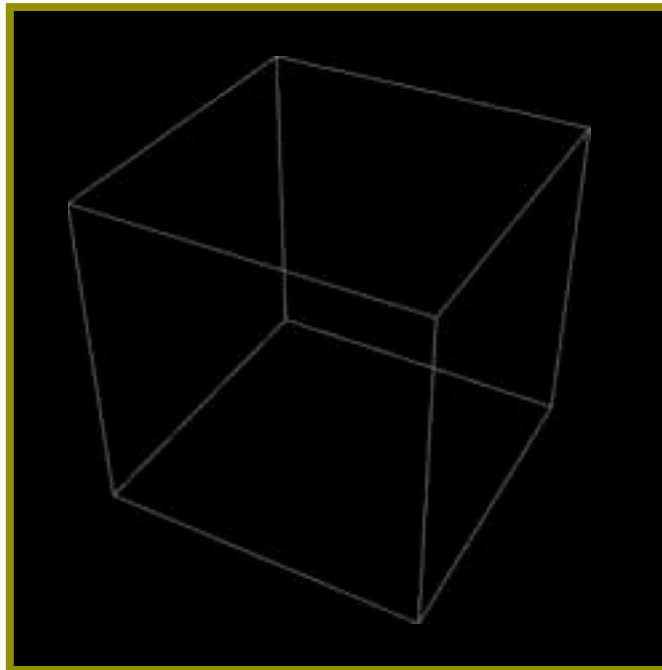
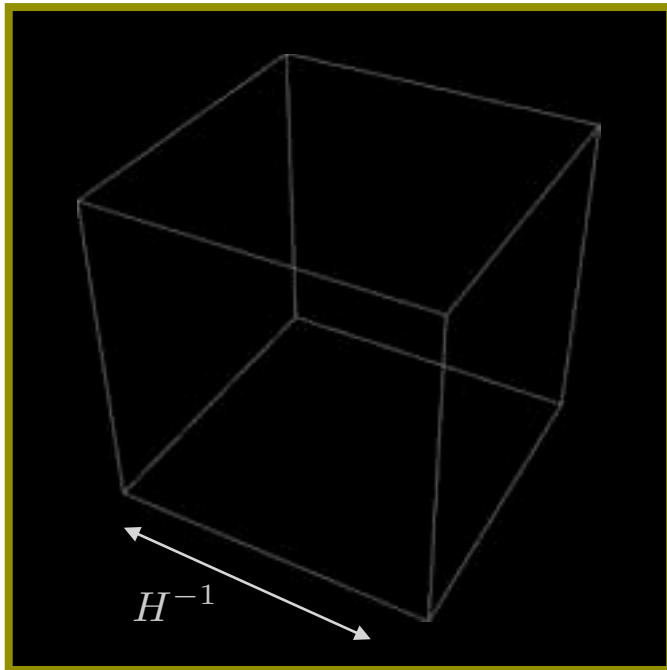
result of fragmented dynamics

* after sufficient time

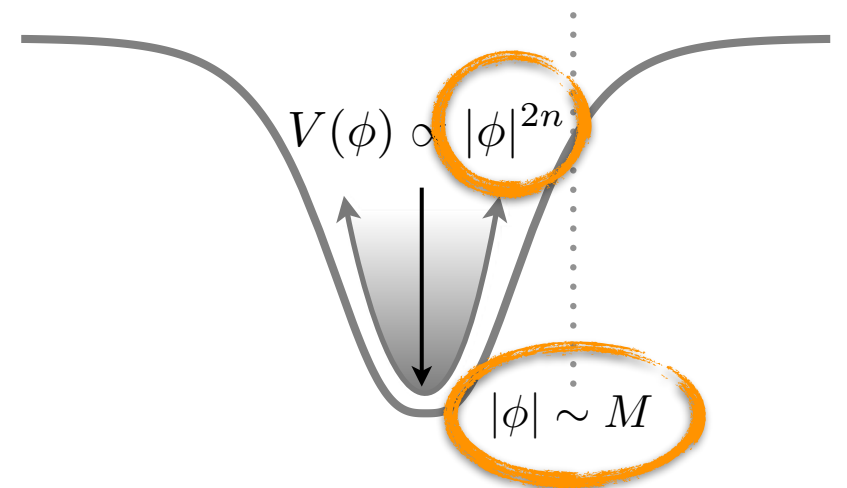
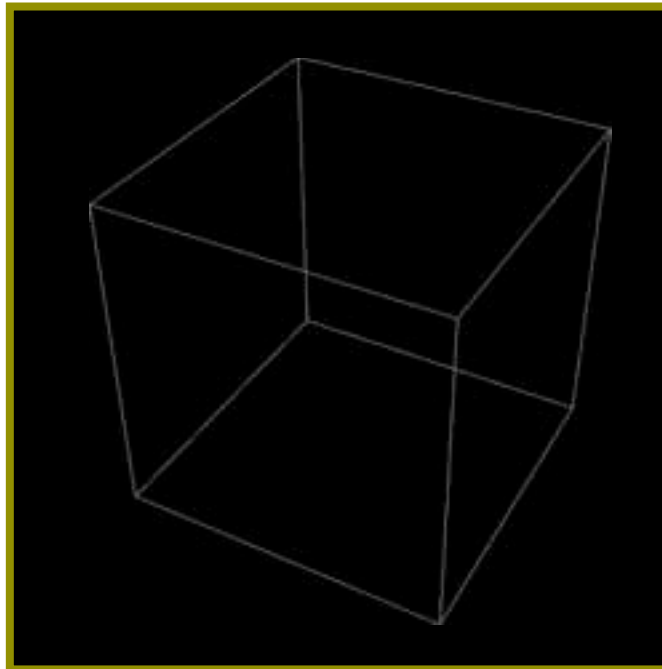
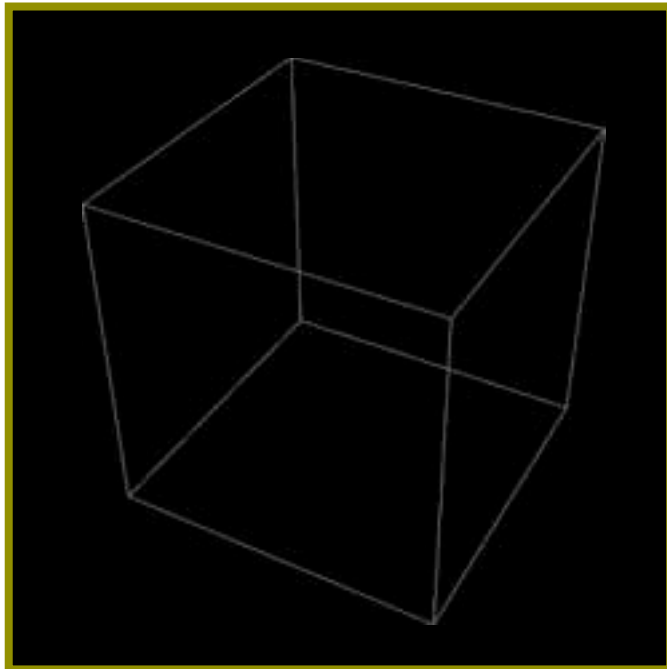
$n = 1$

$n > 1$

$M \sim m_{\text{pl}}$



$M \ll m_{\text{pl}}$



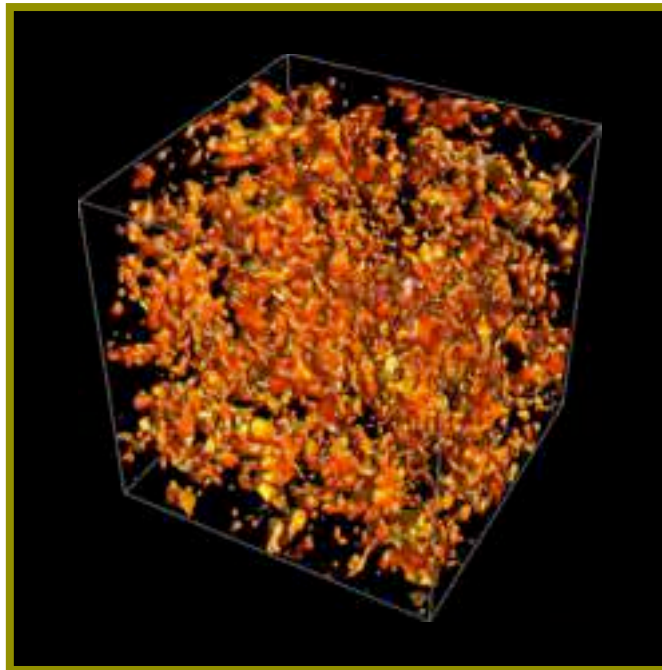
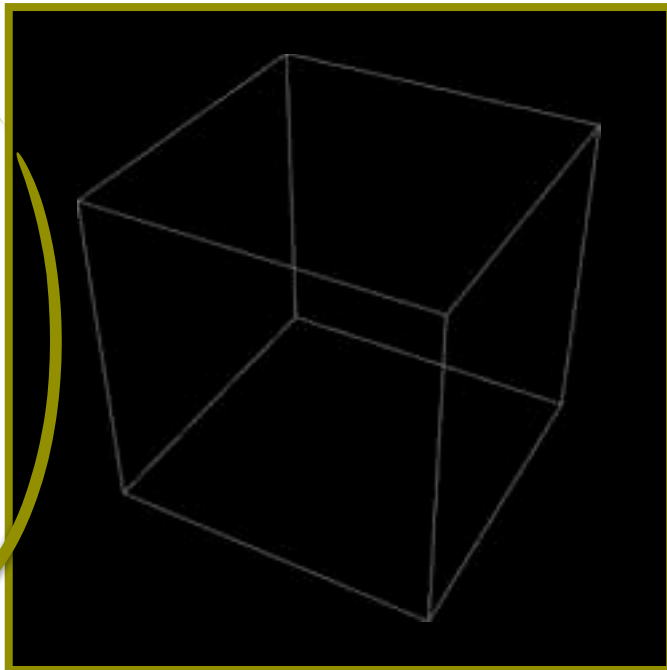
result of fragmented dynamics

* after sufficient time

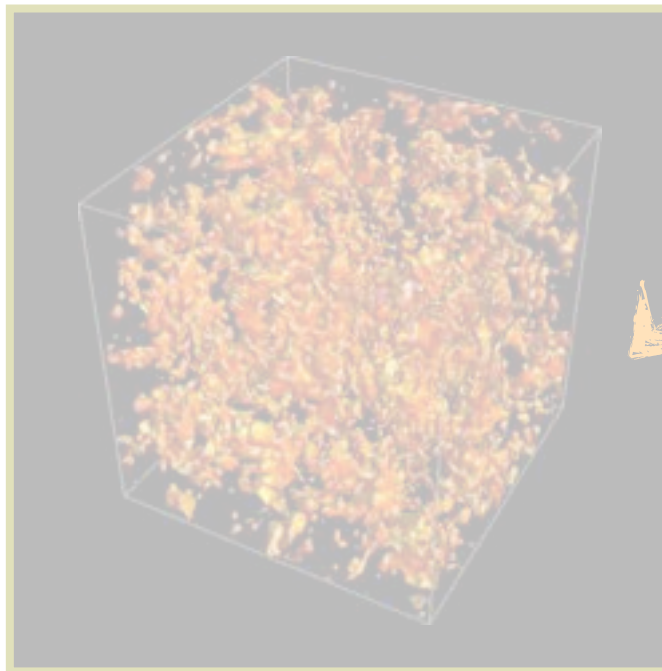
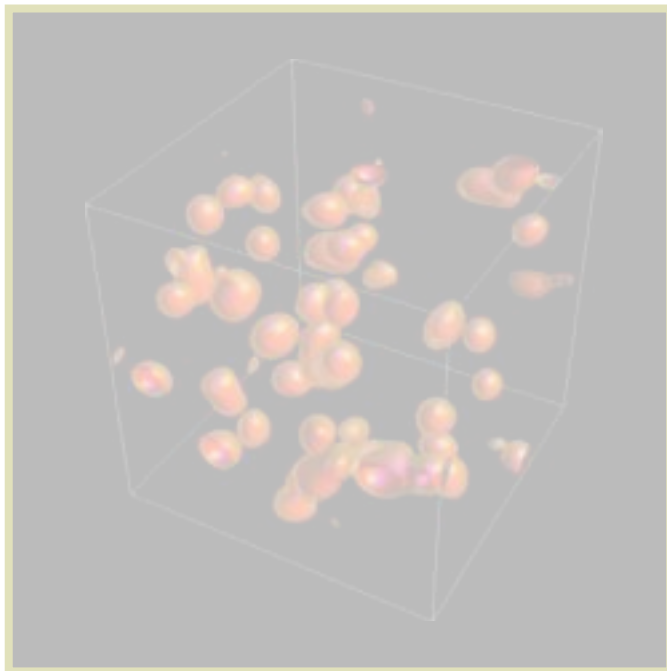
$n = 1$

$n > 1$

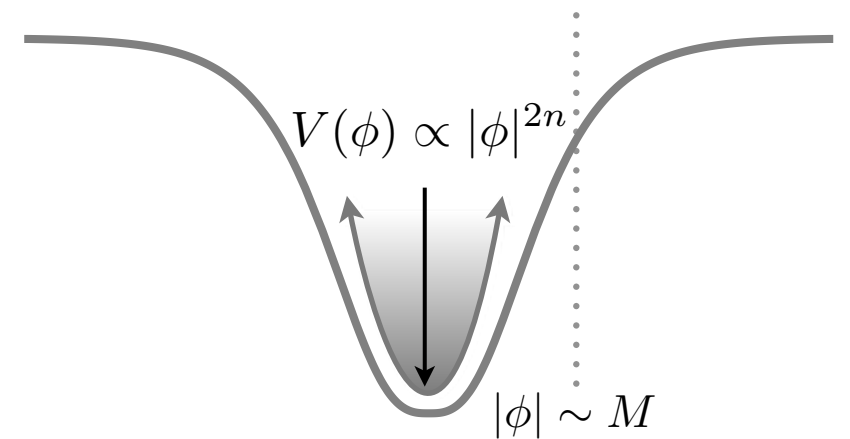
$M \sim m_{\text{pl}}$



$M \ll m_{\text{pl}}$



slow



quickly

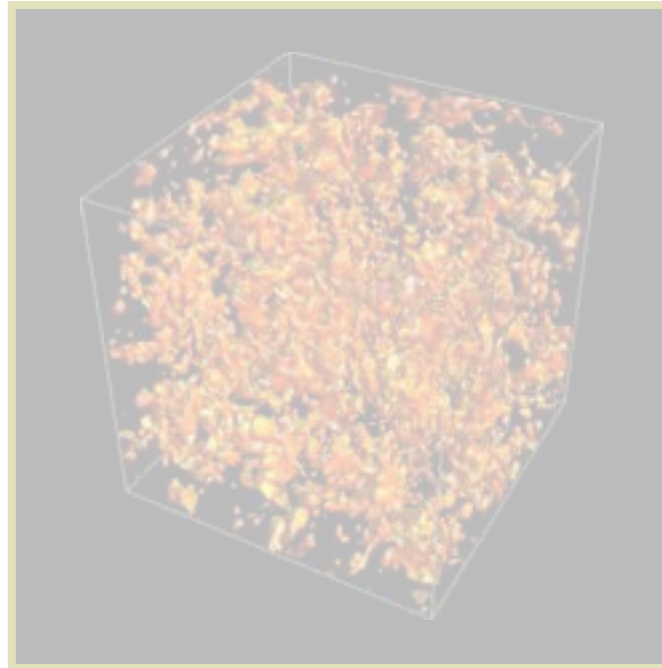
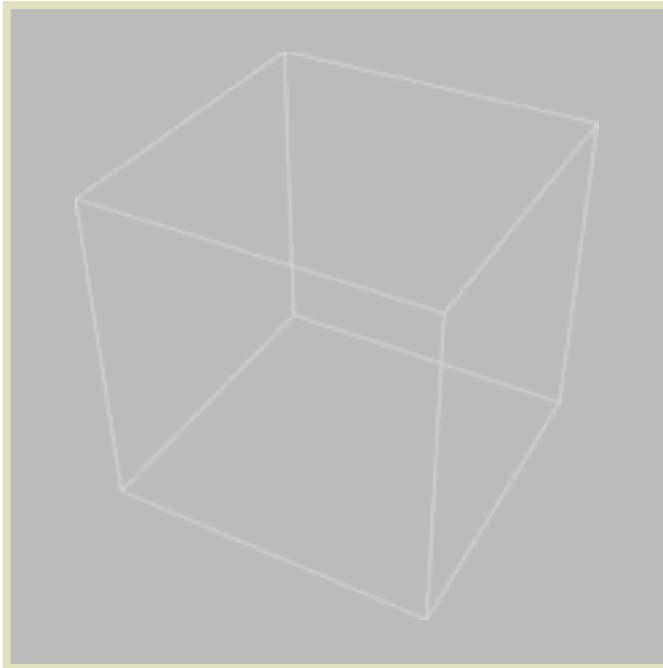
result of fragmented dynamics

* after sufficient time

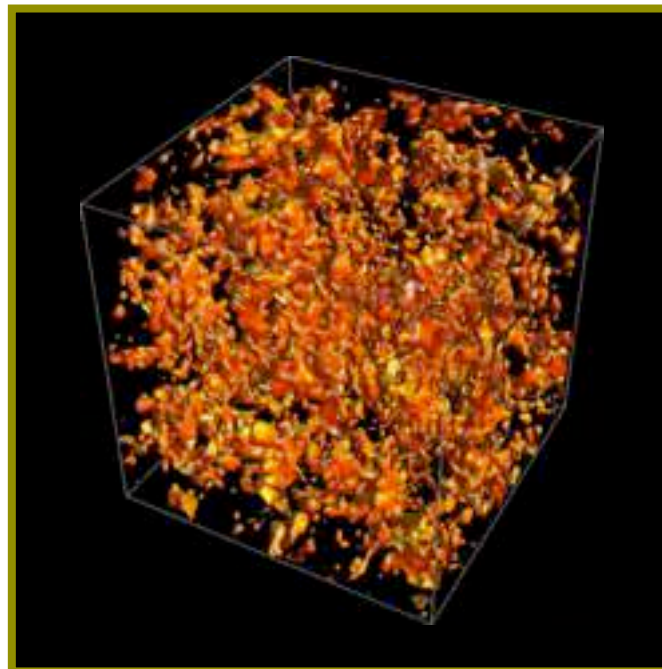
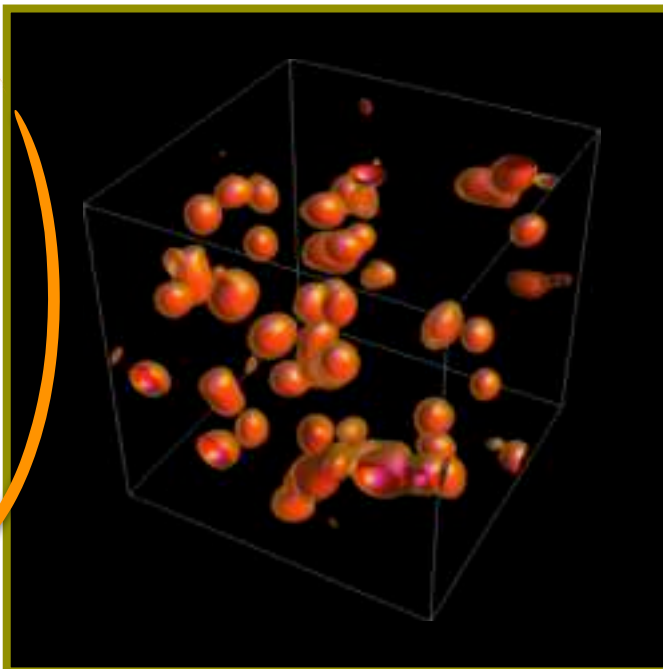
$n = 1$

$n > 1$

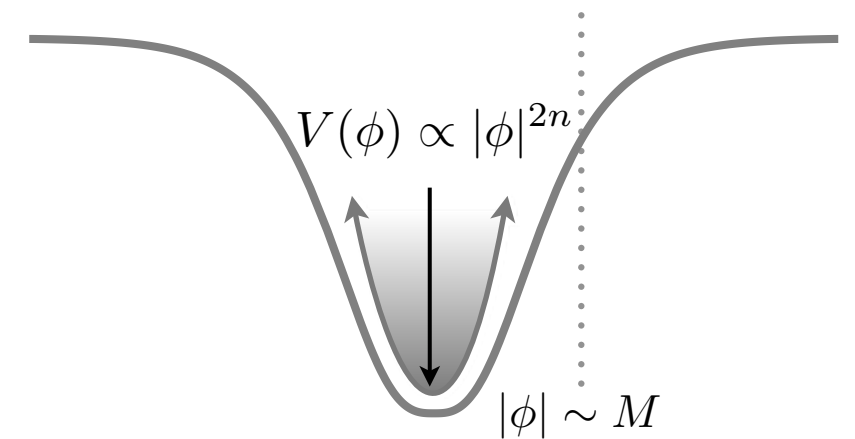
$M \sim m_{\text{pl}}$



$M \ll m_{\text{pl}}$



slow



fast

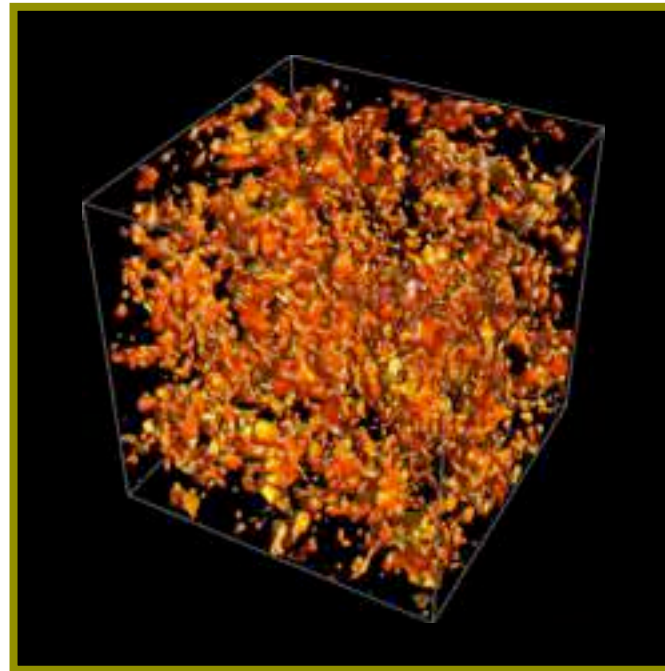
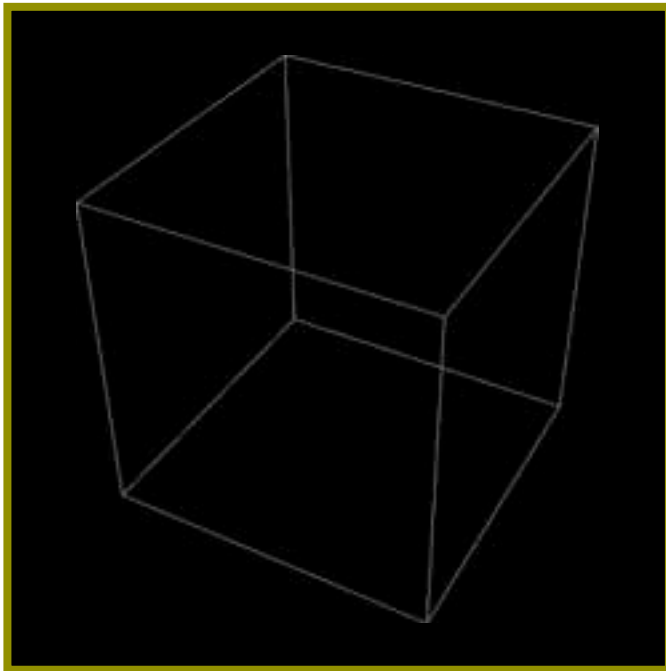
eq. of state

* after sufficient time

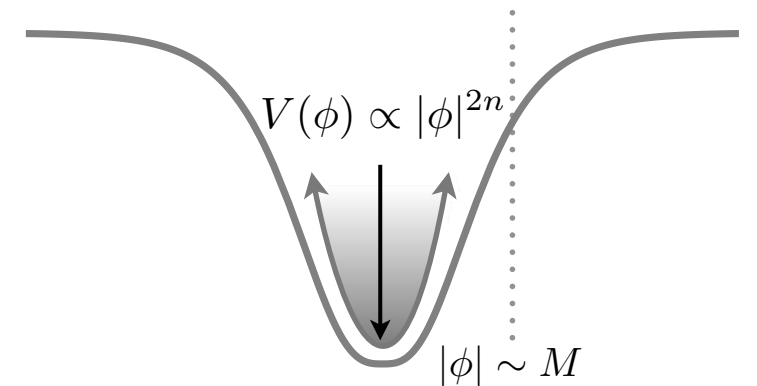
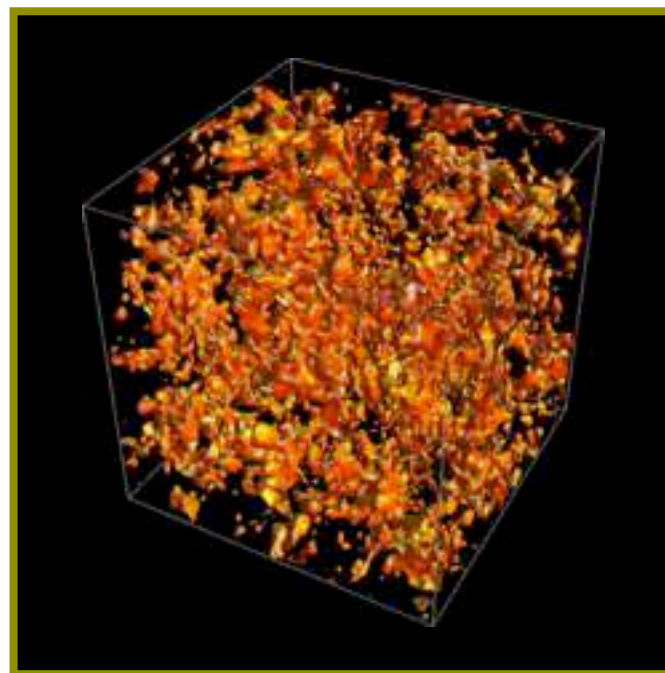
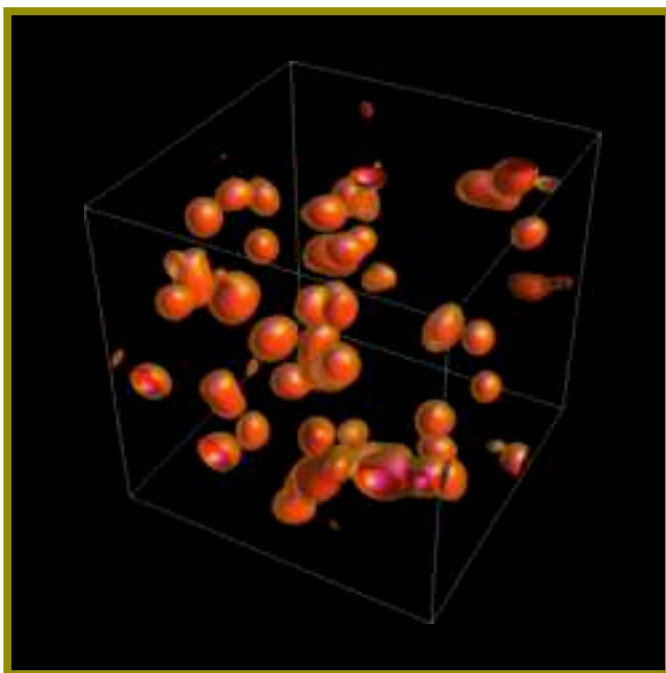
$n = 1$

$n > 1$

$M \sim m_{\text{pl}}$



$M \ll m_{\text{pl}}$



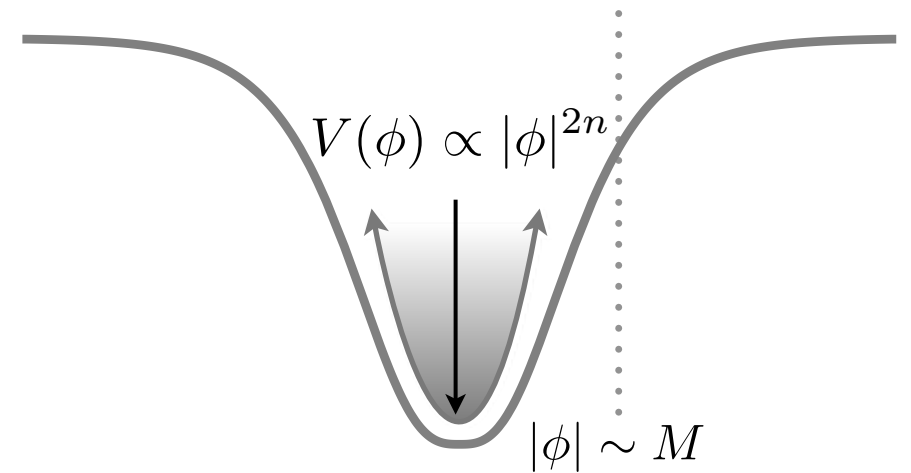
$$w \rightarrow \begin{cases} 0 & \text{if } n = 1 \\ 1/3 & \text{if } n > 1 \end{cases}$$

independent of M

$$w \neq \frac{n-1}{n+1}$$

an upper bound on duration to radiation domination $n > 1$

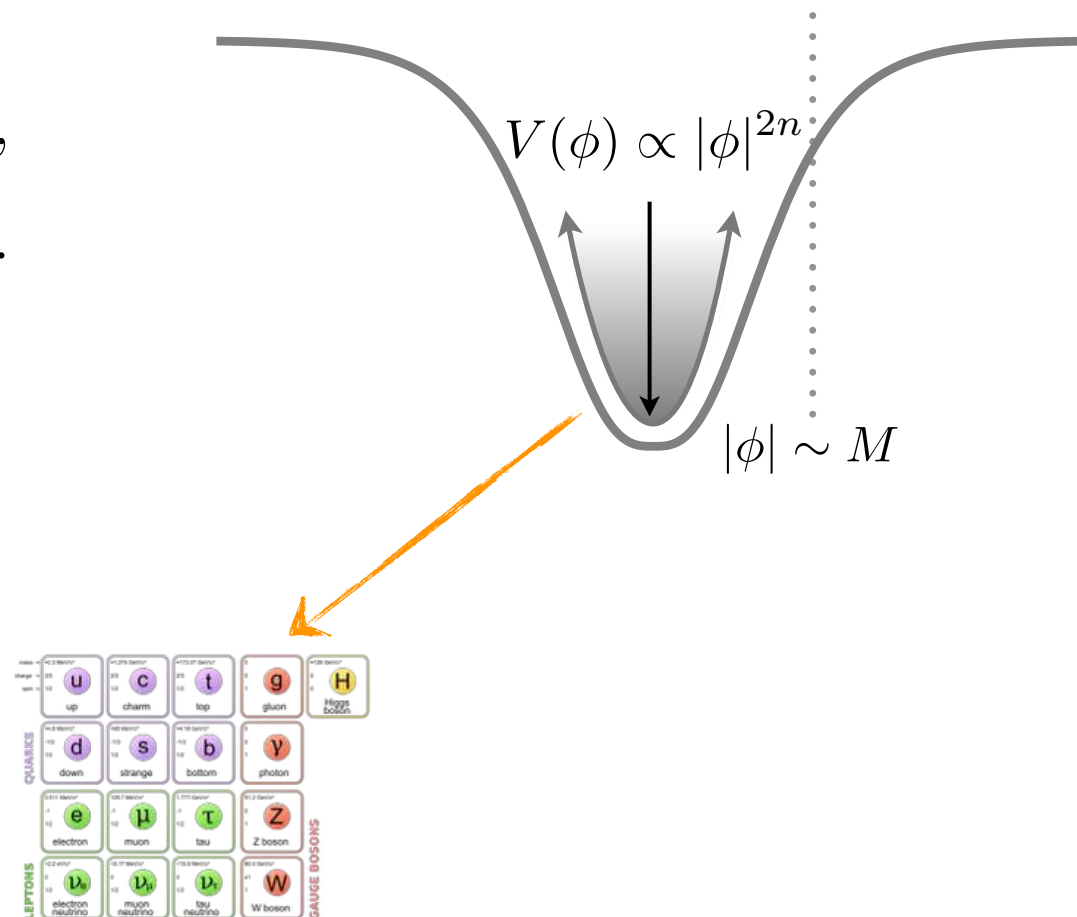
$$\Delta N_{\text{rad}} \sim \begin{cases} 1 & M \lesssim 10^{-2} m_{\text{Pl}} , \\ \frac{n+1}{3} \ln \left(\frac{\kappa}{\Delta\kappa} \frac{10M}{m_{\text{Pl}}} \right) & M \gtrsim 10^{-2} m_{\text{Pl}} . \end{cases}$$



an upper bound on duration to radiation domination

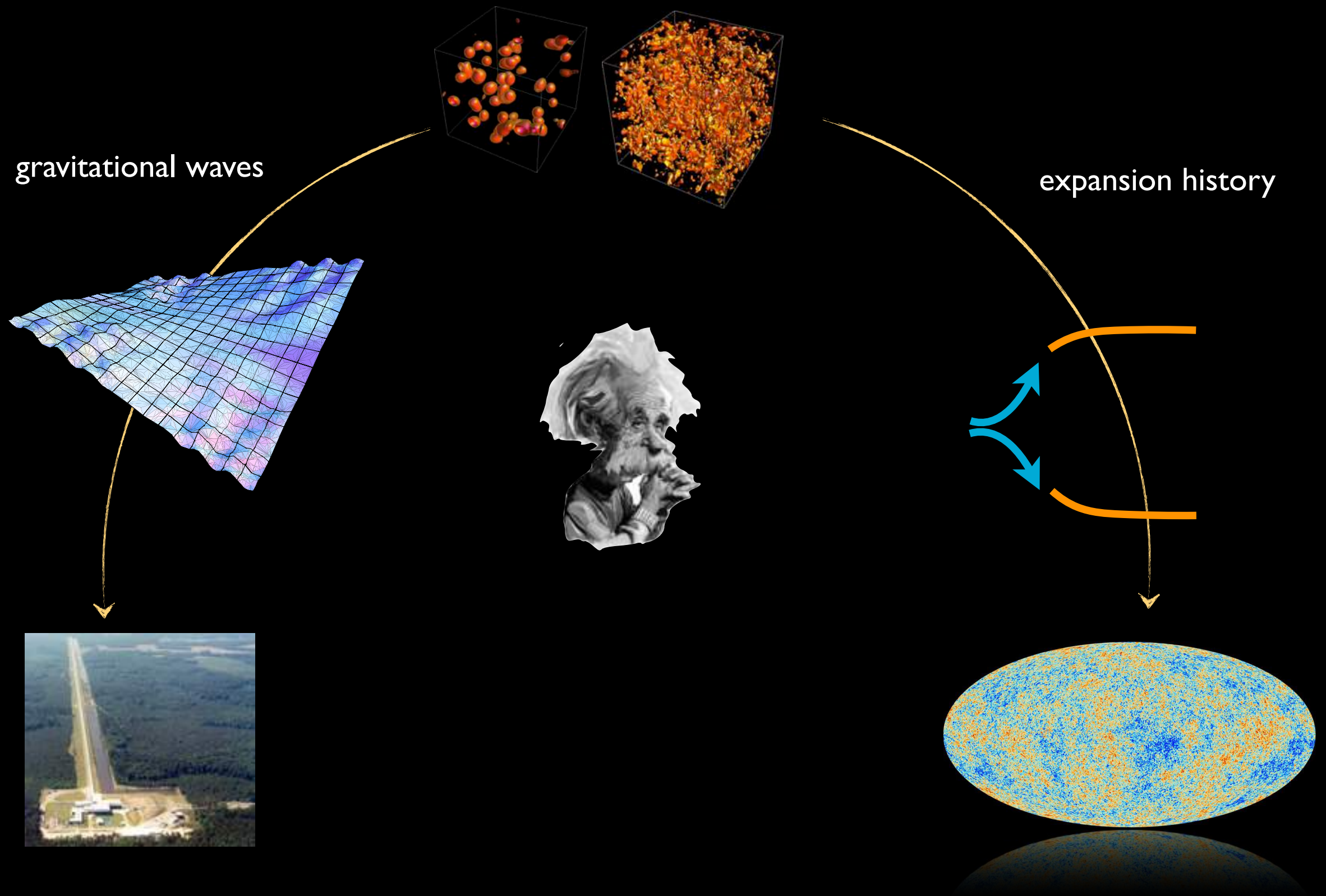
$$\Delta N_{\text{rad}} \sim \begin{cases} 1 & M \lesssim 10^{-2} m_{\text{Pl}} , \\ \frac{n+1}{3} \ln \left(\frac{\kappa}{\Delta\kappa} \frac{10M}{m_{\text{Pl}}} \right) & M \gtrsim 10^{-2} m_{\text{Pl}} . \end{cases}$$

additional *light (massless) fields* can
only decrease the duration!

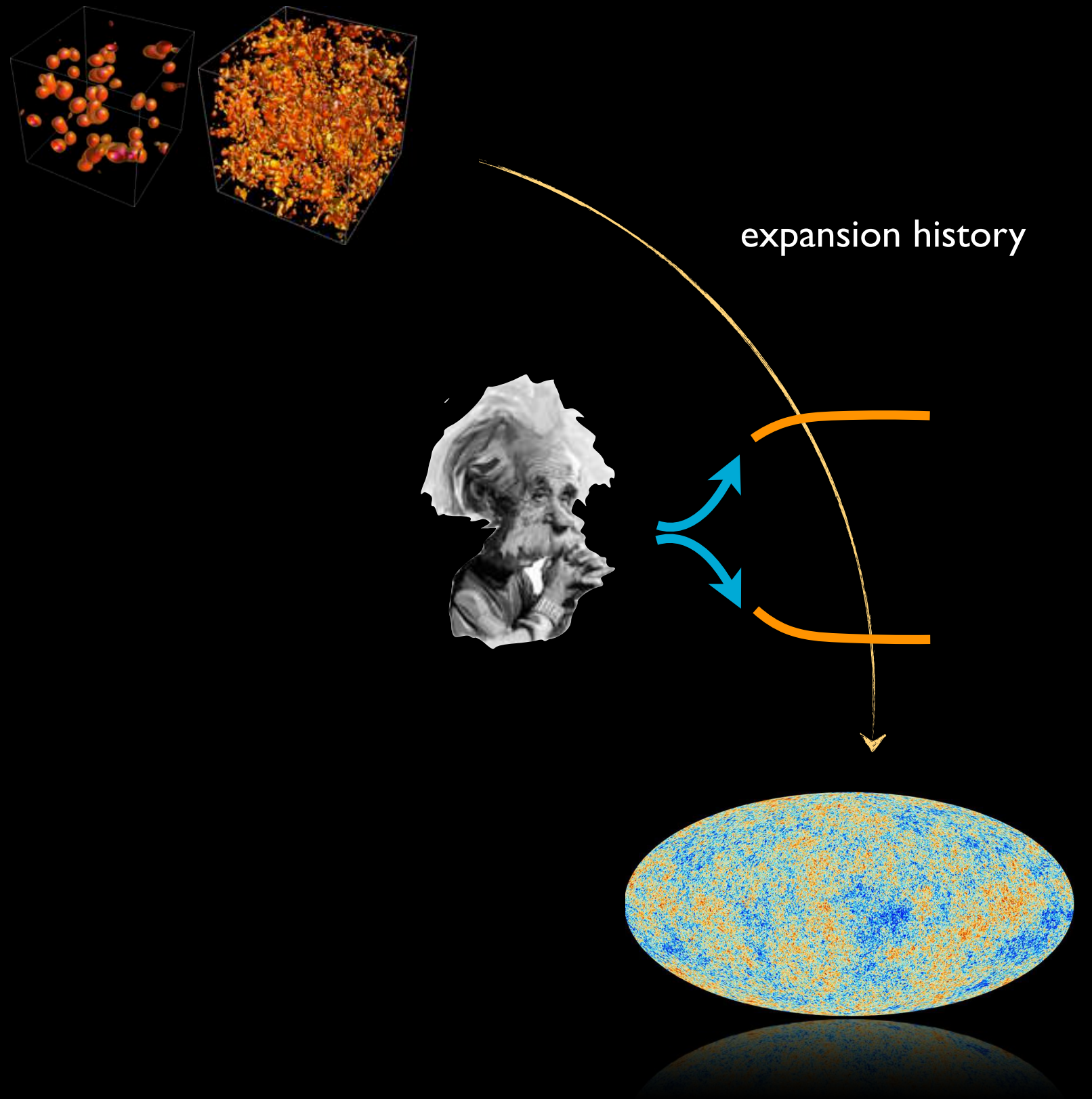


* decay to significantly massive fields can change this conclusion

consequences ?



expansion history



expansion history — eq. of state

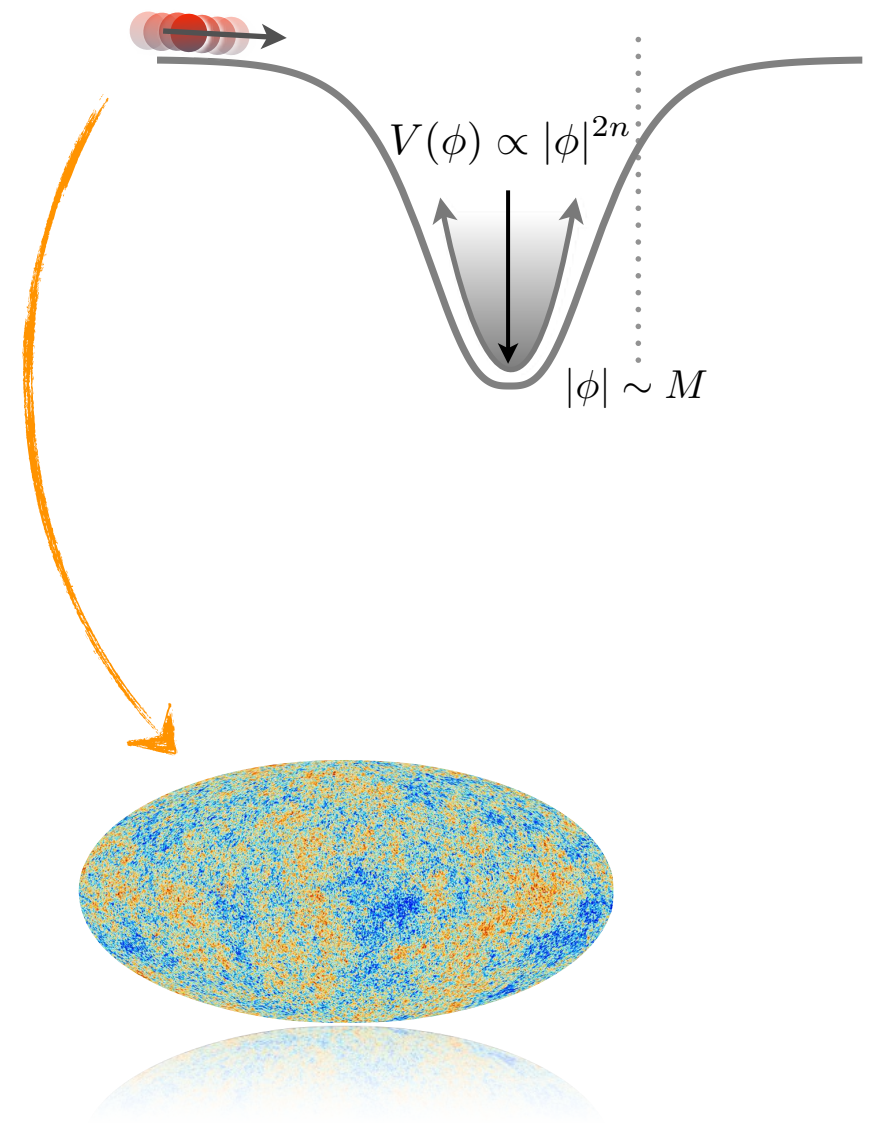
Einstein's Equations

$$G_{\mu\nu} = 8\pi G T_{\mu\nu}$$

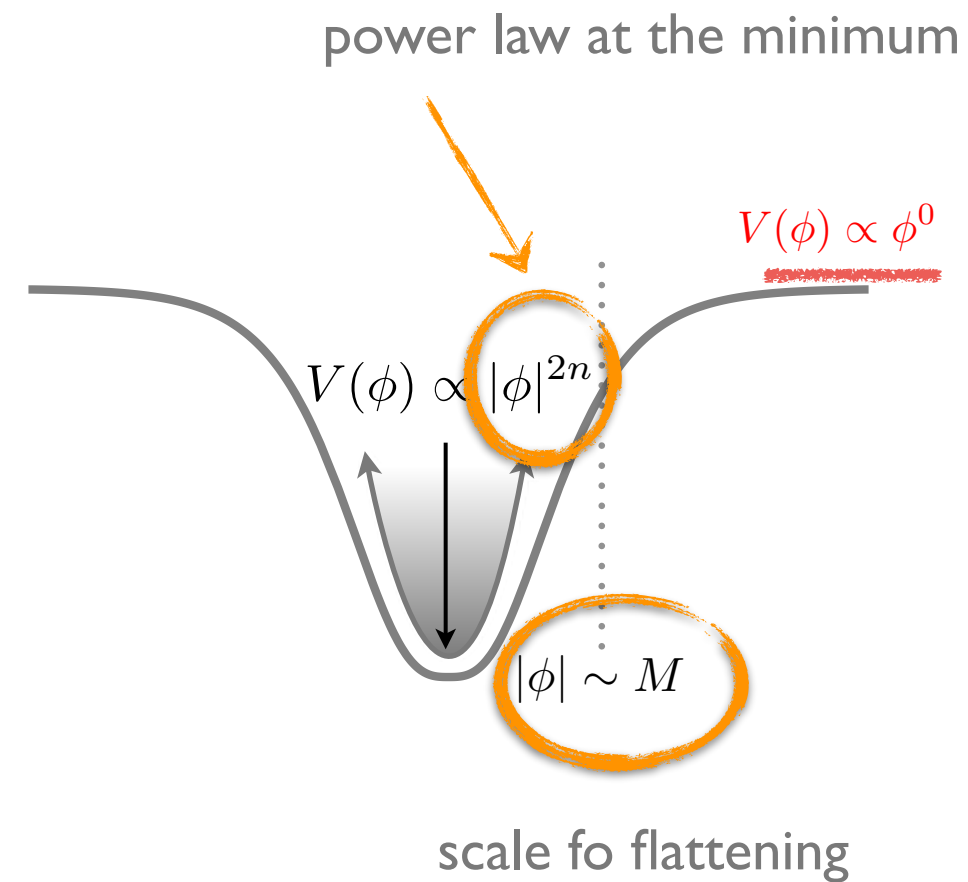
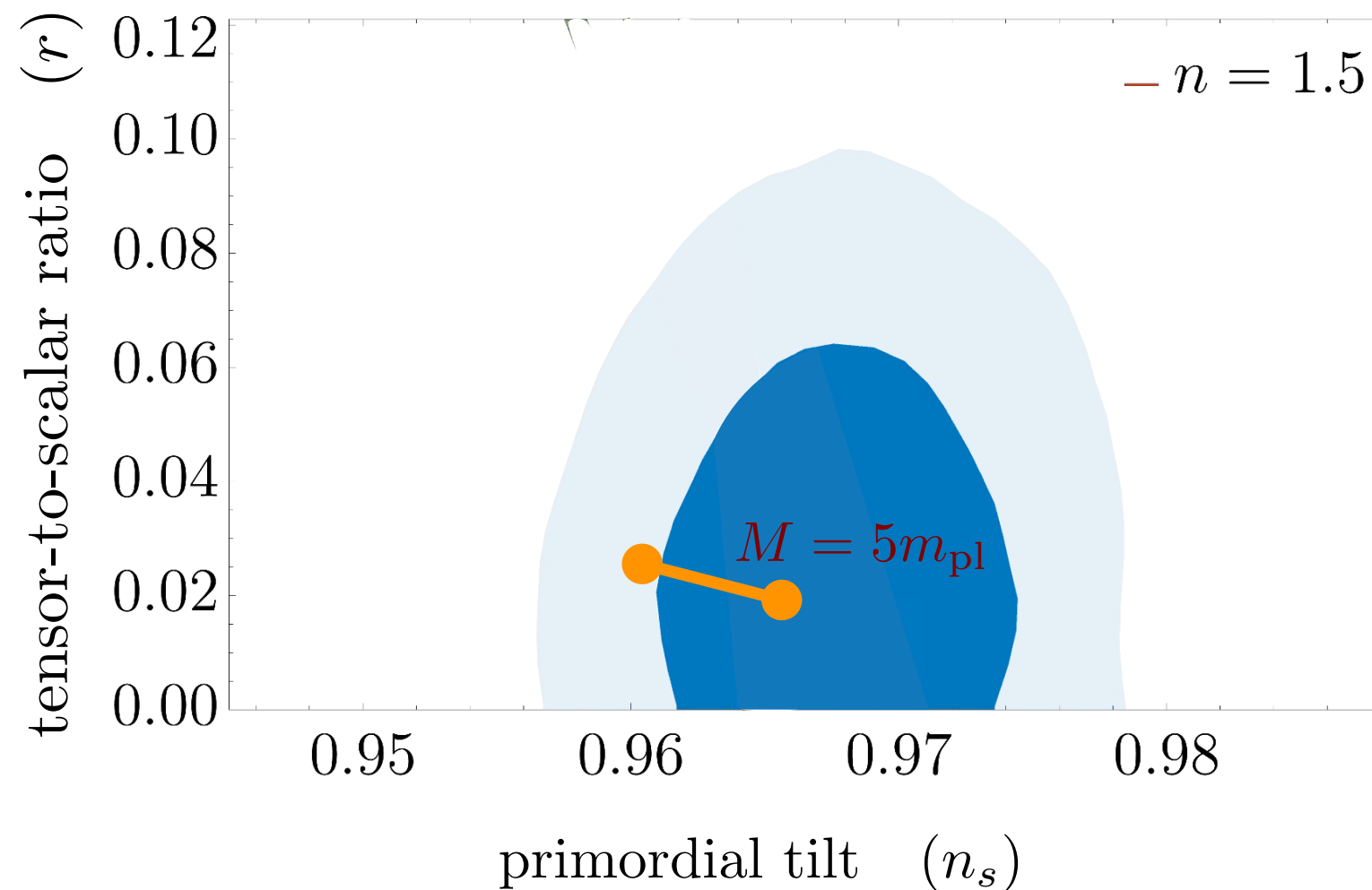
$$a(t) \propto t^{\frac{2}{3(1+w)}}$$

$$\text{eq. of state } w = \frac{\text{pressure}}{\text{density}}$$

impacts
mapping of scales

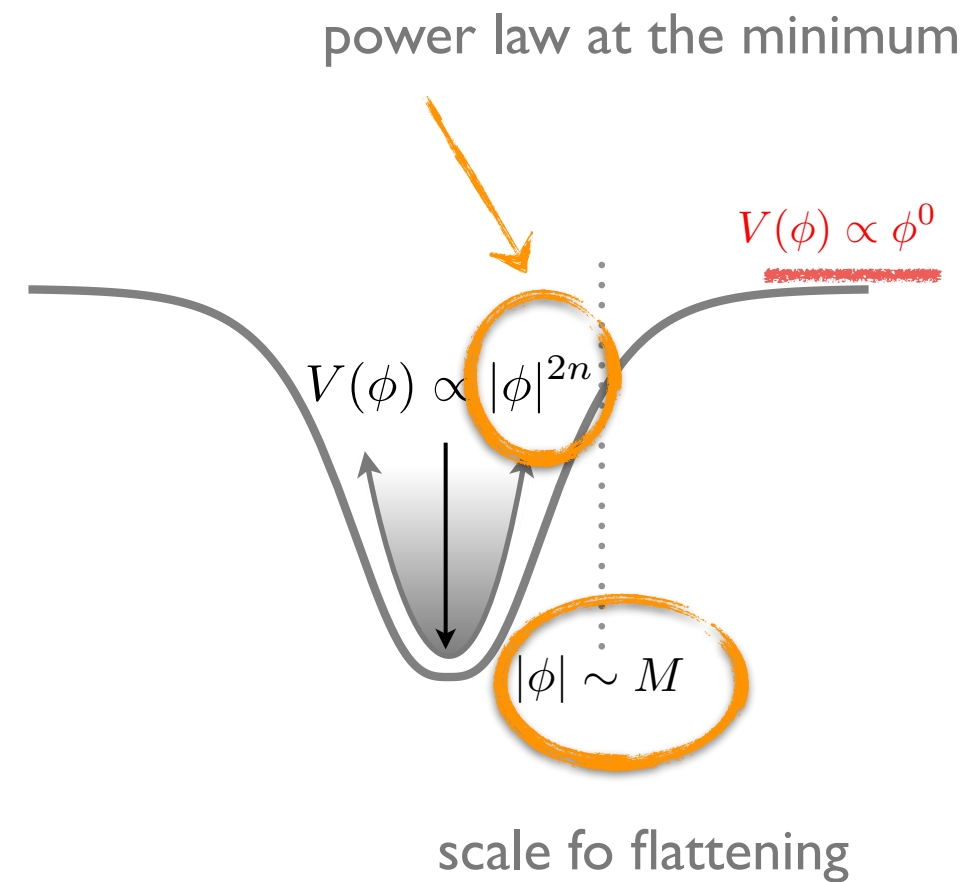
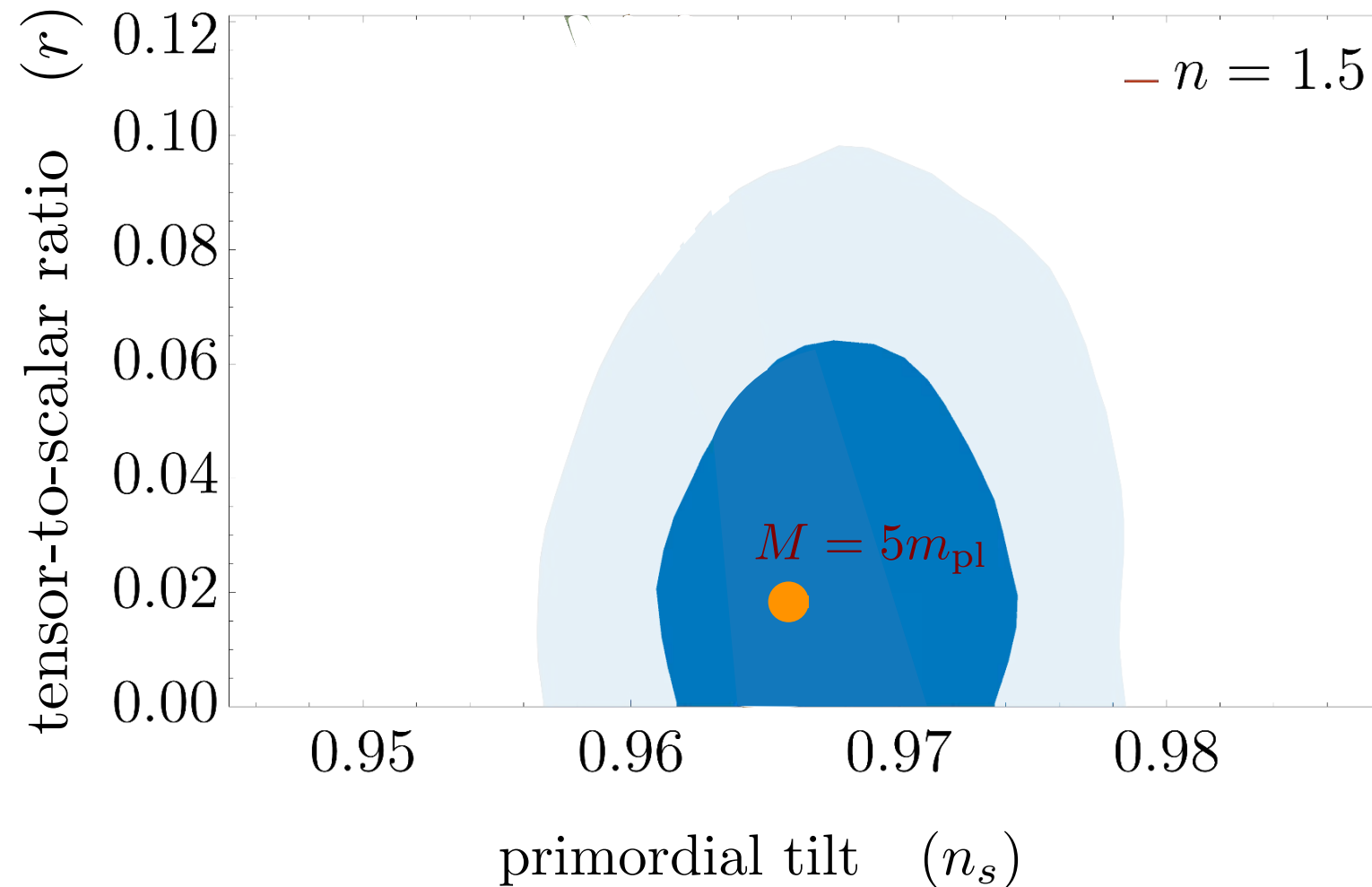


implications for CMB observables



* non-quadratic minimum

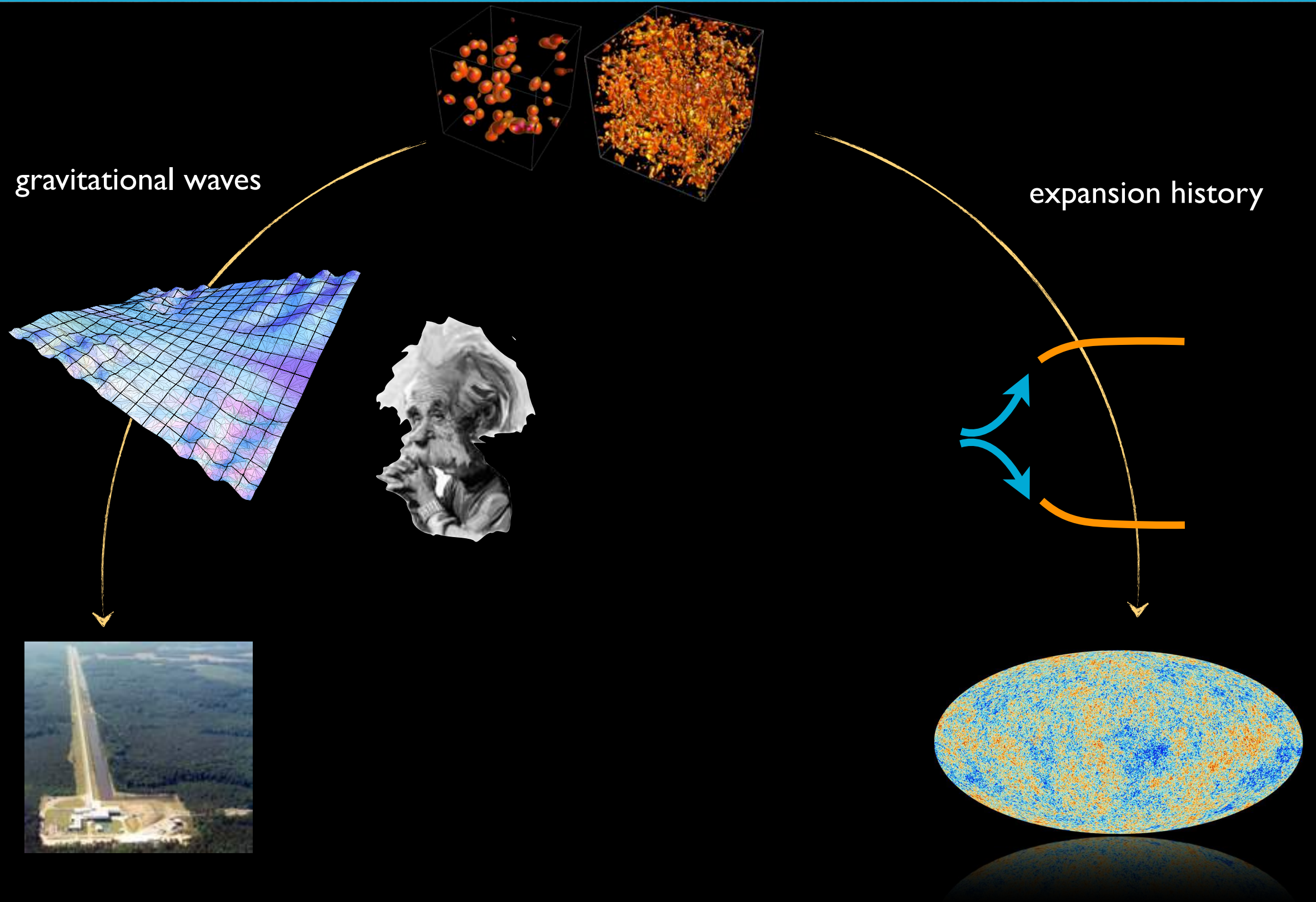
reduction in uncertainty!



$$\Delta N_{\text{rad}} \sim \begin{cases} 1 & M \lesssim 10^{-2} m_{\text{Pl}} , \\ \frac{n+1}{3} \ln \left(\frac{\kappa}{\Delta \kappa} \frac{10M}{m_{\text{Pl}}} \right) & M \gtrsim 10^{-2} m_{\text{Pl}} . \end{cases}$$

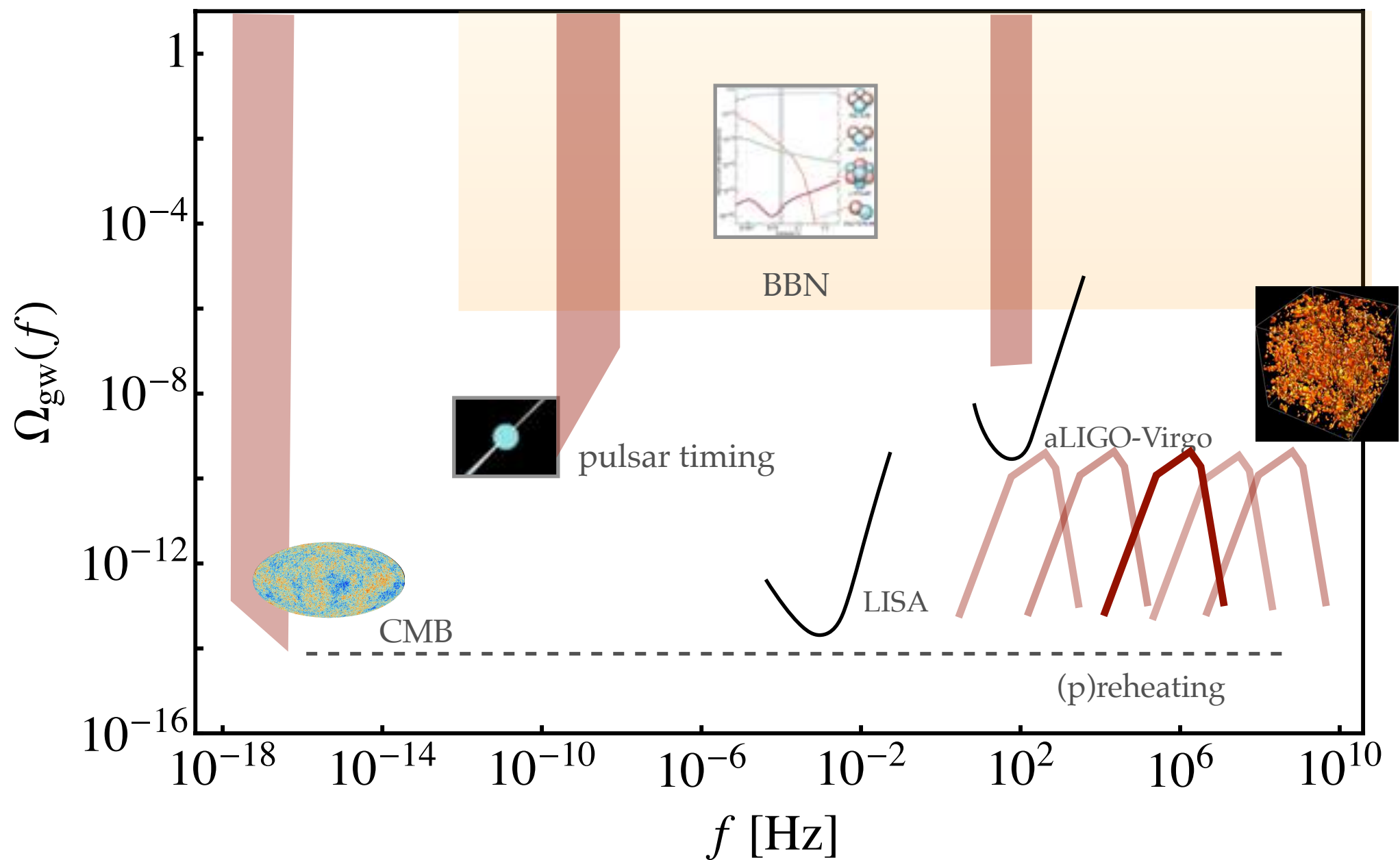
* non-quadratic minimum

consequences ?



stochastic gravitational waves

$$\Omega_{\text{gw}}(f) = \frac{d \ln \rho_{\text{gw}}}{d \ln f} \sim \frac{\rho_{\text{gw}}}{\rho_{\text{crit}}}$$



end of inflation



SIMPLE

- single field
- non-trivial dynamics
- eq. of state + gravitational waves

detailed models

COMPLEX

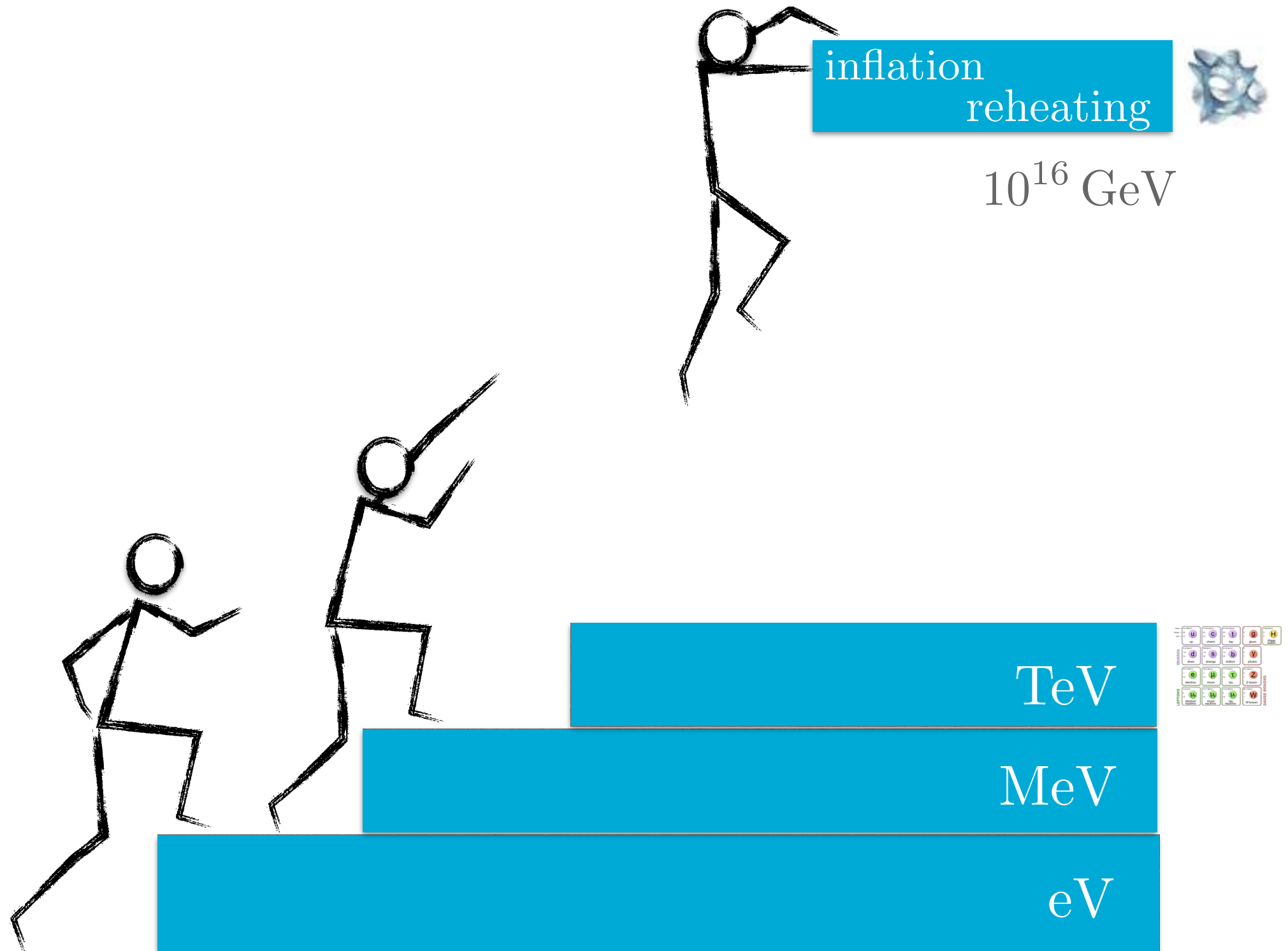
end of inflation

SIMPLE

- single field
- non-trivial dynamics
- eq. of state + gravitational waves

detailed models

COMPLEX



theory : its complicated (probably)

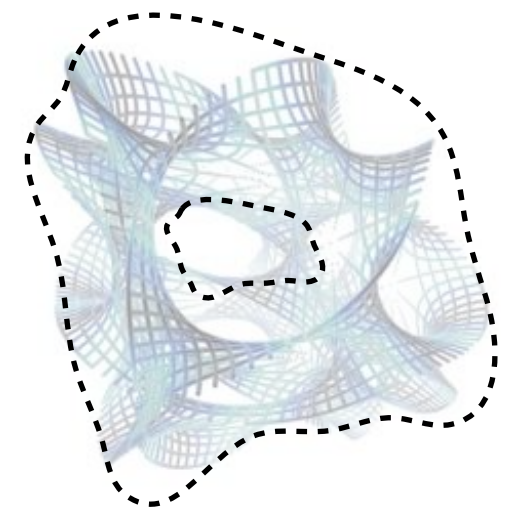
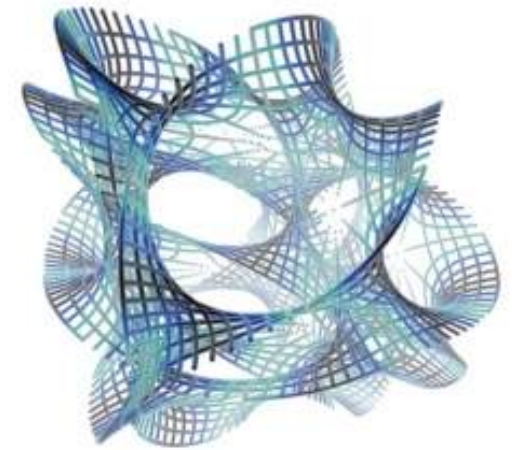
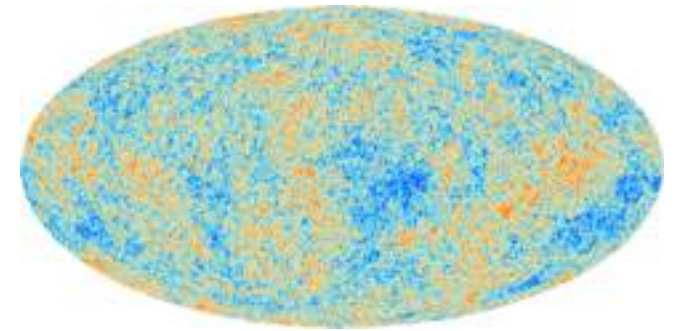
- inflation
- reheating after inflation



QUARKS		GAUGE BOSONS	
UP-QUARKS	DOWN-QUARKS	PHOTONS	GLUONS
<p>mass: 2.3 MeV/c²</p> <p>charge: 2/3</p> <p>spin: 1/2</p> <p>u</p> <p>up</p>	<p>mass: 4.1873 MeV/c²</p> <p>charge: 1/3</p> <p>spin: 1/2</p> <p>c</p> <p>charm</p>	<p>mass: 0</p> <p>charge: 0</p> <p>spin: 1</p> <p>t</p> <p>top</p>	<p>mass: 0</p> <p>charge: 0</p> <p>spin: 1</p> <p>g</p> <p>gluon</p>
<p>mass: 4.1873 MeV/c²</p> <p>charge: 1/3</p> <p>spin: 1/2</p> <p>b</p> <p>bottom</p>	<p>mass: 4.1873 MeV/c²</p> <p>charge: 1/3</p> <p>spin: 1/2</p> <p>s</p> <p>strange</p>	<p>mass: 0</p> <p>charge: 0</p> <p>spin: 1</p> <p>h</p> <p>Higgs boson</p>	<p>mass: 0</p> <p>charge: 0</p> <p>spin: 1</p> <p>γ</p> <p>photon</p>
<p>mass: 0.511 MeV/c²</p> <p>charge: -1</p> <p>spin: 1/2</p> <p>e</p> <p>electron</p>	<p>mass: 105.658 MeV/c²</p> <p>charge: -1</p> <p>spin: 1/2</p> <p>μ</p> <p>muon</p>	<p>mass: 1.777 MeV/c²</p> <p>charge: -1</p> <p>spin: 1/2</p> <p>τ</p> <p>tau</p>	<p>mass: 91.1876 MeV/c²</p> <p>charge: 0</p> <p>spin: 1</p> <p>Z</p> <p>Z boson</p>
<p>mass: 0.511 MeV/c²</p> <p>charge: 0</p> <p>spin: 1/2</p> <p>ν_e</p> <p>electron neutrino</p>	<p>mass: 0.105658 MeV/c²</p> <p>charge: 0</p> <p>spin: 1/2</p> <p>ν_μ</p> <p>muon neutrino</p>	<p>mass: 1.777 MeV/c²</p> <p>charge: 0</p> <p>spin: 1/2</p> <p>ν_τ</p> <p>tau neutrino</p>	<p>mass: 80.379 MeV/c²</p> <p>charge: +1</p> <p>spin: 1</p> <p>W</p> <p>W boson</p>

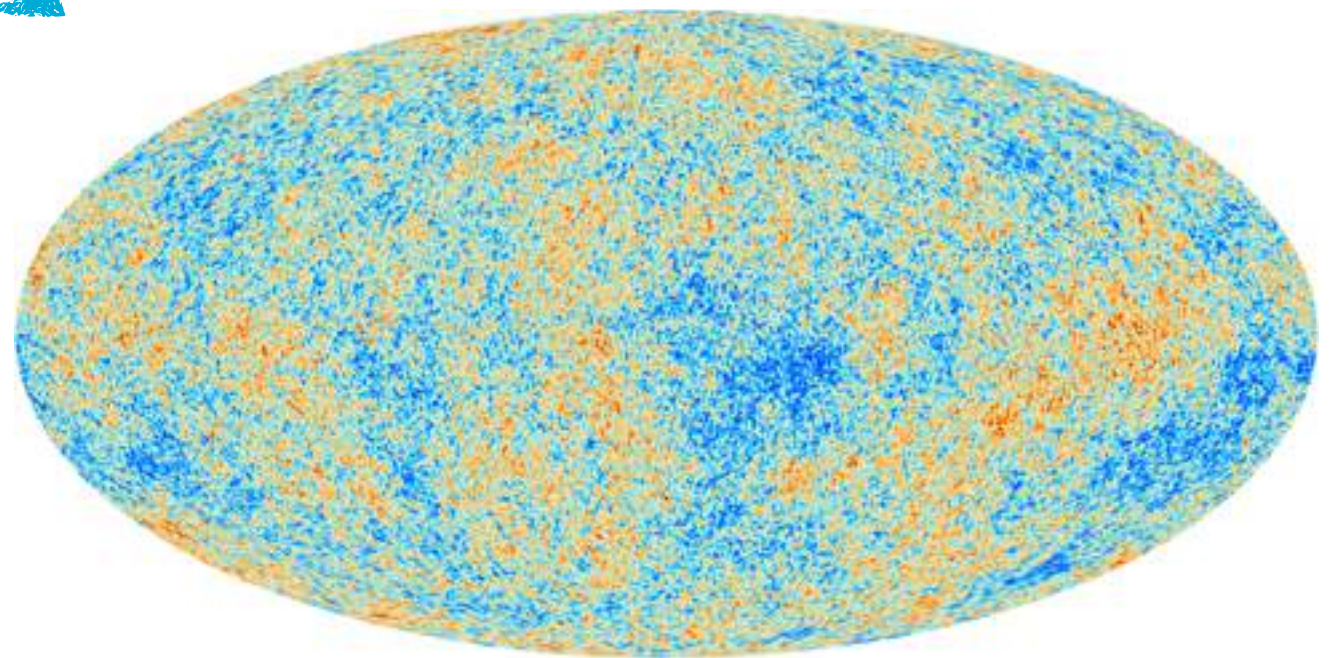
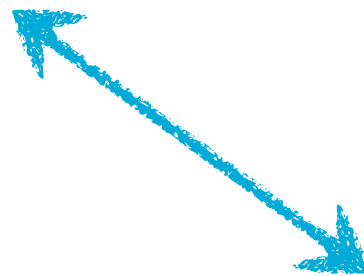
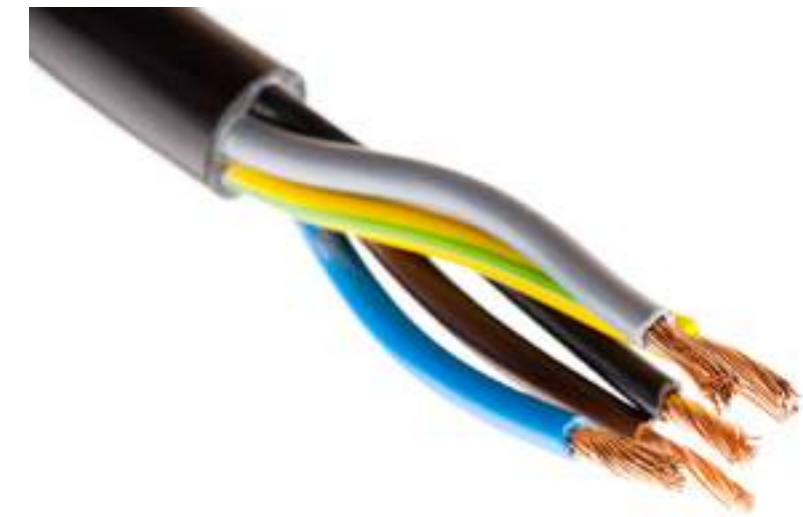
a statistical approach?

- observations: early universe is simple
- theory: not so much ...
- **coarse grained view ?**
- **calculational tools ?**



inspiration from disordered wires

MA & Baumann 2015



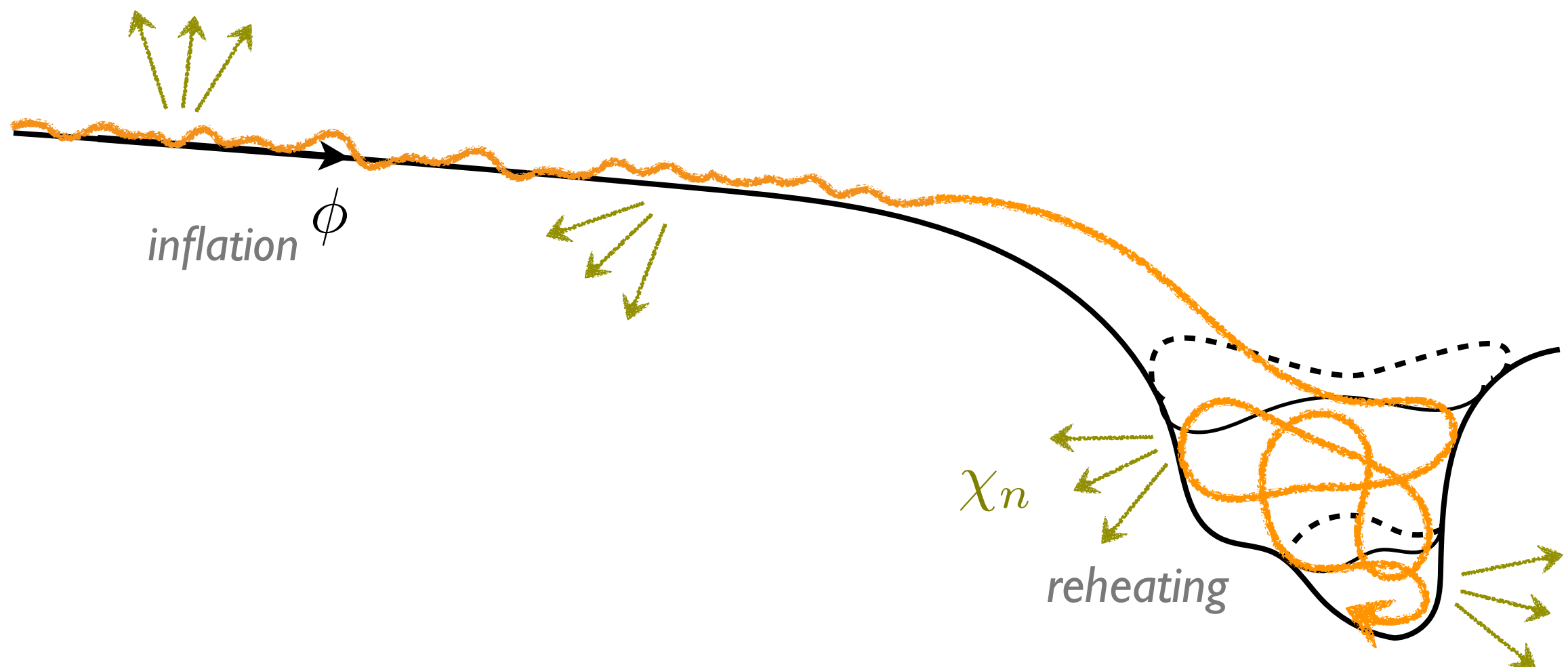


the framework



multifield inflation/reheating

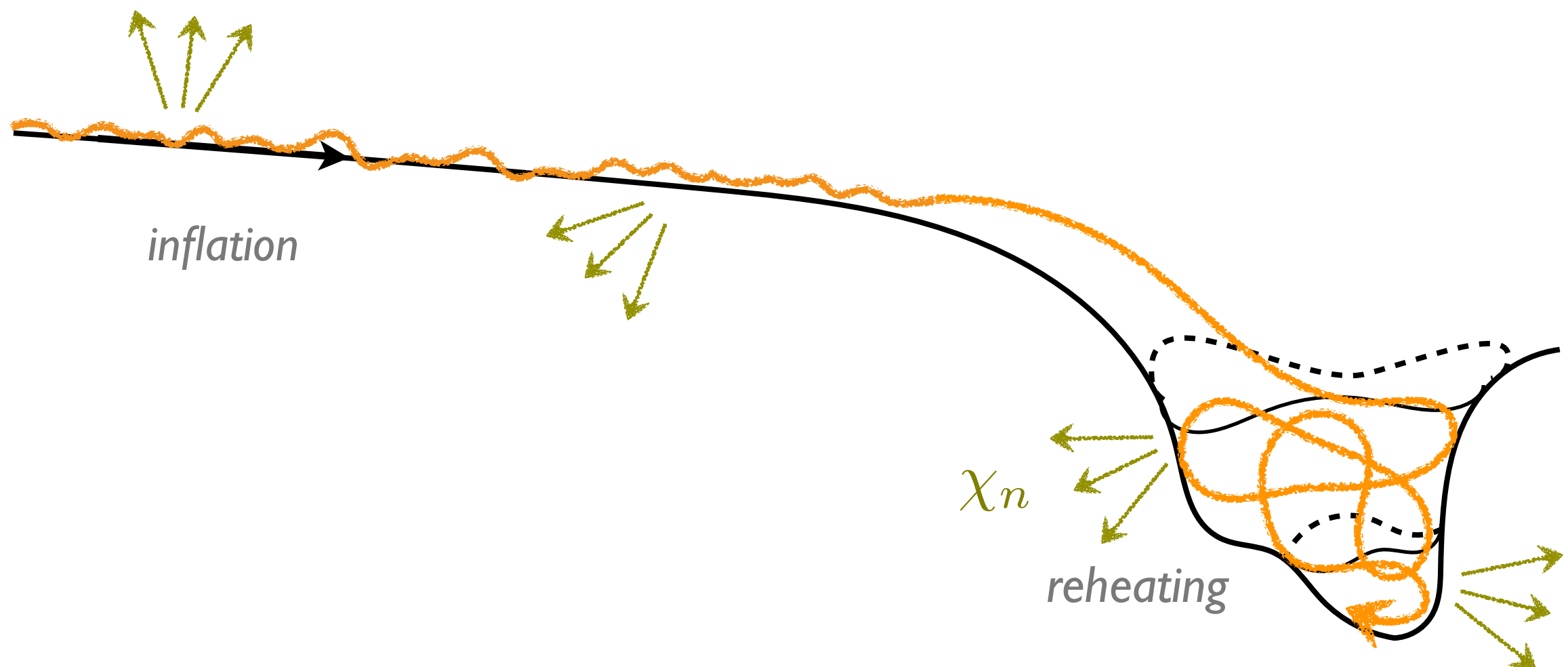
$$S = \int d^4x \sqrt{-g} \left[\frac{M_{\text{pl}}^2}{2} R - \frac{1}{2} G_{ab}(\phi^c) \partial^\mu \phi^a \partial_\mu \phi^b - V(\phi^c) + \dots \right]$$



focus on perturbations

mode functions in Fourier space

$$\left(\frac{d^2}{d\tau^2} + \omega_I^2 \right) \chi_k^I(\tau) + \sum_{J=1}^{N_f} m_{IJ}^s(\tau) \chi_k^J(\tau) = 0,$$
$$\omega_I^2(k) = k^2 + m_I^2,$$

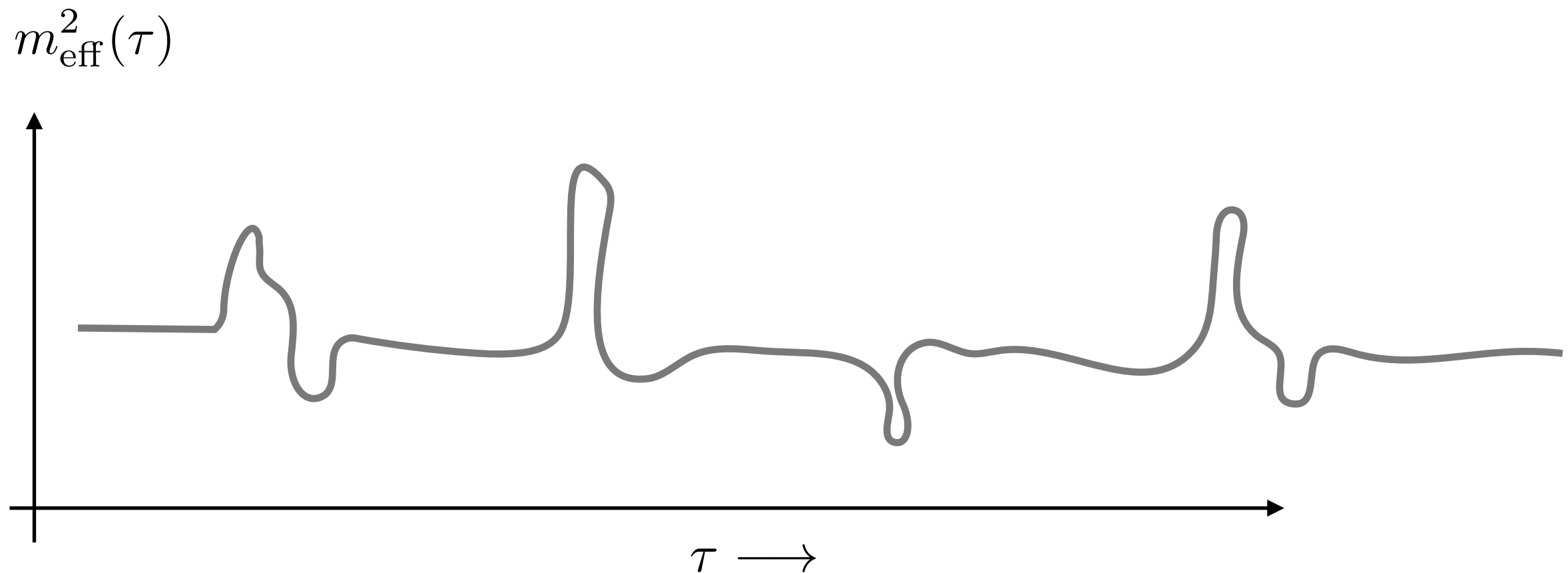


complexity in the “effective mass”/ interactions

simplified version!

$$\ddot{\chi}_k(\tau) + [k^2 + m_{\text{eff}}^2(\tau)] \chi_k(\tau) = 0$$

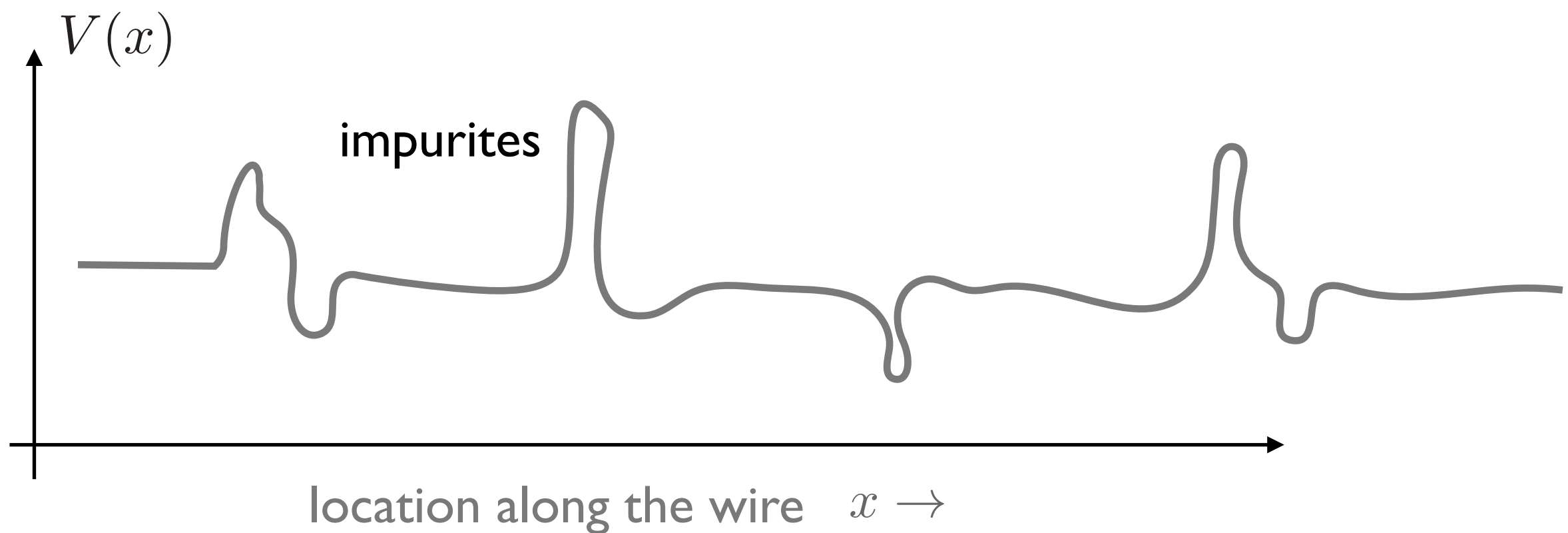
$$m_{\text{eff}}^2(\tau) = -\frac{\ddot{a}(\tau)}{a(\tau)} + a^2(\tau)m_\varphi^2 + a^2(\tau)g^2(\phi(\tau) - \phi_*)^2 + \dots$$



similar problem seen before ...

simplified version!

$$\frac{d^2\psi}{dx^2} + [k^2 - V(x)] \psi = 0$$

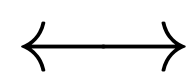


complexity in time
cosmology



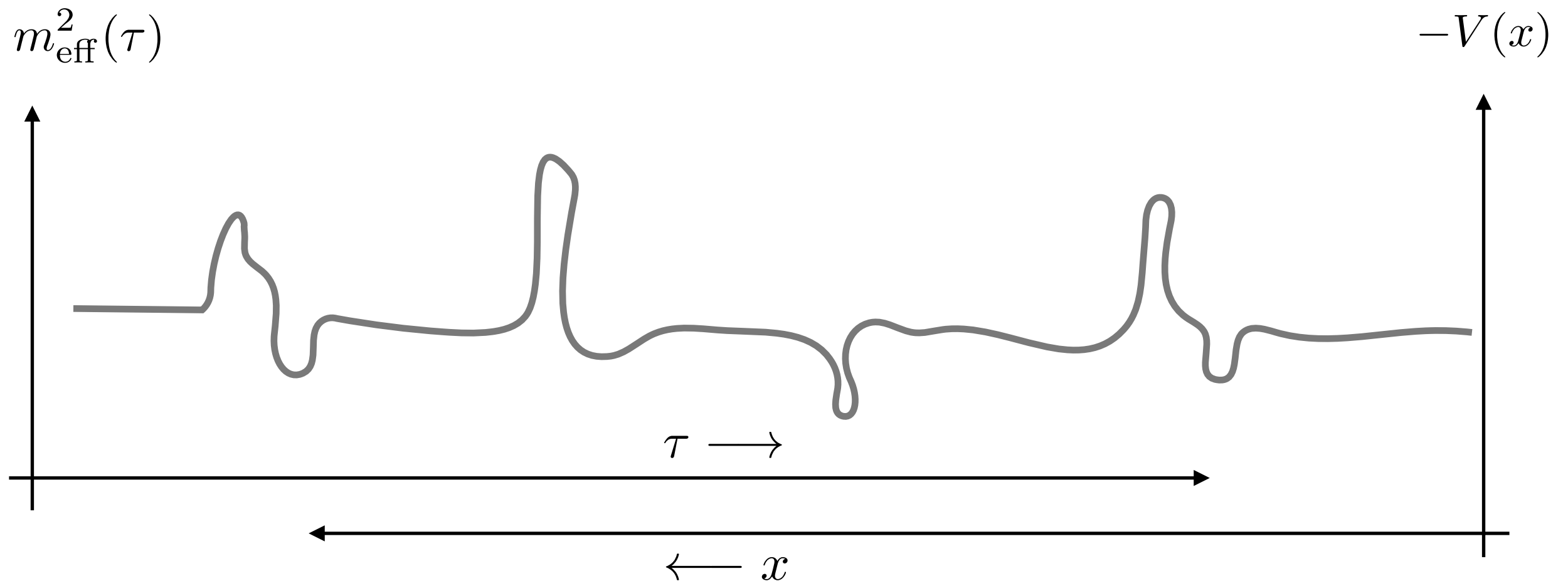
complexity in space
wires

$$\ddot{\chi}_k(\tau) + [k^2 + m_{\text{eff}}^2(\tau)] \chi_k(\tau) = 0$$



$$\psi''(x) + [k^2 - V(x)] \psi(x) = 0$$

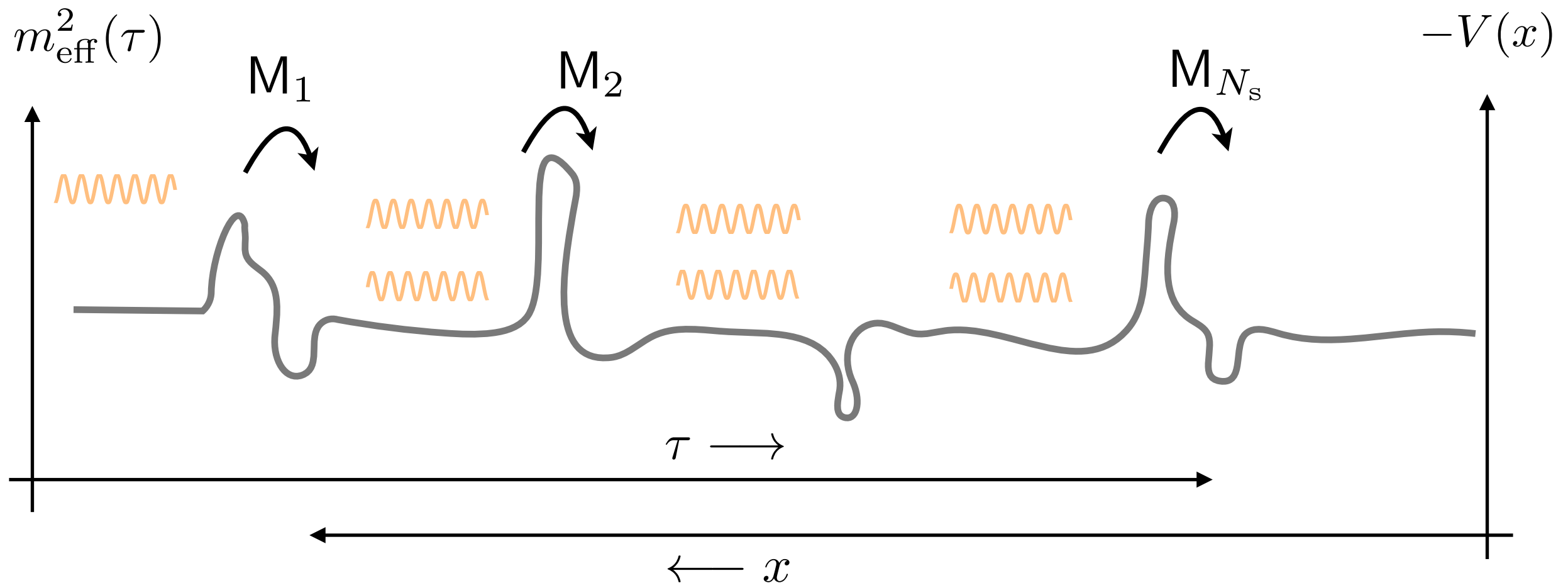
simplified version!



solve using scattering matrices !

$$\ddot{\chi}_k(\tau) + [k^2 + m_{\text{eff}}^2(\tau)] \chi_k(\tau) = 0 \longleftrightarrow \psi''(x) + [k^2 - V(x)] \psi(x) = 0$$

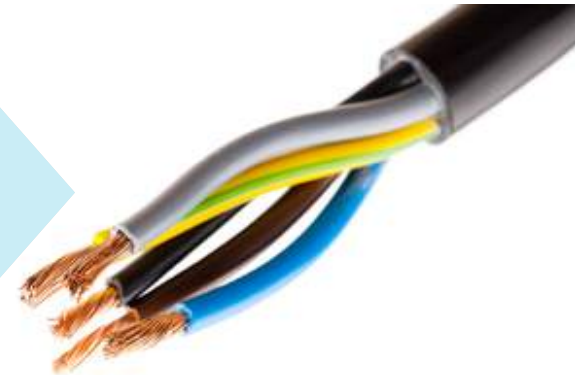
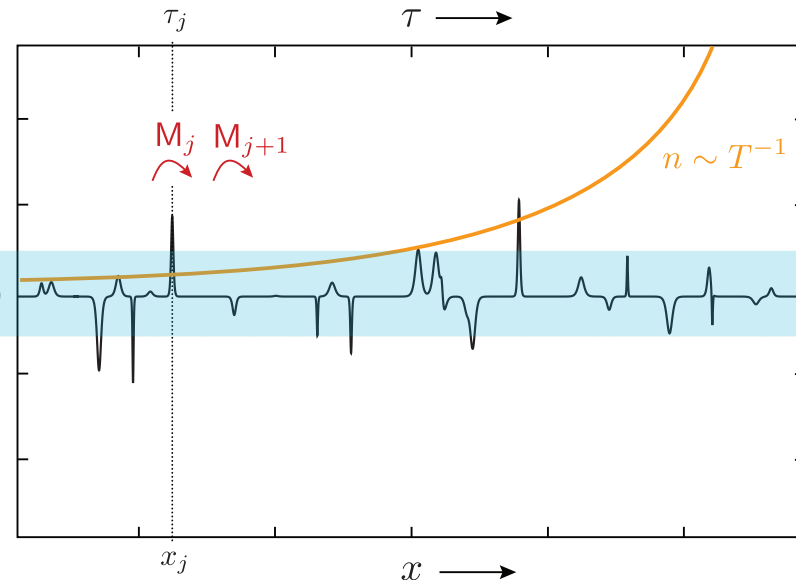
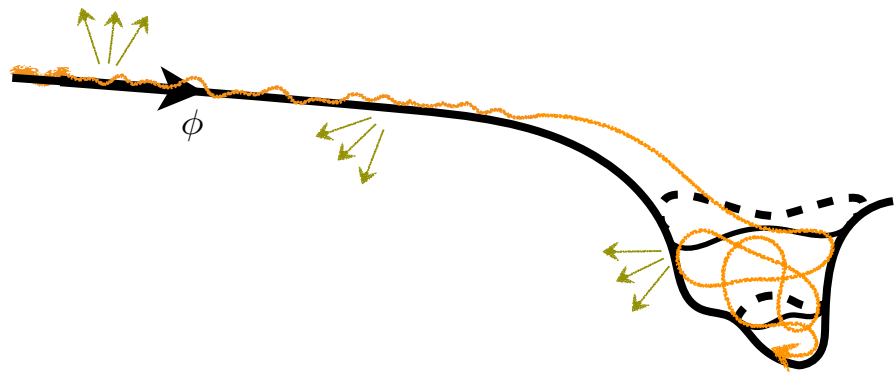
simplified version!



temporal complexity



spatial complexity



$$\frac{d^2 \chi_k}{d\tau^2} + [k^2 + m^2(\tau)] \chi_k = 0$$

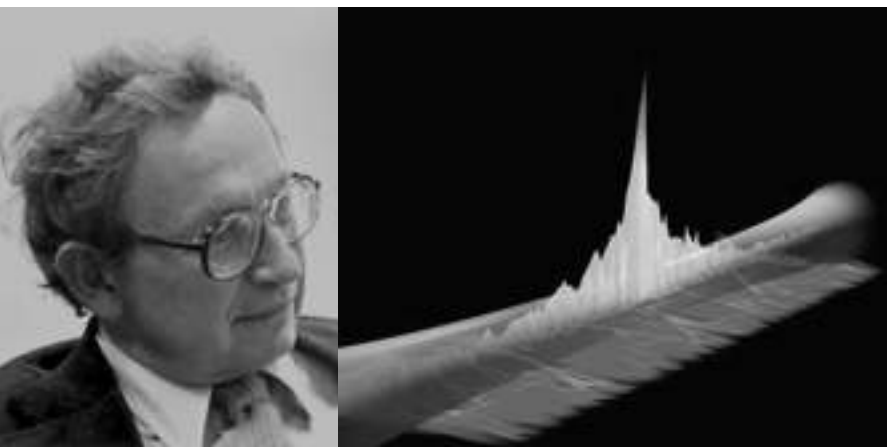
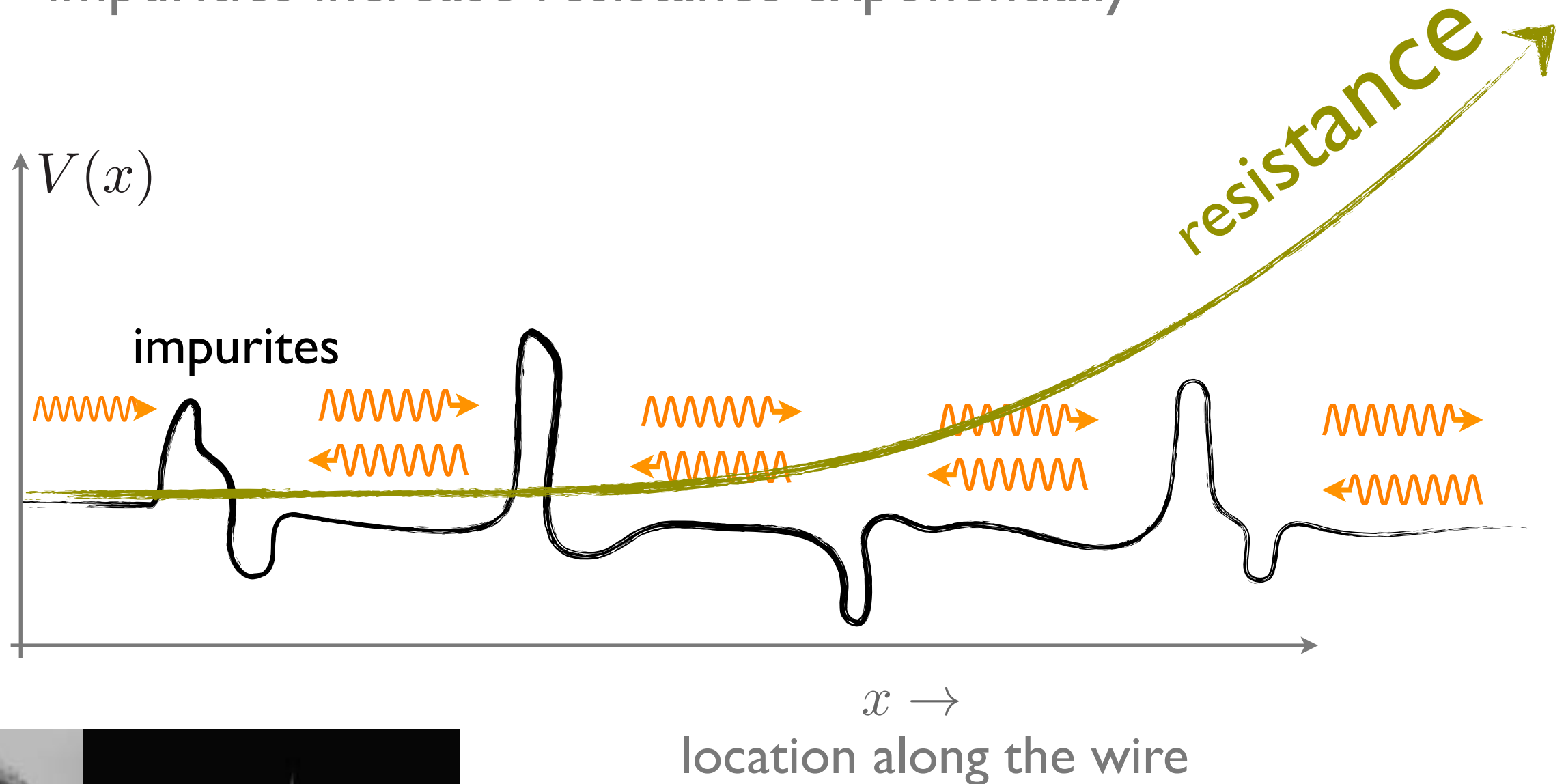
$$\frac{d^2 \psi}{dx^2} + [k^2 - V(x)] \psi = 0$$

- time
- number of particles
- complicated temporal behavior

- position along wire
- resistance
- impurities in wires

universal behavior: Anderson Localization!

- impurities increase resistance exponentially

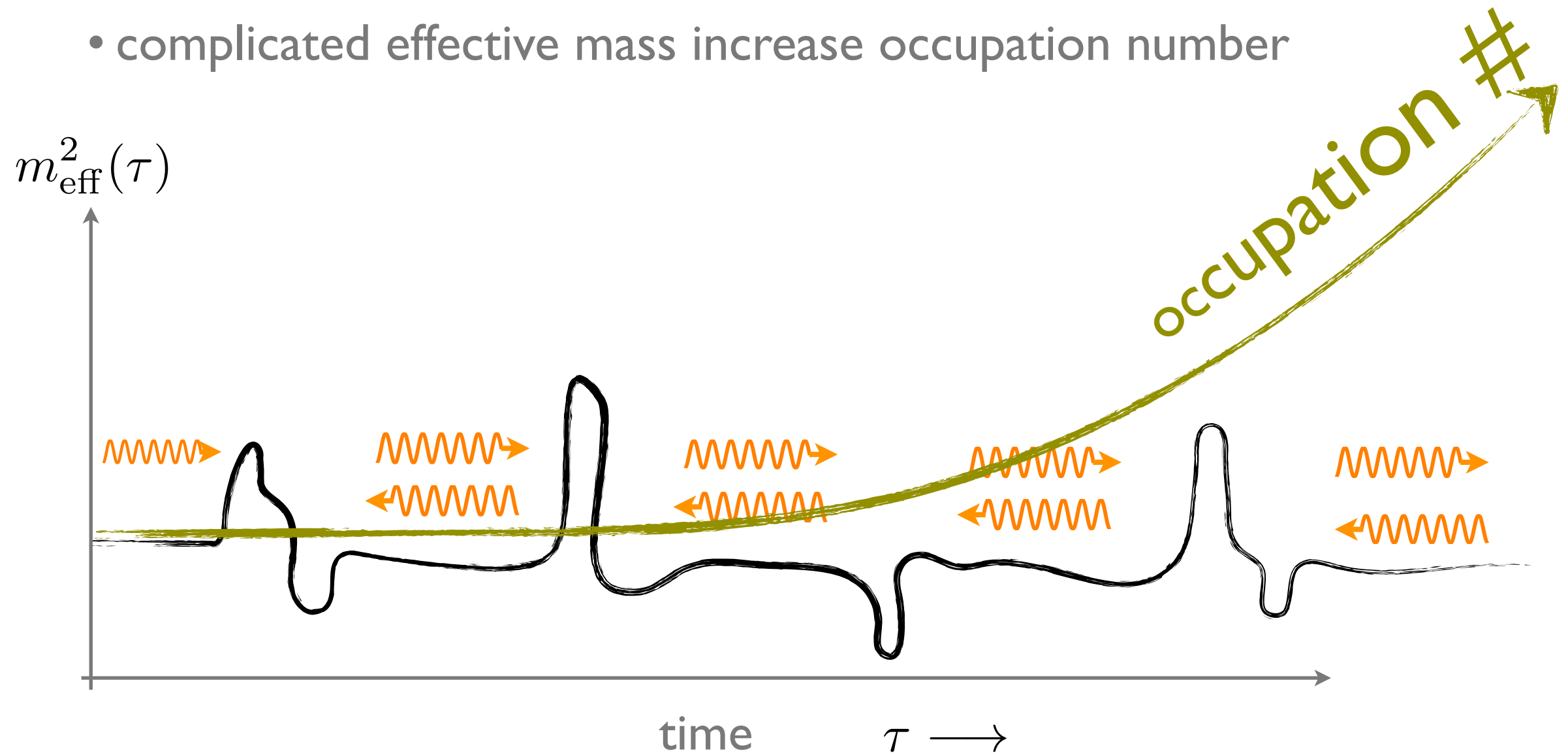


Anderson 1957

at low temperatures, one dimensional wires are insulators

universal behavior: exponential increase in occupation number

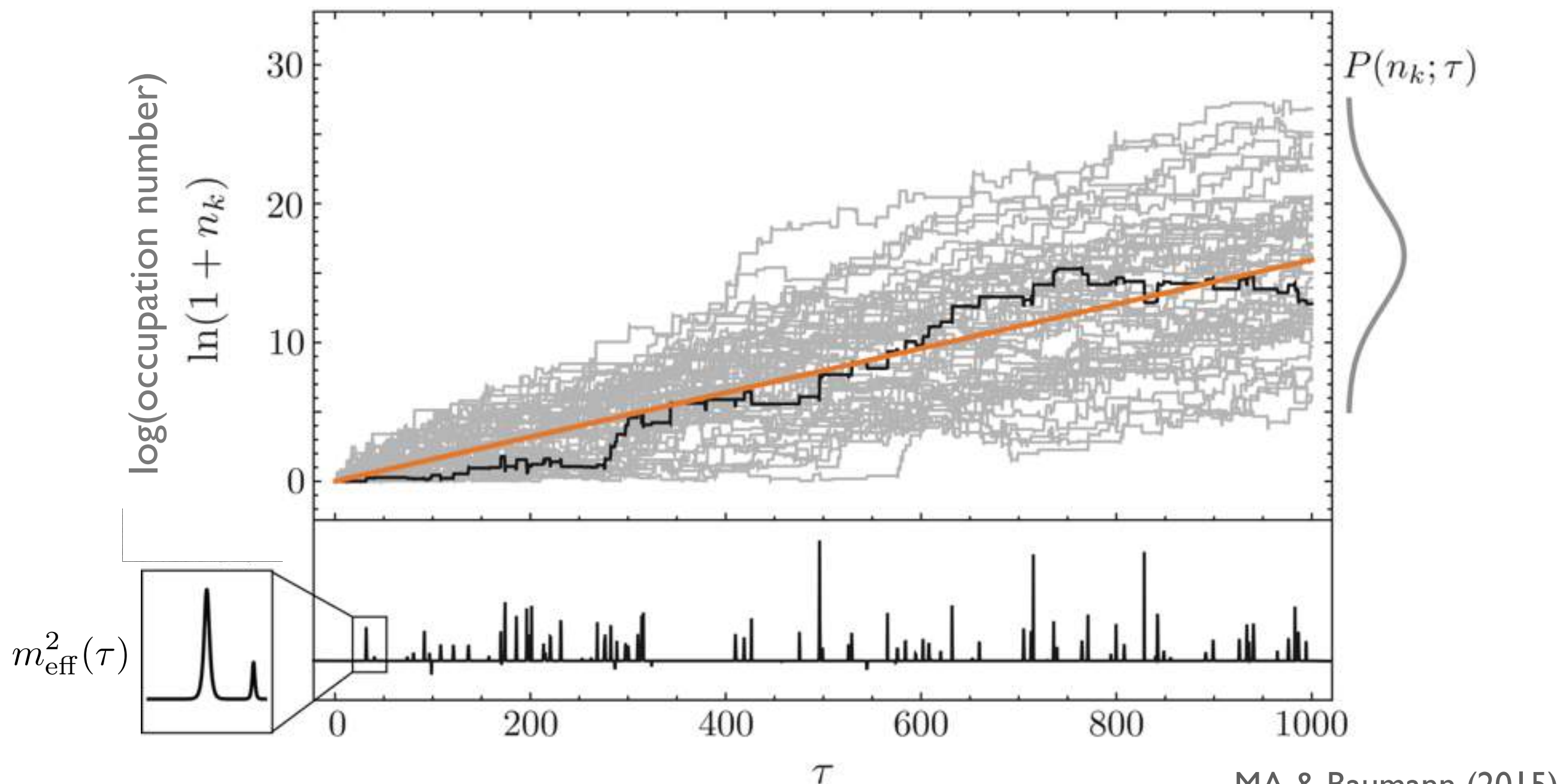
- complicated effective mass increase occupation number



understood in terms of **Bose enhancement**

occupation number performs a drifted random walk

$$n(k, \tau) = \frac{1}{2\omega_k} (|\dot{\chi}_k|^2 + \omega_k^2 |\chi_k|^2)$$



multifield Fokker Planck equation

joint probability for occupation numbers satisfies the **a Fokker Planck-like** equation:

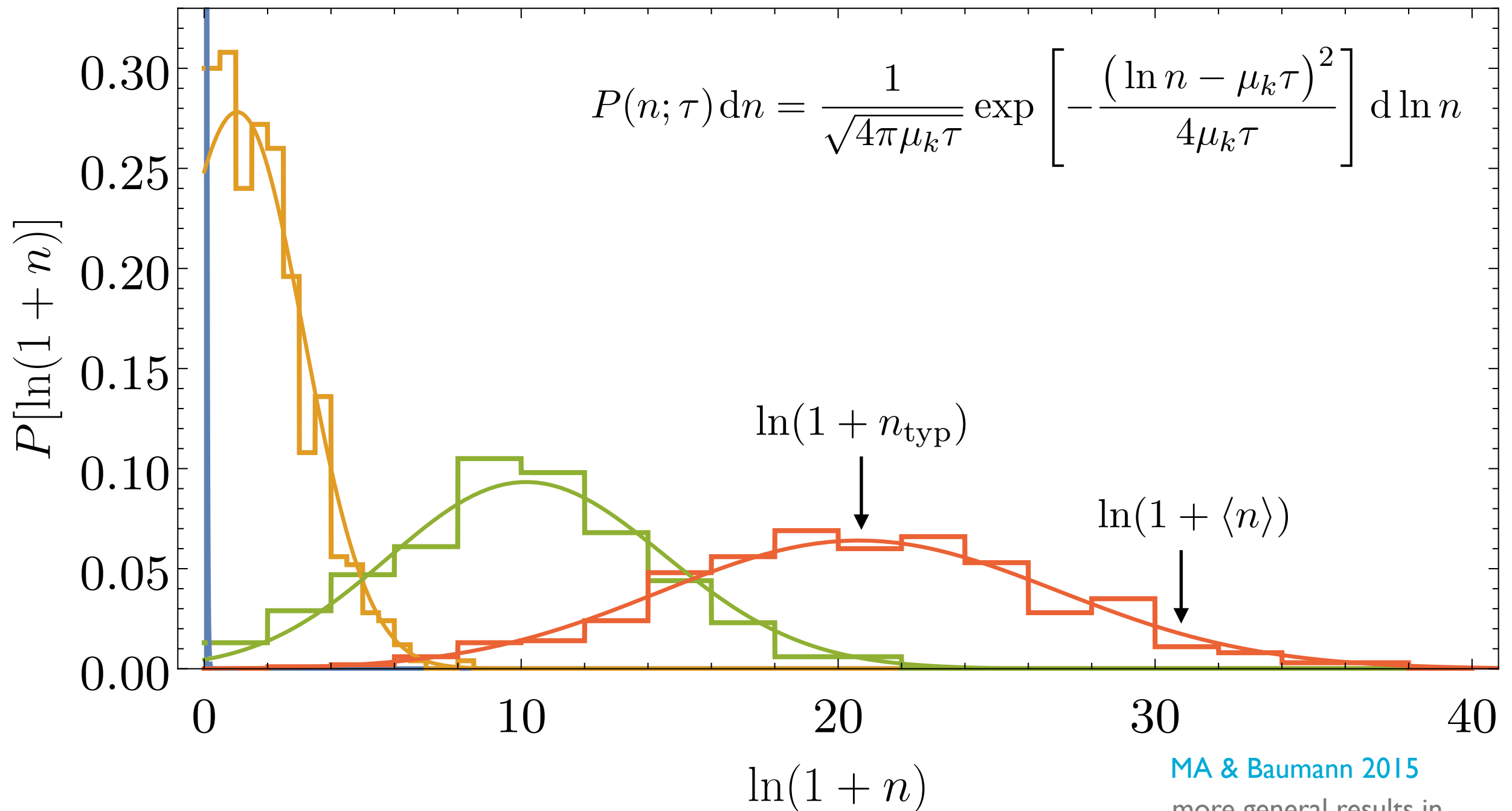
Dokhorov, Mello, Pereyra & Kumar = DMPK eq.

$$\begin{aligned} \frac{1}{\mu_k} \frac{\partial}{\partial \tau} P(n_a; \tau) = & \sum_{a=1}^{N_f} \left[(1 + 2n_a) + \frac{1}{N_f + 1} \sum_{b \neq a} \frac{n_a + n_b + 2n_a n_b}{n_a - n_b} \right] \frac{\partial P}{\partial n_a} \\ & + \frac{2}{N_f + 1} \sum_{a=1}^{N_f} n_a (1 + n_a) \frac{\partial^2 P}{\partial n_a^2} \end{aligned}$$

MA & Baumann 2015

For more general results in
beyond statistical isotropy, see
MA, Garcia, Xie and Wen 2017

solution: “universal” distributions



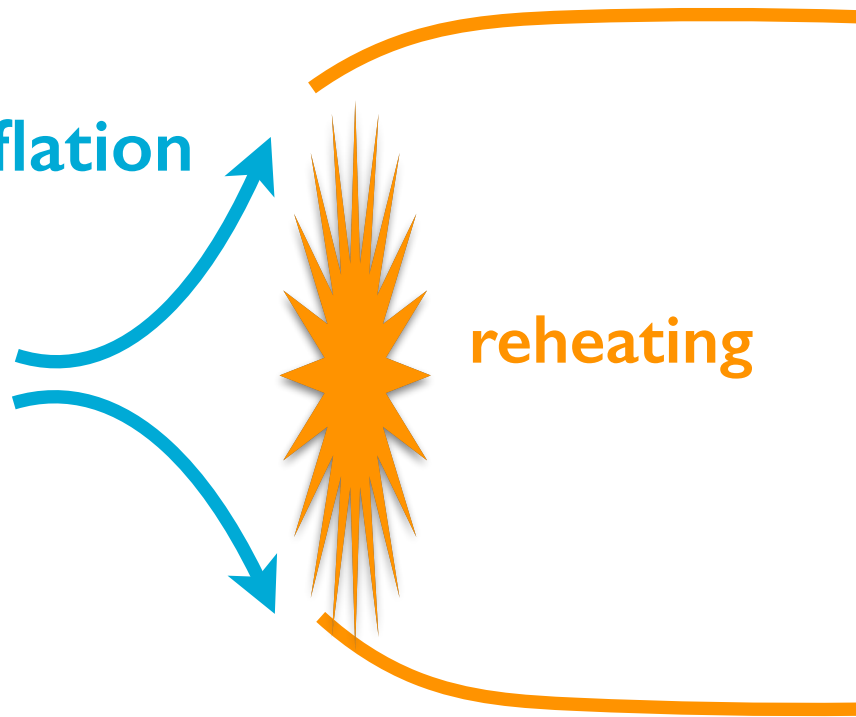
MA & Baumann 2015

more general results in
MA, Garcia, Xie and Wen 2017

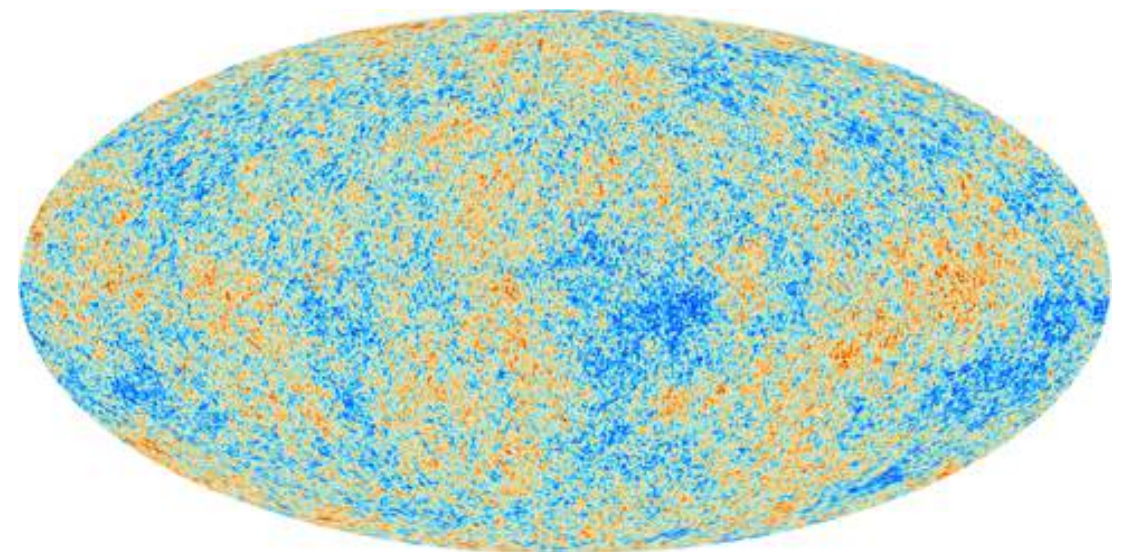
MA, Carlsten, Garcia, Green, Baumann ...

for superhorizon scales, use field amplitude rather than occupation number

inflation



applications



WORK IN
PROGRESS

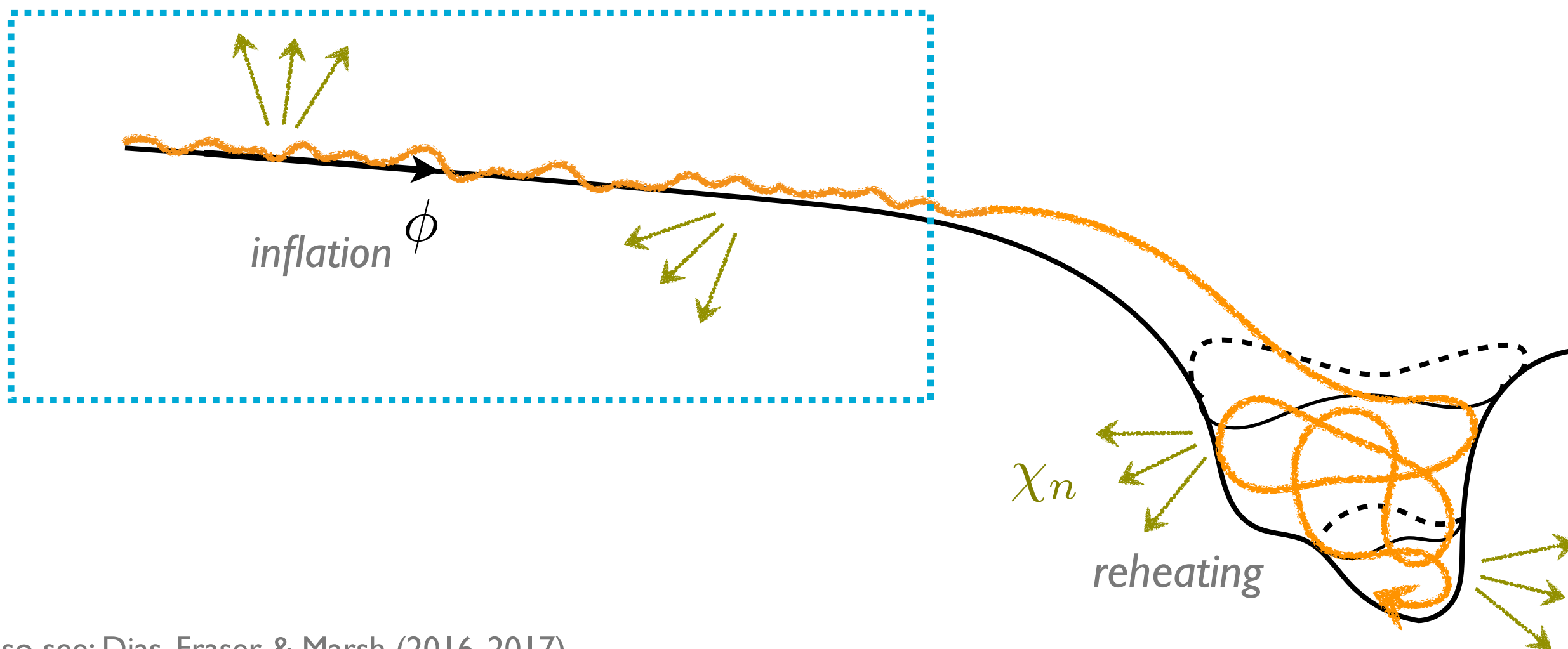
applications: inflation

MA, Garcia, Baumann, Carlsten, Chia & Green

background dynamics \rightarrow particle production \leftrightarrow curvature fluctuations

$$\langle n_{k_1} n_{k_2} \dots \rangle$$

$$\langle \zeta_{k_1} \zeta_{k_2} \dots \rangle$$



also see: Dias, Fraser & Marsh (2016, 2017)

WORK IN
PROGRESS

combine particle production with driving and dissipation

background dynamics \rightarrow particle production \leftrightarrow curvature fluctuations

$$\langle n_{k_1} n_{k_2} \dots \rangle$$

$$\langle \zeta_{k_1} \zeta_{k_2} \dots \rangle$$

$$\ddot{\pi}_k + [3H + \mathcal{O}_d] \pi_k + \frac{k^2}{a^2} \pi_k = \mathcal{O}_s(\langle \chi \chi \dots \rangle_k)$$

$$\zeta_k = -H \pi_k$$

dissipation

driving

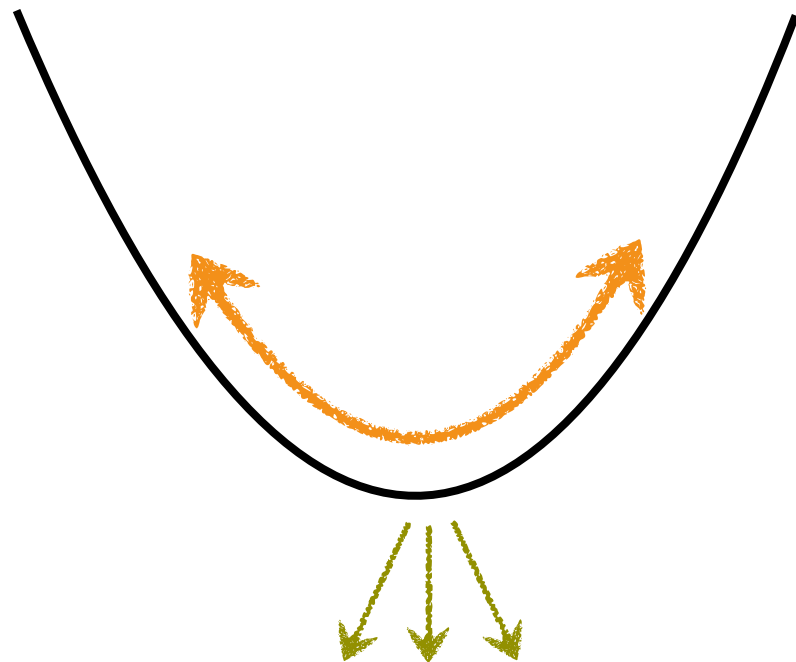
MA, Garcia, Baumann, Carlsten, Chia & Green

Green, Horn, Senatore, and Silverstein (2009),
Nacir, Porto, Senatore, and Zaldarriaga (2012),
Flauger, Mirbabayi, Senatore, Silverstein (2016),
Chen, Namjoo and Wang (2015,16), Dias, Fraser
& Marsh (2016, 17)

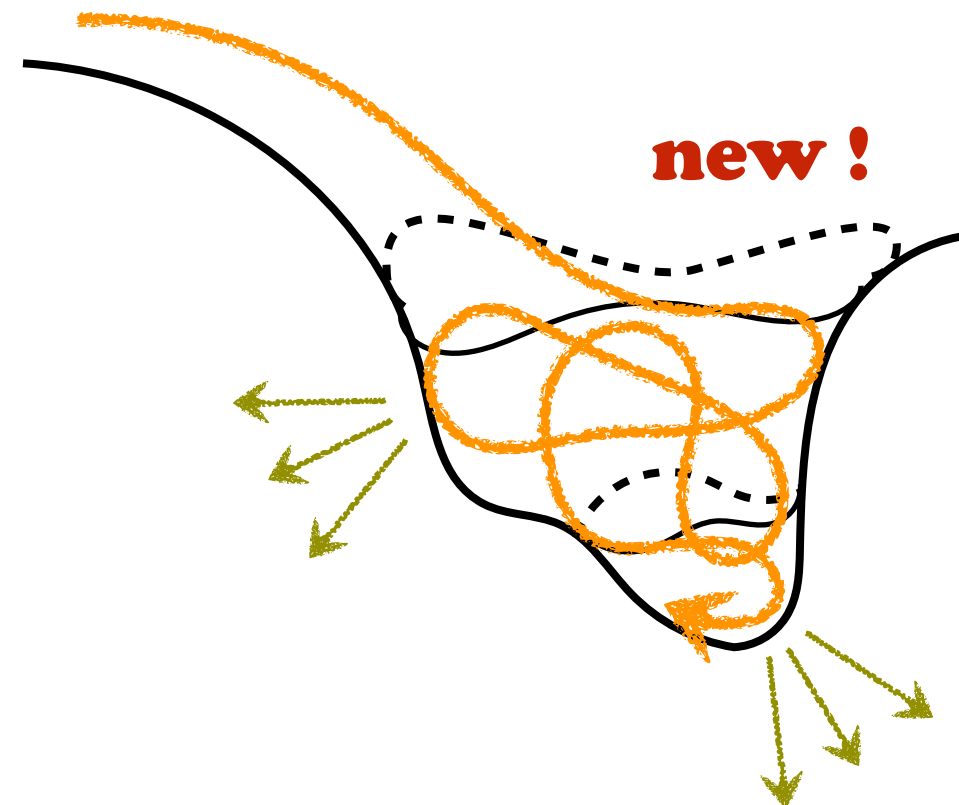


**WORK IN
PROGRESS**

applications : reheating



multichannel — multifield — statistical



**model-insensitive description of a
complicated reheating process.**

for example:

Shtanov, Traschen & Brandenberger (1995)

Kofman, Linde & Starobinsky (1997)

Zanchin et. al (1998) & Bassett (1998) [with noise]

Barnaby, Kofman & Braden et. al 2010 [quasiperiodic]

Giblin, Nesbit, Ozsoy, Sengor & Watson (2016-17)

MA, Garcia & Shen



end of inflation



SIMPLE

- single field
- non-trivial dynamics
- eq. of state + gravitational waves

detailed models

- universal dynamics independent of details
- Fokker-Plank/random matrix theory
- statistical properties of correlation functions



COMPLEX

end of inflation



SIMPLE

- single field
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COMPLEX



non-perturbative dynamics to solve problems

- Higgs vacuum stability and inflation [East, Kearny, Shakya, Yoo, Zurek]
- Higgs fine tuning problem ? [MA, Fan, Lozanov, Reece (in progress)]
- matter-antimatter asymmetry ? [Lozanov & MA 2014, Hertzberg & Karouby 2014]
- primordial seeds for cosmological magnetic fields ? [MA & Lozanov 2016]
- primordial black holes ? [MA & Lozanov (in progress)]

non-perturbative dynamics can lead to model specific observational signatures

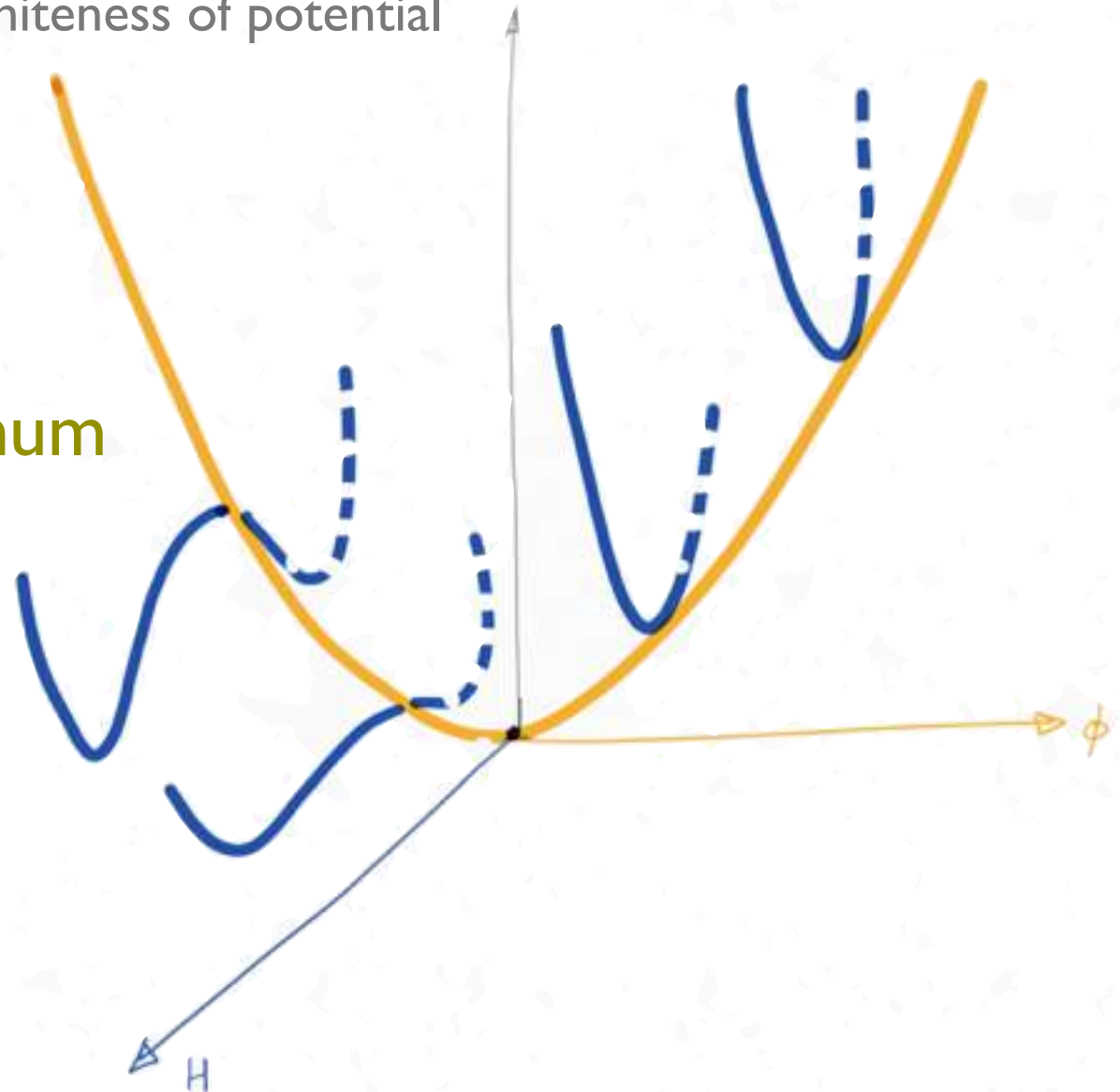
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- primordial black holes ? [MA & Lozanov (in progress)]

inflaton (modulus) - Higgs system

$$V(\phi, H) = -\mu^2 H^\dagger H + \lambda (H^\dagger H)^2 + \frac{M^2}{f} \phi (H^\dagger H - v^2)$$

$$b \equiv \frac{M^4}{2\lambda m_\phi^2 f^2} \leq 1 \quad \text{for positive definiteness of potential}$$

arrange the potentials to yield
small Higgs masses at the global minimum



implications?

$$V(\phi, H) = -\mu^2 H^\dagger H + \lambda (H^\dagger H)^2 + \frac{M^2}{f} \phi (H^\dagger H - v^2)$$

$$b \equiv \frac{M^4}{2\lambda m_\phi^2 f^2} \leq 1$$

for positive definiteness of potential

$$b \equiv \frac{M^4}{2\lambda m_\phi^2 f^2} \rightarrow 1$$

back-reaction
efficiency parameter

- inhomogeneous fragmentation of Higgs-modulus system with $\rho_H \sim \rho_\phi$
- generation of stochastic g-waves
- non-trivial equation of state

end of inflation



SIMPLE

- single field
- non-trivial dynamics
- eq. of state + gravitational waves



detailed models

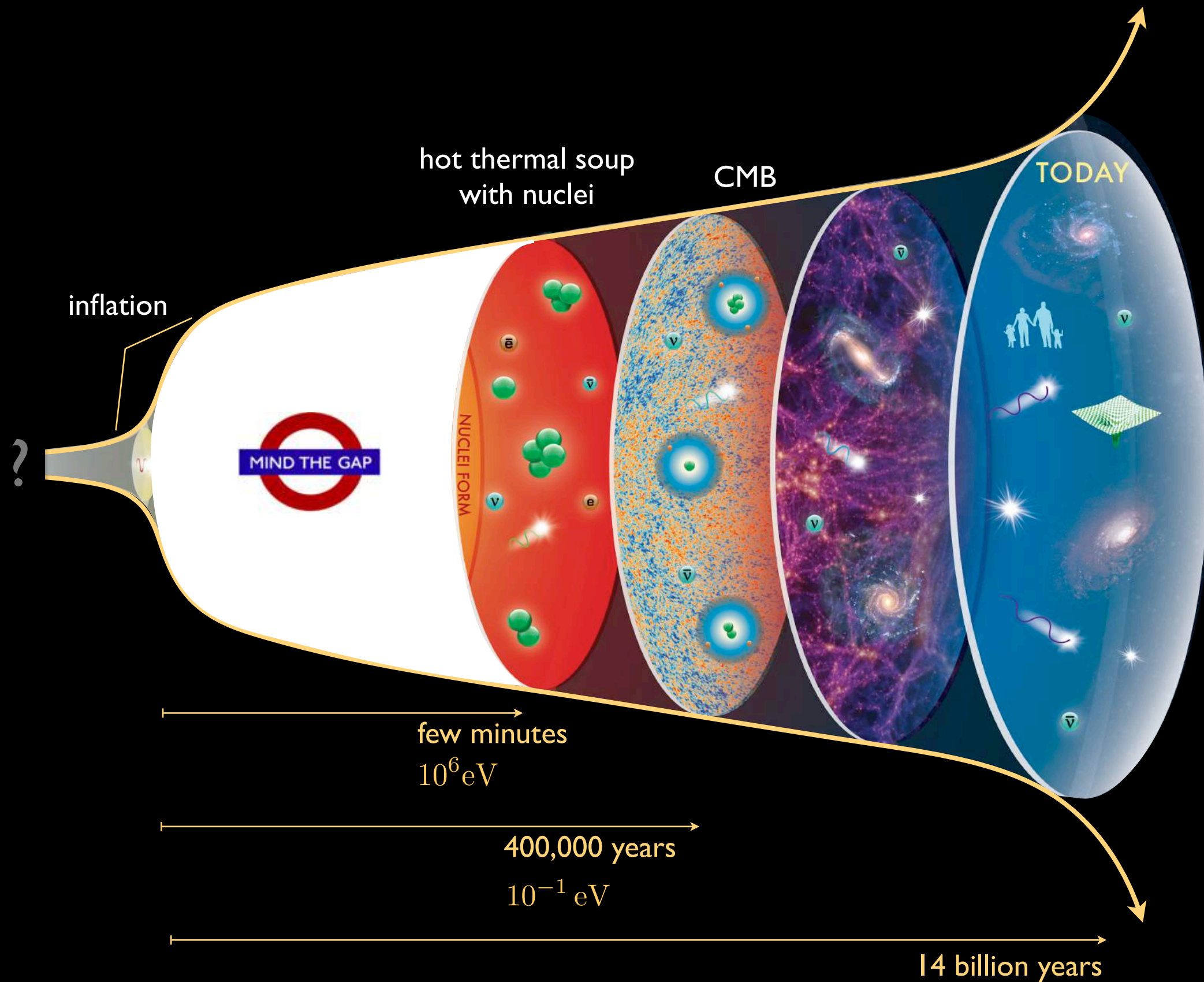
- hierarchy problem
- matter-antimatter asymmetry
- primordial magnetic fields

- universal dynamics independent of details
- Fokker-Plank/Random Matrix theory
- statistical properties of correlation functions

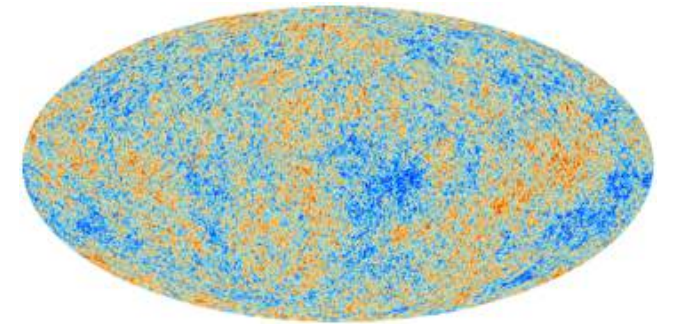
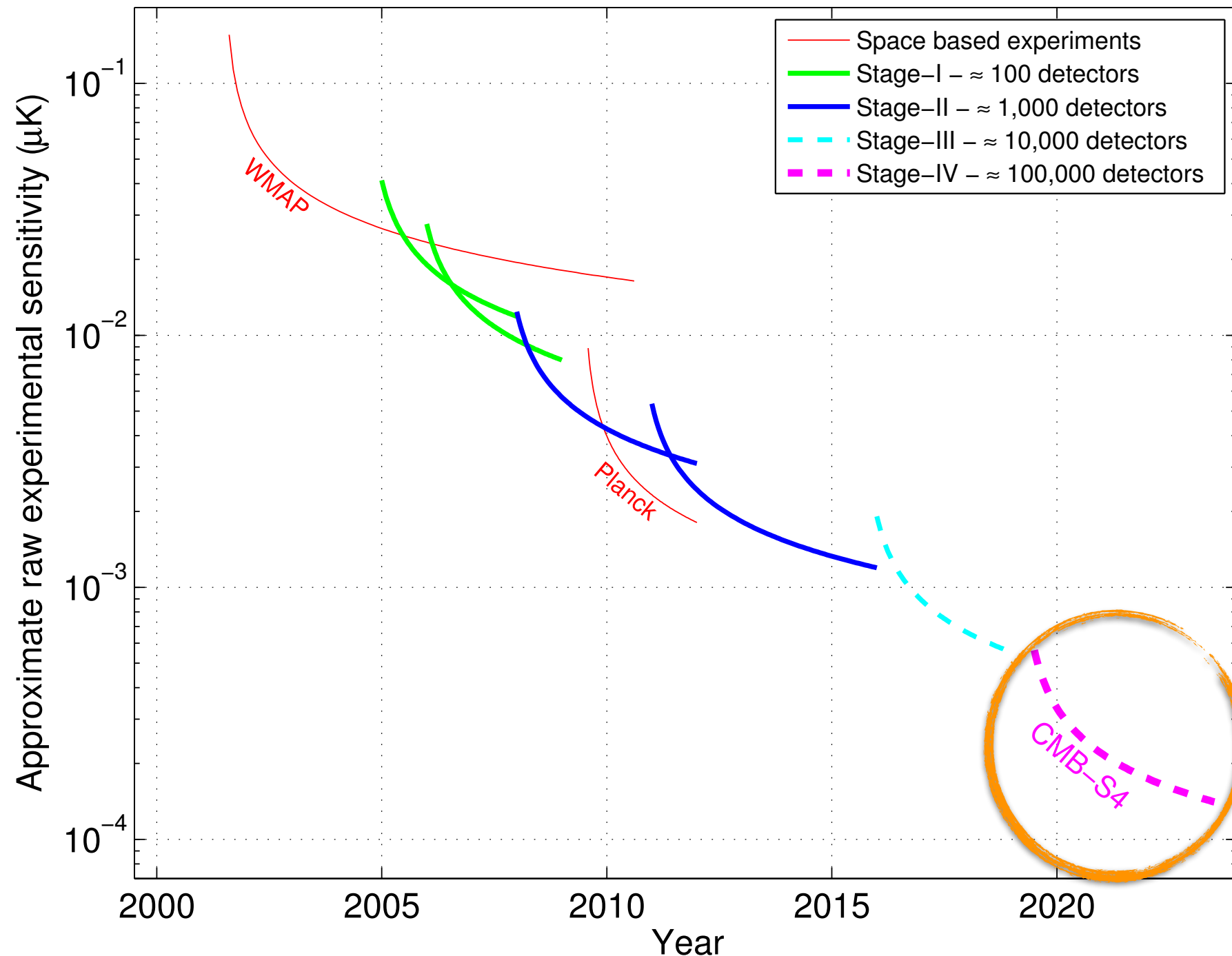


COMPLEX

What's Next? — upcoming observations



significant improvement expected

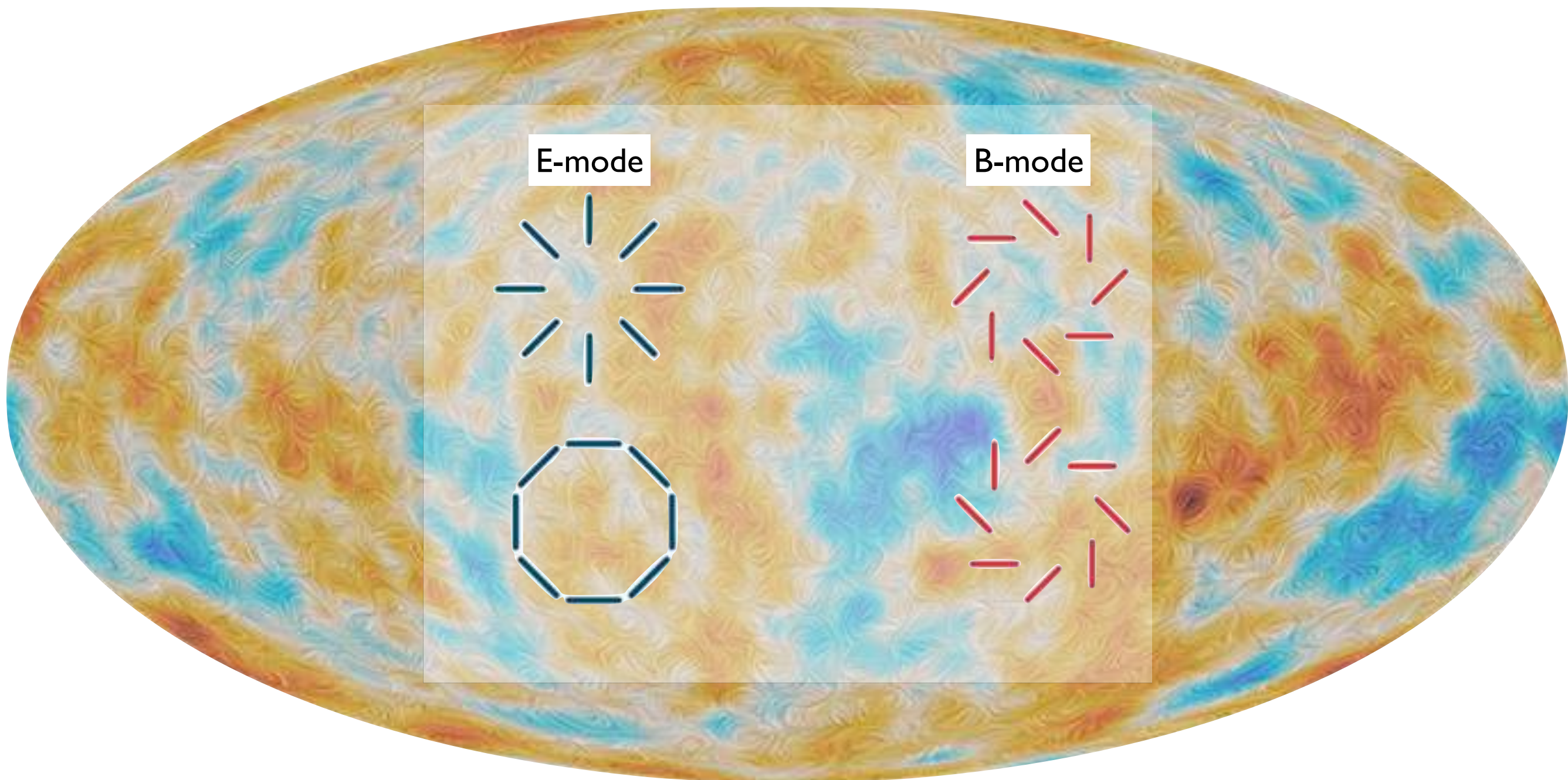


ground based

- Chilean Atacama Plateau
- South Pole
- Northern Hemisphere ?

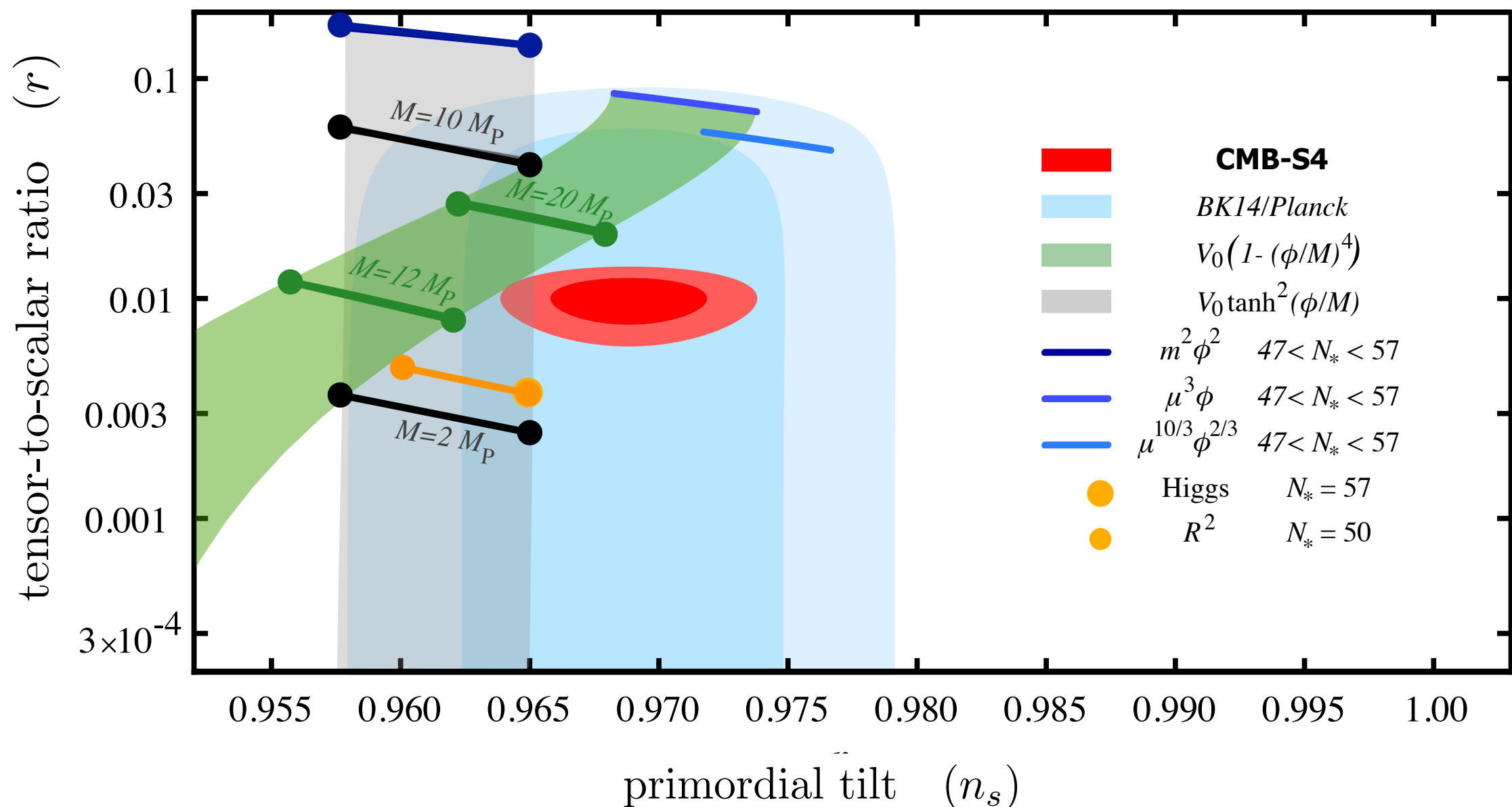
additional information in CMB ?

CMB is polarized



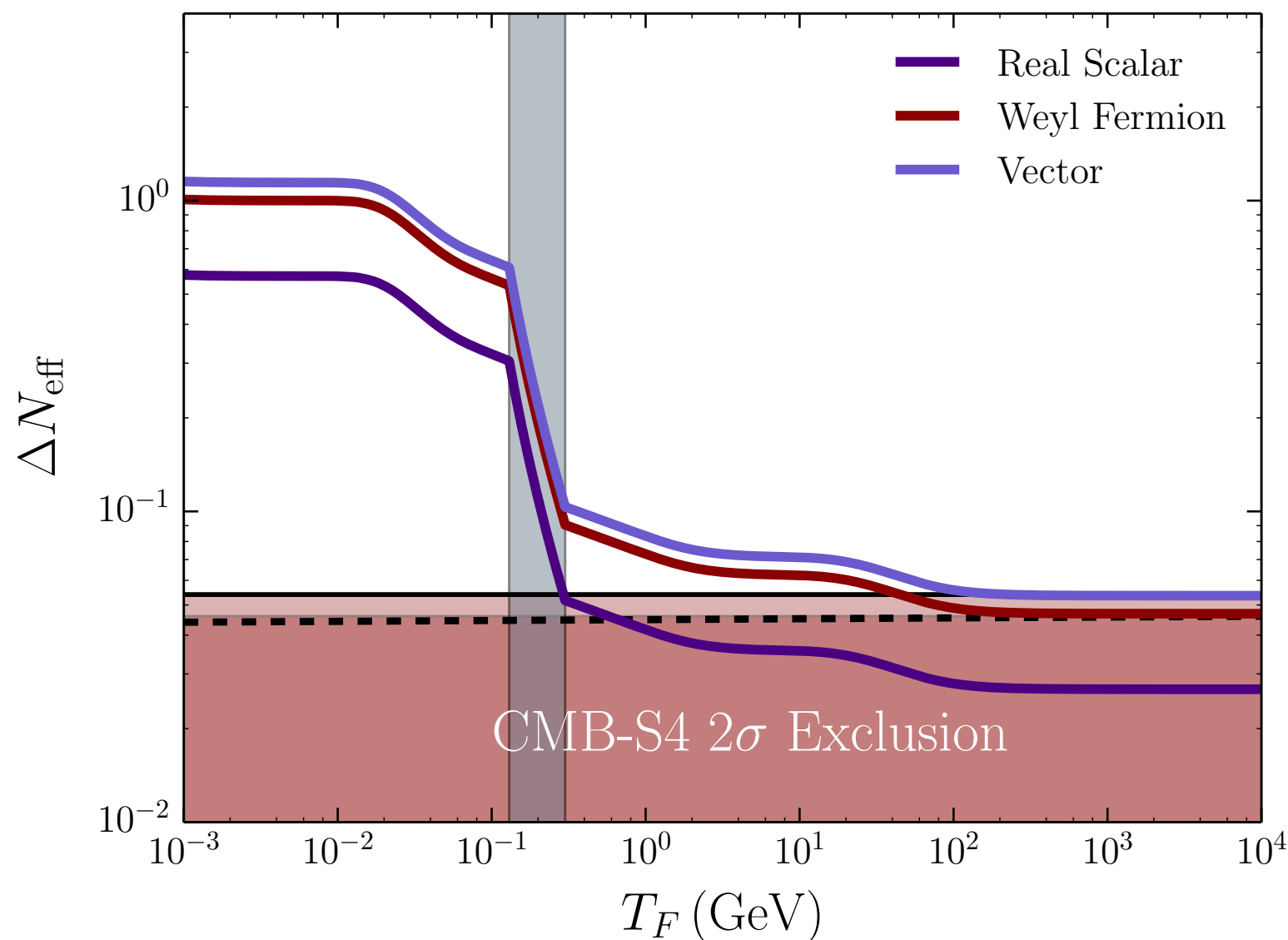
improvement expected necessitates better theoretical understanding

Will also require better understanding of the aftermath of inflation



constraints on additional light species

will also require better understanding of the aftermath of inflation



will be able to constrain the the effective number of additional light species, beyond neutrinos!

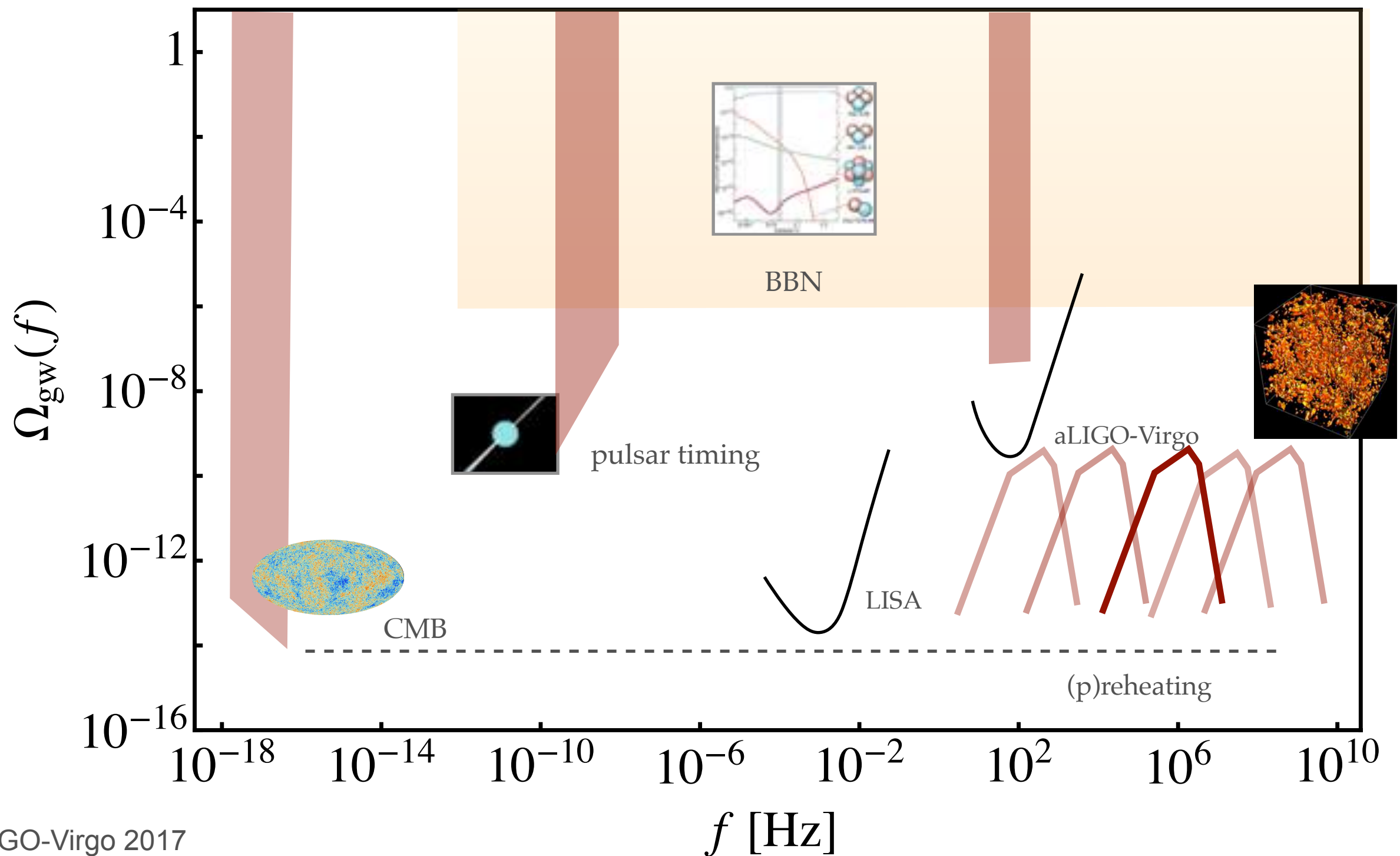
large scale structure

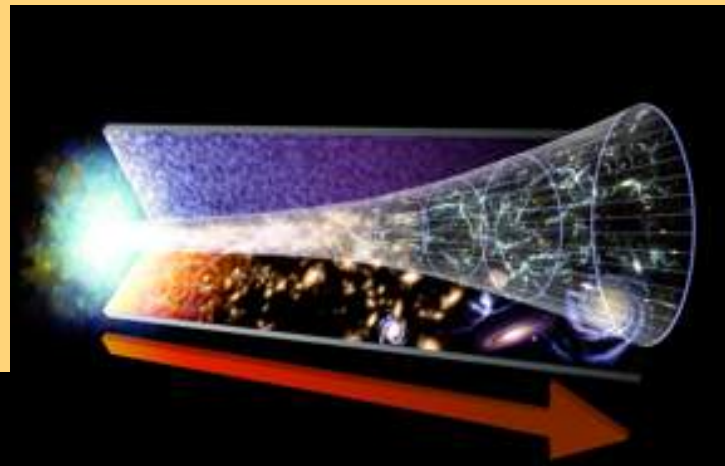
- 3D, more modes — lensing, spectroscopic, 21 cm etc.)



	LSST	DESI	Euclid	SPHEREx	CHIME
Survey type	photo	spectro	photo+spectro	low-res spectro	21-cm
Ground or space	ground	ground	space	space	ground
Previous surveys	CFHTLS, DES, HSC	BOSS, eBOSS, PFS	no direct precursor	PRIMUS, COMBO-17, COSMOS	GBT HIM
Survey start	2020	2020	2018	2020	2016
Redshift-range	$z < 3$ (1% sources above 3)	$z < 1.4$, $2 < z < 3.5$ (Ly α)	$z < 3$	$z < 1.5$	$0.75 < z < 2.5$
Survey area [deg ²]	20k	14k	15k	40k	20k
Approximate number of objects	2×10^9 (WL sources)	22×10^6 gal., $\sim 2.4 \times 10^5$ QSOs	40×10^6 redshifts, 1.5×10^9 photo-zs	15×10^9 pixels	10^7 pixels
Galaxy clustering	✓✓◇	✓	✓	✓	✓
Weak lensing	✓		✓		✓
RSD		✓	✓	✓✓	✓✓
Multi-tracer	✓✓	✓✓	✓✓	✓	

stochastic gravitational waves





summary

◎ **inflation** - origin of density perturbations

◎ **aftermath** - *gap* in our cosmic history



- simple models (solitons, eq. of state, g-waves)
- complex (universal behavior)
- model-specific (hierarchy problem etc.)

◎ **what's next** — light species ? g-waves ? polarization ? relics ?



Assistant Professor in Theoretical Astro-Particle Physics/Cosmology Position

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inflation ends, what's next ?

